# Point Processes and Multiple SLE/GFF Coupling <sup>∗</sup>

Makoto Katori †

12 August 2022

#### Abstract

In the series of lectures, we will discuss probability laws of random points, curves, and surfaces.

Starting from a brief review of the notion of martingales, one-dimensional Brownian motion (BM), and the D-dimensional Bessel processes,  $BES_D$ ,  $D \geq 1$ , first we study Dyson's Brownian motion model with parameter  $\beta > 0$ , DYS<sub>β</sub>, which is a one-parameter family of repulsively interacting N Brownian motions on  $\mathbb{R}, N \in \mathbb{N} := \{1, 2, \ldots\}$ , and is regarded as multivariate extensions of BES<sub>D</sub> with the relation  $\beta = D - 1$ . In particular, the determinantal structures are proved for  $DYS<sub>2</sub>$ , which is realized as the eigenvalue process of Hermitian-matrix-valued BM studied in random matrix theory.

Next, using the reproducing kernels of Hilbert function spaces, the Gaussian analytic functions (GAFs) are defined on a unit disk  $\mathbb{D}$  and an annulus  $\mathbb{A}_q := \{z \in \mathbb{C} : q < |z| < 1\},\$  $q \in (0, 1)$ , which provide models of random surfaces. As zeros of the GAFs, determinantal point processes and permanental-determinantal point processes are obtained, which have symmetry and invariance associated with conformal transformations.

Then, the Schramm–Loewner evolution with parameter  $\kappa > 0$ , SLE<sub> $\kappa$ </sub>, is introduced, which is driven by a BM on  $\mathbb R$  and generates a family of conformally invariant probability laws of random curves on the upper half complex plane  $\mathbb{H}$ . We regard  $\text{SLE}_{\kappa}$  as a complexification of  $BES_D$  with the relation  $\kappa = 4/(D-1)$ .

The last topic of lectures is the construction of the multiple  $SLE_{\kappa}$ , which is driven by the N-particle DYS<sub>β</sub> on R and generates N interacting random curves in H. There, we define the Gaussian free field (GFF) and its generalization called the imaginary surface with parameter  $\chi$ , which are considered as the distribution-valued random fields on  $\mathbb{H}$ . Under the relation  $\chi = 2/\sqrt{\kappa} - \kappa/\sqrt{2}$ , we characterize the SLE/GFF coupling studied by Dubédat, Sheffield, and Miller by its temporal stationarity, and extend it to multiple cases. We prove that the multiple SLE/GFF coupling is established, if and only if the driving N-particle process on  $\mathbb R$  is identified with DYS<sub>β</sub> with the relation  $\beta = 8/\kappa$ . This relation between parameters is very simple, but highly nontrivial, since if we simply combine the relations  $\beta = D-1$  and  $\kappa = 4/(D-1)$  mentioned above, we will have a different relation  $\beta = 4/\kappa$ . Under the present multiple SLE/GFF coupling with  $\beta = 8/\kappa$ , we can prove that the multiple SLE driven by DYS<sub>β</sub> exhibits the phase transitions at  $\kappa = 4$  and  $\kappa = 8$  in the similar way to the original  $SLE_{\kappa}$  with a single SLE curve.

Keywords Itô's formula, martingales, conformal transformations, point processes, SLE curves, imaginary surfaces, multiple SLE/GFF coupling

<sup>∗</sup>This manuscript was prepared for the lectures given in the 4th ZiF Summer School 'Randomness in Physics and Mathematics: From Integrable Probability to Disordered Systems' held at ZiF–Center for Interdisciplinary Research, Bielefeld University, Germany, from 1st to 13th August 2022, which has been organized by Gernot Akemann and Friedrich Götze.

<sup>†</sup>Department of Physics, Faculty of Science and Engineering, Chuo University, Kasuga, Bunkyo-ku, Tokyo 112- 8551, Japan; e-mail: katori@phys.chuo-u.ac.jp

# Contents





# <span id="page-2-0"></span>1 Determinantal Martingales and Determinantal Stochastic Processes

# <span id="page-2-1"></span>1.1 One-dimensional standard Brownian motion (BM)

First of all, we consider the motion of a Brownian particle in one-dimensional space R starting from the origin 0 at time  $t = 0$ . At each time  $t > 0$ , the particle position is randomly distributed, and each realization of its path (trajectory) is denoted by  $\omega$  and called a **sample path** or simply a **path**. Let  $\Omega$  be the collection of all sample paths and we call it the **sample path space**. The position of the Brownian particle at time  $t \geq 0$  in a path  $\omega \in \Omega$  is written as  $B(t, \omega)$ . Usually we omit  $\omega$  and simply write it as  $B(t)$ ,  $t \geq 0$ .

We represent each event associated with the process by a subset of  $\Omega$ , and the collection of all events is denoted by F. The whole sample path space  $\Omega$  and the empty set  $\emptyset$  are in F. For any two sets  $A, B \in \mathcal{F}$ , we assume that  $A \cup B \in \mathcal{F}$  and  $A \cap B \in \mathcal{F}$ . If  $A \in \mathcal{F}$ , then its complement  $A^c$  is also in F. It is closed for any infinite sum of events in the sense that, if  $A_n \in \mathcal{F}, n = 1, 2, \ldots$ , then  $\cup_{n>1}A_n \in \mathcal{F}$ . Such a collection of events F is said to be a  $\sigma$ -field (sigma-field). The symbol  $\sigma$  is for 'sum'.

A **probability measure** P is a nonnegative function defined on the  $\sigma$ -field F. Since any element of  $\mathcal F$  is given by a set as above, any input of P is a set; P is a set function. It satisfies the following properties:  $P[A] \geq 0$  for all  $A \in \mathcal{F}$ ,  $P[\Omega] = 1$ ,  $P[\emptyset] = 0$ , and if  $A, B \in \mathcal{F}$  are disjoint,  $A \cap B = \emptyset$ , then  $P[A \cup B] = P[A] + P[B]$ . In particular,  $P[A^c] = 1 - P[A]$  for all  $A \in \mathcal{F}$ . The triplet  $(\Omega, \mathcal{F}, P)$  is called the **probability space**.

The smallest  $\sigma$ -field containing all intervals on R is called the **Borel**  $\sigma$ **-field** and denoted by  $\mathcal{B}(\mathbb{R})$ . A random variable or measurable function is a real-valued function  $f(\omega)$  on  $\Omega$ such that, for every Borel set  $A \in \mathcal{B}(\mathbb{R})$ ,  $f^{-1}(A) \in \mathcal{F}$ . Two events A and B are said to be independent if  $P[A \cap B] = P[A]P[B]$ . Two random variables X and Y are independent if the events  $A = \{X : X \in A\}$  and  $B = \{Y : Y \in B\}$  are independent for any  $A, B \in \mathcal{B}(\mathbb{R})$ .

The one-dimensional standard Brownian motion,  $\{B(t, \omega) : t \geq 0\}$ , has the following three properties:

- (BM1)  $B(0, \omega) = 0$  with probability one.
- (BM2) There is a subset of the sample path space  $\widetilde{\Omega} \subset \Omega$ , such that  $P[\widetilde{\Omega}] = 1$  and for any  $\omega \in \Omega$ ,  $B(t, \omega)$  is a real continuous function of t. We say that  $B(t)$  has a **continuous path** almost surely (a.s., for short).

(BM3) For an arbitrary  $M \in \mathbb{N} := \{1, 2, 3, \dots\}$ , and for any sequence of times,  $t_0 := 0$  $t_1 < \cdots < t_M$ , the increments  $B(t_m) - B(t_{m-1}), m = 1, 2, \ldots, M$ , are independent, and each increment is in the normal distribution (the Gaussian distribution) with mean 0 and variance  $\sigma^2 = t_m - t_{m-1}$ . It means that for any  $1 \leq m \leq M$  and  $a < b$ ,

$$
P[B(t_m) - B(t_{m-1}) \in [a, b]] = \int_a^b p(t_m - t_{m-1}, z|0) dz,
$$

where we define for  $x, y \in \mathbb{R}$ 

<span id="page-3-0"></span>
$$
p(t,y|x) = \begin{cases} \frac{1}{\sqrt{2\pi t}} e^{-(x-y)^2/2t}, & \text{for } t > 0, \\ \delta(x-y), & \text{for } t = 0. \end{cases}
$$
 (1.1)

Unless otherwise noted, the one-dimensional standard Brownian motion is simply abbreviated to BM in this lecture note. The probability measure P for the BM in particular is called the Wiener measure. The expectation with respect to the probability measure P is denoted by E. We write the conditional probability as  $P[\cdot|C]$ , where C denotes the condition. The conditional expectation is similarly written as  $E[\cdot|C]$ .

The third property (**BM3**) implies that for any  $0 \le s \le t < \infty$ 

$$
P[B(t) \in A|B(s) = x] = \int_{A} p(t - s, y|x)dy
$$
\n(1.2)

holds,  $\forall A \in \mathcal{B}(\mathbb{R})$ ,  $\forall x \in \mathbb{R}$ . Therefore the integral kernel  $p(t, y|x)$  given by  $(1.1)$  is called the transition probability density function of Brownian motion starting from  $x$ . The probability that the BM is observed in a region  $A_m \in \mathcal{B}(\mathbb{R})$  at time  $t_m$  for each  $m = 1, 2, ..., M$  is then given by

<span id="page-3-1"></span>
$$
P[B(t_m) \in A_m, m = 1, 2, ..., M] = \int_{A_1} dx_1 \cdots \int_{A_M} dx_M \prod_{m=1}^M p(t_m - t_{m-1}, x_m | x_{m-1}), \qquad (1.3)
$$

where  $x_0 := 0$ .

By (BM3), we can see that, for any  $c > 0$ , the probability distribution of  $B(c^2t)/c$  is equivalent with that of  $B(t)$  at arbitrary time  $t \geq 0$ . It is written as

$$
\frac{1}{c}B(c^2t) \stackrel{\text{d}}{=} B(t), \quad \forall c > 0,
$$

where the symbol  $\stackrel{d}{=}$  is for **equivalence in distribution**. Moreover, [\(1.3\)](#page-3-1) implies that, for any  $c > 0$ ,  $B(t)$ ,  $t \ge 0$  and its time changed process with  $t \mapsto c^2t$  multiplied by a factor  $1/c$  (dilatation) follow the same probability law. This **equivalence in probability law** of stochastic processes is expressed as

<span id="page-3-2"></span>
$$
(B(t))_{t\geq 0} \stackrel{\text{(law)}}{=} \left(\frac{1}{c}B(c^2t)\right)_{t\geq 0}, \quad \forall c \geq 0,
$$
\n
$$
(1.4)
$$

and called the scaling property of Brownian motion.

For  $a > 0$ , let  $T_a = \inf\{t \geq 0 : B(t) = a\}$ . Then for any  $t \geq 0$ ,

<span id="page-3-3"></span>
$$
P[T_a < t, B(t) < a] = P[T_a < t, B(t) > a],
$$
\n(1.5)

since the transition probability density [\(1.1\)](#page-3-0) is a symmetric function of the increment  $y - x$ . This property is called the **reflection principle** of BM. For  $\{\omega : B(t) > a\} \subset \{\omega : T_a < t\}, a > 0$ , the above is equal to  $P[B(t) > a]$ .

The formula [\(1.3\)](#page-3-1) also means that for any fixed  $s \geq 0$ , under the condition that  $B(s)$  is given,  ${B(t): t < s}$  and  ${B(t): t > s}$  are independent. This independence of the events in the future and those in the past is called the **Markov property**. A positive random variable  $\tau$  is called stopping time (or Markov time), if the event  $\{\omega : \tau \leq t\}$  is determined by the behavior of the process until time t and independent of that after t. For any stopping time  $\tau$ ,  $\{B(t): t \leq \tau\}$  and  ${B(t): t > \tau}$  are independent. It is called the **strong Markov property**. A stochastic process which has the strong Markov property and has a continuous path almost surely is generally called a diffusion process.

For each time  $t \in [0,\infty)$ , we write the smallest  $\sigma$ -field generated by the BM up to time  $t \geq 0$ as  $\sigma(B(s): 0 \leq s \leq t)$  and define

$$
\mathcal{F}_t := \sigma(B(s) : 0 \le s \le t), \quad t \ge 0. \tag{1.6}
$$

By definition, with respect to any event in  $\mathcal{F}_t$ ,  $B(s)$  is measurable at every  $s \in [0, t]$ . Then we have a nondecreasing family  $\{\mathcal{F}_t : t \ge 0\}$  of sub- $\sigma$ -fields of the original  $\sigma$ -field  $\mathcal{F}$  in the probability space  $(\Omega, \mathcal{F}, P)$  such that  $\mathcal{F}_s \subset \mathcal{F}_t \subset \mathcal{F}$  for  $0 \leq s < t < \infty$ . We call this family of  $\sigma$ -fields a **filtration**.

Since the probability density of increment in any time-interval  $t - s > 0$ ,  $p(t - s, z|0)$ , has mean zero, BM satisfies the equality

$$
E[B(t)|\mathcal{F}_s] = B(s), \quad 0 \le s < t < \infty, \quad \text{a.s.} \tag{1.7}
$$

That is, the mean is constant in time, even though the variance increases in time as  $\sigma^2 = t$ . Processes with such a property are called **martingales**. We note that for  $0 \le s < t < \infty$ 

$$
E[B(t)^{2}|\mathcal{F}_{s}] = E[(B(t) - B(s))^{2} + 2(B(t) - B(s))B(s) + B(s)^{2}|\mathcal{F}_{s}]
$$
  
= 
$$
E[(B(t) - B(s))^{2}|\mathcal{F}_{s}] + 2E[(B(t) - B(s))B(s)|\mathcal{F}_{s}] + E[B(s)^{2}|\mathcal{F}_{s}].
$$

By the property (BM3) and the definition of  $\mathcal{F}_s$ ,

$$
E[(B(t) - B(s))^2 | \mathcal{F}_s] = t - s,
$$
  
\n
$$
E[(B(t) - B(s))B(s) | \mathcal{F}_s] = E[B(t) - B(s) | \mathcal{F}_s] B(s) = 0,
$$
  
\n
$$
E[B(s)^2 | \mathcal{F}_s] = B(s)^2.
$$

Then we have the equality

<span id="page-4-1"></span>
$$
E[B(t)^{2} - t | \mathcal{F}_{s}] = B(s)^{2} - s, \quad 0 \le s < t < \infty, \quad \text{a.s.}
$$
 (1.8)

It means that  $B(t)^2 - t$  is a martingale.

For the transition probability density of BM [\(1.1\)](#page-3-0), it should be noted that  $p(\cdot, y|x) = p(\cdot, x|y)$ for any  $x, y \in \mathbb{R}$ , and  $u_t(x) := p(t, y|x)$  is a unique solution of the **heat equation (diffusion**) equation)

<span id="page-4-0"></span>
$$
\frac{\partial}{\partial t}u_t(x) = \frac{1}{2}\frac{\partial^2}{\partial x^2}u_t(x), \quad x \in \mathbb{R}, \quad t \ge 0
$$
\n(1.9)

with the initial condition  $u_0(x) = \delta(x - y)$ . The solution of [\(1.9\)](#page-4-0) with the initial condition  $u_0^f$  $_{0}^{J}(x) =$  $f(x), x \in \mathbb{R}$  is then given by

<span id="page-5-1"></span>
$$
u_t^f(x) = \mathcal{E}^x[f(B(t))] = \int_{-\infty}^{\infty} f(y)p(t, y|x)dy,
$$
\n(1.10)

if f is a measurable function satisfying the condition  $\int_{-\infty}^{\infty} e^{-ax^2} |f(x)| dx < \infty$  for some  $a > 0$ . Since  $p(t, y|x)$  plays as an integral kernel in [\(1.10\)](#page-5-1), it is also called the **heat kernel**.

For  $0 \leq s < t < \infty$ ,  $\xi \in \mathbb{R}$ , consider  $\mathbb{E}[e^{\sqrt{-1}\xi(B(t)-B(s))}|\mathcal{F}_s]$ . It is calculated as the **Fourier** transform of the transition probability density  $p(t - s, \cdot | 0)$ ;

$$
\int_{-\infty}^{\infty} e^{\sqrt{-1}\xi z} p(t-s, z|0) dz = \int_{-\infty}^{\infty} e^{\sqrt{-1}\xi z} \frac{e^{-z^2/2(t-s)}}{\sqrt{2\pi(t-s)}} dz
$$
  
=  $e^{-\xi^2(t-s)/2}$ . (1.11)

The obtained function of  $\xi \in \mathbb{R}$ ,

$$
E[e^{\sqrt{-1}\xi(B(t)-B(s))}|\mathcal{F}_s] = e^{-\xi^2(t-s)/2}, \quad 0 \le s < t < \infty,
$$
\n(1.12)

<span id="page-5-0"></span>is called the characteristic function of BM.

### 1.2 Itô's formula and Kolmogorov equation

Let  $X(t), t \geq 0$  be a one-dimensional diffusion process on the probability space  $(\Omega_X, \mathcal{F}_X, \mathcal{P}_X)$ , where the expectation is written as  $E<sub>X</sub>$  and the filtration is given by the **natural filtration** of  $X; (\mathcal{F}_X)_t = \sigma(X(s): 0 \le s \le t), t \ge 0$ . For each time interval  $[0, t], t > 0$ , put  $n \in \mathbb{N}$  and let  $\Delta_n = \Delta_n([0, t])$  be a subdivision of  $[0, t]$  with  $t_0 := 0 < t_1 < \cdots < t_{n-1} < t_n := t$ . Then we define  $Q^{\Delta_n}(t) = \sum_{m=1}^n (X(t_m) - X(t_{m-1}))^2$ . If there is a process  $Q(t), t \ge 0$  such that

$$
\lim_{n \to \infty} \mathcal{P}_X[|Q^{\Delta_n}(t) - Q(t)| > \varepsilon] = 0, \quad \forall \varepsilon > 0 \tag{1.13}
$$

holds provided max<sub>1≤m≤n</sub>  $|t_m - t_{m-1}| \to 0$  as  $n \to \infty$ , then we call  $Q(t)$ ,  $t \geq 0$ , the **quadratic variation** of  $X(t)$ ,  $t \ge 0$  and express it by  $\langle X, X \rangle_t$ ,  $t \ge 0$ .

For BM,  $B(t)$ ,  $t \geq 0$ , set  $n \in \mathbb{N}$ ,  $t_0 := 0 < t_1 < \cdots < t_{n-1} < t_n := t$  and put  $Q_{BM}^{\Delta_n}(t) =$  $\sum_{m=1}^{n} (B(t_m) - B(t_{m-1}))^2$ . By the property (BM3), the mean is given by  $E[Q_{BM}^{\Delta_n}(t)] = \sum_{m=1}^{n} (t_m (t_{m-1}) = t$ . The variance of  $Q_{BM}^{\Delta_n}(t)$ 

$$
\sigma_{BM}^{\Delta_n}(t)^2 := \mathbb{E}[(Q_{BM}^{\Delta_n}(t) - t)^2]
$$
  
=  $\mathbb{E}\left[\left\{\sum_{m=1}^n \left\{(B(t_m) - B(t_{m-1}))^2 - (t_m - t_{m-1})\right\}\right\}^2\right]$ 

is calculated as

<span id="page-6-0"></span>
$$
\sum_{m=1}^{n} \left\{ \mathbb{E}[(B(t_m) - B(t_{m-1}))^4] - 2(t_m - t_{m-1})\mathbb{E}[(B(t_m) - B(t_{m-1}))^2] + (t_m - t_{m-1})^2 \right\}
$$
  
= 
$$
\sum_{m=1}^{n} \left\{ 3(t_m - t_{m-1})^2 - 2(t_m - t_{m-1})^2 + (t_m - t_{m-1})^2 \right\}
$$
  
= 
$$
2\sum_{m=1}^{n} (t_m - t_{m-1})^2 \le 2t \max_{1 \le m \le n} |t_m - t_{m-1}|,
$$
 (1.14)

where independence of increments of BM mentioned in **(BM3)** and

$$
E[(B(t) - B(s))^4] = 3(t - s)^2, \quad 0 \le s \le t
$$
\n(1.15)

were used. By Chebyshev's inequality, we have

$$
\mathbf{P}[|Q_{\text{BM}}^{\Delta_n}(t) - t| > \varepsilon] \le \frac{\sigma_{\text{BM}}^{\Delta_n}(t)^2}{\varepsilon^2}, \quad \forall \varepsilon > 0.
$$

Provided that  $\max_{1 \le m \le n} |t_m - t_{m-1}| \to 0$  as  $n \to \infty$ , [\(1.14\)](#page-6-0) gives  $\lim_{n \to \infty} \sigma_{BM}^{\Delta_n}(t) = 0$  and it proves

<span id="page-6-1"></span>
$$
\langle B, B \rangle_t = t, \qquad t \ge 0. \tag{1.16}
$$

For a stopping time  $\tau$ , we put  $X^{\tau}(t) := X(t \wedge \tau), t \geq 0$ , where  $t \wedge \tau := \min\{t, \tau\}$ . We define a diffusion process  $X(t), t \geq 0$  as a **local martingale**, if there exists stopping times  $\tau_n, n \in \mathbb{N}$ such that (i) the sequence  ${\tau_n}_{n\in\mathbb{N}}$  is nondecreasing and  $\lim_{n\to\infty}\tau_n = \infty$  a.s., and (ii) for every n, the process  $X^{\tau_n}(t)$ ,  $t \geq 0$  is a martingale. When  $X(t)$  is a local martingale, we can prove that a unique increasing continuous process is given by  $\langle X, X \rangle_t, t \geq 0$  such that  $\langle X, X \rangle_0 = 0$  and  $X(t)^2 - \langle X, X \rangle_t, t \geq 0$  provides a local martingale.

Assume that  $X(t)$  and  $Y(t)$ ,  $t \geq 0$  are both local martingales. Then  $(X(t) + Y(t))^2 - \langle X + \rangle$  $Y, X + Y$ <sub>t</sub> and  $(X(t) - Y(t))^2 - \langle X - Y, X - Y \rangle_t$ ,  $t \ge 0$  are local martingales. Therefore, their difference  $4X(t)Y(t) - \{\langle X + Y, X + Y \rangle_t - \langle X - Y, X - Y \rangle_t\}, t \ge 0$  is also a local martingale. For any pair of local martingales  $X(t)$  and  $Y(t)$ ,  $t \geq 0$ , we define the **mutual quadratic variation** (cross variation) process as

$$
\langle X, Y \rangle_t := \frac{1}{4} \{ \langle X + Y, X + Y \rangle_t - \langle X - Y, X - Y \rangle_t \}, \quad t \ge 0. \tag{1.17}
$$

We can prove that, if  $B_i(t)$ ,  $t \geq 0$ ,  $i = 1, 2, \ldots, D$  are independent BMs, then

<span id="page-6-2"></span>
$$
\langle B_i, B_j \rangle_t = \delta_{ij} t, \quad i, j = 1, \dots, D, \quad t \ge 0. \tag{1.18}
$$

For a continuous process  $A(t)$ ,  $t \geq 0$ , here we consider the quantity

$$
S^{\Delta_n}(t) = \sum_{m=1}^{n} |A(t_m) - A(t_{m-1})|
$$

instead of  $Q^{\Delta_n}(t)$ , where  $\Delta_n, n \in \mathbb{N}$  is a subdivision of the time interval [0, t]. We can see that if  $\Delta_n \subset \Delta_{n+1}$ , then  $S^{\Delta_n}(t) \leq S^{\Delta_{n+1}}(t)$ ,  $\forall t \geq 0$ . Assume that  $\max_{1 \leq m \leq n} |t_m - t_{m-1}| \to 0$  as  $n \to \infty$ . Let  $\lim_{n\to\infty} \sup_{\Delta_n} S^{\Delta_n}(t) = S(t) \leq \infty$  and call it the **variation** of A on [0, t]. If  $S(t) < \infty$  for every t, then the process  $A(t)$ ,  $t \geq 0$  is of **finite variation**. Let  $S_{BM}(t)$ ,  $t \geq 0$  be the variation of BM. We have

$$
Q_{\text{BM}}^{\Delta_n}(t) \le \sup_{1 \le m \le n} |B(t_m) - B(t_{m-1})| S_{\text{BM}}(t), \quad t \ge 0, \quad n \in \mathbb{N}.
$$

By the property (BM2), RHS becomes 0 as  $n \to \infty$  if  $S_{BM}(t)$  is finite. On the other hand, we have proved  $Q_{BM}^{\Delta_n}(t) \to \langle B, B \rangle_t$  as  $n \to \infty$  in probability and the fact [\(1.16\)](#page-6-1). Hence  $S_{BM}(t) = \infty$ a.s.  $\forall t \geq 0$ .

Let  $N \in \mathbb{N}$  and  $\{X_1(t),...,X_N(t)\}\,$ ,  $t \geq 0$  be a set of diffusion processes. Put  $\mathbf{X}(t) =$  $(X_1(t),...,X_N(t))$ . Let F be a real function of  $(t, x) \in [0, \infty) \times \mathbb{R}^N$ , which is bounded and has a bounded first-order derivative with respect to  $t$  and bounded first- and second-order derivatives with respect to  $x_j, 1 \leq j \leq N$ , and we denote this by  $F \in C^{1,2}_b$  $b^{\text{1,2}}$ . We know that, in order to describe the statistics of a function of several random variables, we have to take into account the 'propagation of error'. For the process  $F(t, X(t))$  that is defined as a function of t as well as a functional of processes  $X_1(t), \ldots, X_N(t), t \geq 0$ , Itô's formula gives an equation which governs the difference of  $F(t, \mathbf{X}(\cdot))$  as

<span id="page-7-0"></span>
$$
dF(t, \mathbf{X}(t)) = \sum_{i=1}^{N} \frac{\partial F}{\partial x_i}(t, \mathbf{X}(t)) dX_i(t) + \frac{\partial F}{\partial t}(t, \mathbf{X}(t)) dt + \frac{1}{2} \sum_{1 \le i, j \le N} \frac{\partial^2 F}{\partial x_i \partial x_j}(t, \mathbf{X}(t)) d\langle X_i, X_j \rangle_t, \quad t \ge 0.
$$
 (1.19)

The first term gives a local martingale and the second and third terms give processes of finite variations. A continuous process  $X$  given by the sum of a local martingale  $M$  and a finite-variation process A,

$$
X(t) = M(t) + A(t), \quad t \ge 0,
$$

is called a **semi-martingale**. We will use the fact that if the continuous process  $F(t, \mathbf{X}(t))$ ,  $t > 0$ is a local martingale, then its finite-variation part should vanish, and vice versa.

Let  $X(t)$ ,  $t \geq 0$  be a one-dimensional diffusion in  $(\Omega_X, \mathcal{F}_X, P_X)$ , which satisfies the following SDE

<span id="page-7-2"></span>
$$
dX(t) = \sigma(X(t))dB(t) + b(X(t))dt, \quad X(0) = x.
$$
\n(1.20)

Here  $B(t)$ ,  $t \geq 0$  is BM and the functions  $\sigma, b : \mathbb{R} \to \mathbb{R}$  satisfy the condition that  $\exists K \geq 0$ , s.t.  $|\sigma(x) - \sigma(y)| \leq K|x - y|, |b(x) - b(y)| \leq K|x - y|, x, y \in \mathbb{R}$ . (This is called the **Lipschitz** continuity.) Put

$$
u(s,x) = \mathcal{E}^x[f(X(T-s))], \quad 0 \le s < T < \infty,\tag{1.21}
$$

with an  $(\mathcal{F}_X)_T$ -measurable bounded function f. By the Markov property, for  $0 \leq s < t < T < \infty$ ,

<span id="page-7-1"></span>
$$
u(s,x) = \mathbf{E}^x \Big[ \mathbf{E}^{X(t-s)} [f(X(T-t))]\Big]
$$
  
= 
$$
\mathbf{E}^x [u(t, X(t-s))].
$$
 (1.22)

Assume that  $u(s, x) \in C_6^{1,2}$  $b^{\text{1,2}}$ . Then by Itô's formula  $(1.19)$ ,

<span id="page-8-1"></span>
$$
u(t, X(t-s)) - u(s, x) = \int_0^{t-s} \left(\frac{\partial u}{\partial s} + \mathcal{L}u\right) (s+r, X(r)) dr
$$

$$
+ \int_0^{t-s} \sigma(X(r)) \frac{\partial u}{\partial x} (s+r, X(r)) dB(r), \tag{1.23}
$$

where

<span id="page-8-2"></span>
$$
(Lf)(x) = \frac{1}{2}a(x)\frac{d^2}{dx^2}f(x) + b(x)\frac{d}{dx}f(x)
$$
\n(1.24)

with

$$
a(x) = \sigma(x)^2. \tag{1.25}
$$

Since [\(1.22\)](#page-7-1) holds, taking expectation of [\(1.23\)](#page-8-1) gives

$$
\mathcal{E}^{x}\left[\int_{0}^{t-s}\left(\frac{\partial u}{\partial s} + \mathsf{L}u\right)(s+r, X(r))dr\right] = 0,
$$

where the expectation of the second term in RHS of [\(1.23\)](#page-8-1) vanished for it is a martingale. Set  $h = t - s$ , divide the both side of the above equation by h and take the limit  $h \to 0$ . Then we have

<span id="page-8-4"></span>
$$
\frac{\partial u(s,x)}{\partial s} + \mathsf{L}u(s,x) = 0. \tag{1.26}
$$

Since L acts as a differential operator with respect to the initial value  $x$ , it is called the **backward** Kolmogorov equation. The differential operator [\(1.24\)](#page-8-2) is called the generator of the diffusion process.

# <span id="page-8-0"></span>1.3 Martingale polynomials of BM

Let

<span id="page-8-3"></span>
$$
G_{\alpha}(t,x) = e^{\alpha x - \alpha^2 t/2}.
$$
\n(1.27)

For BM, we perform the transformation with the parameter  $\alpha \in \mathbb{C} := \{z = x + \sqrt{-1}y : x, y \in \mathbb{R}\},\$  $B \mapsto \check{B}_{\alpha}$  defined by

$$
\check{B}_{\alpha}(t) = G_{\alpha}(t, B(t)), \quad t \ge 0,
$$
\n(1.28)

which is called the Esscher transformation. Since

$$
E[e^{\alpha B(t)}] = \int_{-\infty}^{\infty} e^{\alpha x} p(t, x | 0) dx = e^{\alpha^2 t/2}, \quad t \ge 0.
$$

we can see that

$$
G_{\alpha}(t, B(t)) = \frac{e^{\alpha B(t)}}{\mathcal{E}[e^{\alpha B(t)}]}, \quad t \ge 0,
$$
\n(1.29)

and hence

$$
E[G_{\alpha}(t, B(t)) | \mathcal{F}_s] = \frac{E[e^{\alpha B(t)} | \mathcal{F}_s]}{E[e^{\alpha B(t)}]}
$$
  
= 
$$
\frac{E[e^{\alpha B(s)} e^{\alpha(B(t) - B(s))} | \mathcal{F}_s]}{E[e^{\alpha B(s)} e^{\alpha(B(t) - B(s))}]}, \quad 0 < s < t.
$$

By the definition of  $\mathcal{F}_s$  and independence of increment of BM (the property (BM3)), the numerator is equal to  $e^{\alpha B(s)} \mathbb{E}[e^{\alpha(B(t)-B(s))}]$ , and the denominator is equal to  $\mathbb{E}[e^{\alpha B(s)}]\mathbb{E}[e^{\alpha(B(t)-B(s))}]$ . Hence the above equals  $e^{\alpha B(s)} / E[e^{\alpha B(s)}] = G_{\alpha}(s, B(s))$ . Therefore,  $G_{\alpha}(t, B(t))$  is a martingale:

$$
E[G_{\alpha}(t, B(t)) | \mathcal{F}_s] = G_{\alpha}(s, B(s)), \quad 0 \le s \le t.
$$
\n(1.30)

The function [\(1.27\)](#page-8-3) is expanded as

$$
G_{\alpha}(t,x) = \sum_{n=0}^{\infty} m_n(t,x) \frac{\alpha^n}{n!}
$$
\n(1.31)

with

<span id="page-9-1"></span>
$$
m_n(t, x) = \left(\frac{t}{2}\right)^{n/2} H_n\left(\frac{x}{\sqrt{2t}}\right), \quad n \in \mathbb{N}_0 := \{0, 1, 2, \dots\}.
$$
 (1.32)

Here  ${H_n(x)}_{n \in \mathbb{N}_0}$  are the **Hermite polynomials** of degrees  $n \in \mathbb{N}_0$ ,

<span id="page-9-0"></span>
$$
H_n(x) := (-1)^n e^{x^2} \frac{d^n e^{-x^2}}{dx^n} = \sum_{k=0}^{[n/2]} (-1)^k \frac{n!}{k!(n-2k)!} (2x)^{n-2k},\tag{1.33}
$$

where for  $a \geq 0$ , [a] denotes the largest integer which is not larger than a. We can show the following contour integral representations of the Hermite polynomials,

<span id="page-9-3"></span>
$$
H_n(z) = \frac{n!}{2\pi\sqrt{-1}} \oint_{C(\delta_0)} d\eta \, \frac{e^{2\eta z - \eta^2}}{\eta^{n+1}}, \quad n \in \mathbb{N}_0,
$$
\n(1.34)

<span id="page-9-4"></span>where  $C(\delta_0)$  is a closed contour on the complex plane C encircling the origin 0 once in the positive direction.

**Lemma 1.1** The functions  $\{m_n(t,x)\}_{n\in\mathbb{N}_0}$  satisfy the following.

(i) They are monic polynomials of degrees  $n \in \mathbb{N}_0$  with time-dependent coefficients:

$$
m_n(t,x) = x^n + \sum_{k=0}^{n-1} c_n^{(k)}(t) x^k, \quad t \ge 0.
$$

(ii) For  $0 \le k \le n-1$ ,  $c_n^{(k)}(0) = 0$ . That is,

$$
m_n(0, x) = x^n, \quad n \in \mathbb{N}_0.
$$

(iii) If we set  $x = B(t)$ , they provide martingales:

$$
\mathcal{E}[m_n(t, B(t)) | \mathcal{F}_s] = m_n(s, B(s)), \quad 0 \le s \le t, \quad n \in \mathbb{N}_0.
$$
\n
$$
(1.35)
$$

We call  ${m_n(t, x)}_{n \in \mathbb{N}_0}$  the fundamental martingale polynomials associated with BM [\[30\]](#page-88-0). For  $n = 2$ , [\(1.33\)](#page-9-0) gives  $H_2(x) = 4x^2 - 2$ , and then  $m_2(t, x) = x^2 - t$  by [\(1.32\)](#page-9-1). We already proved in [\(1.8\)](#page-4-1) that  $m_2(t, B(t)) = B(t)^2 - t$  is a martingale.

The Fourier transformation of  $G_{\alpha}(t, x)$  with respect to the parameter  $\alpha \in \mathbb{R}$  is calculated as

<span id="page-9-2"></span>
$$
\widehat{G}_w(t,x) := \int_{-\infty}^{\infty} \frac{e^{-\sqrt{-1}\alpha w}}{2\pi} G_\alpha(t,x) d\alpha = \frac{e^{-(\sqrt{-1}x+w)^2/2t}}{\sqrt{2\pi t}}.
$$
\n(1.36)

Owing to the factor  $e^{-w^2/2t}$  in  $\widehat{G}_w(t, x)$ , the following calculations are justified,

$$
G_{\alpha}(t,x) = \int_{-\infty}^{\infty} e^{\sqrt{-1}\alpha w} \widehat{G}_{w}(t,x) dw = \sum_{n=0}^{\infty} \frac{\alpha^{n}}{n!} \int_{-\infty}^{\infty} (\sqrt{-1}w)^{n} \widehat{G}_{w}(t,x) dw, \qquad (1.37)
$$

which proves the integral representation of  $m_n(t, x)$ ,

$$
m_n(t,x) = \int_{-\infty}^{\infty} (\sqrt{-1}w)^n \hat{G}_w(t,x) dw
$$
  
= 
$$
\int_{-\infty}^{\infty} (\sqrt{-1}w)^n \frac{e^{-(\sqrt{-1}x+w)^2/2t}}{\sqrt{2\pi t}} dw, \quad t \ge 0, \quad n \in \mathbb{N}_0.
$$
 (1.38)

We define this type of integral transformation of a function  $f$  as

<span id="page-10-4"></span>
$$
\mathcal{I}[f(W)|(t,x)] = \int_{-\infty}^{\infty} f(\sqrt{-1}w)\widehat{G}_w(t,x)dw.
$$
 (1.39)

Then the above results are written as

<span id="page-10-5"></span>
$$
m_n(t,x) = \mathcal{I}[W^n|(t,x)], \quad t \ge 0, \quad n \in \mathbb{N}_0.
$$
\n
$$
(1.40)
$$

# <span id="page-10-1"></span><span id="page-10-0"></span>1.4 Bessel processes  $(BES<sub>D</sub>)$

#### 1.4.1 Radial part of D-dimensional BM

Let  $D \in \mathbb{N}$  denote the spatial dimension. For  $D \geq 2$ , the D-dimensional BM in  $\mathbb{R}^D$  starting from the position  $\mathbf{x} = (x_1, \dots, x_D) \in \mathbb{R}^D$  is defined by the D-dimensional vector-valued diffusion process,

<span id="page-10-3"></span>
$$
\boldsymbol{B}^{\boldsymbol{x}}(t) = (B_1^{x_1}(t), B_2^{x_2}(t), \dots, B_D^{x_D}(t)), \quad t \ge 0,
$$
\n(1.41)

where

$$
B_i^{x_i}(t) := x_i + B_i(t), \quad t \ge 0, \quad i = 1, \dots, D
$$

are independent one-dimensional BMs starting from  $x_i \in \mathbb{R}$ .

The D-dimensional **Bessel process** is defined as the absolute value (*i.e.*, the radial coordinate) of the D-dimensional Brownian motion,

<span id="page-10-2"></span>
$$
R^{x}(t) := |\mathbf{B}^{x}(t)|
$$
  
=  $\sqrt{B_1^{x_1}(t)^2 + \dots + B_D^{x_D}(t)^2}, \quad t \ge 0,$  (1.42)

where the initial value is given by  $R^x(0) = x = |\mathbf{x}| = \sqrt{x_1^2 + \cdots + x_D^2} \ge 0$ . By definition  $R^x(t)$  is nonnegative,  $R^x(t) \in \mathbb{R}_+ \cup \{0\}$ , where  $\mathbb{R}_+ := \{x \in \mathbb{R} : x > 0\}$ . We will abbreviate the D-dimensional Bessel process to  $BES<sub>D</sub>$ .

By this definition,  $R^x(t)$  is a functional of D-tuples of diffusion processes  $\{B_i^{x_i}(t)\}_{i=1}^D$ ,  $t \geq 0$ . Now we apply Itô's formula [\(1.19\)](#page-7-0) to  $BES_D$  [\(1.42\)](#page-10-2). Assume  $x = |\mathbf{x}| > 0$ . In this case

$$
R^{x}(t) = F(\mathbf{B}^{x}(t)), \quad t \ge 0 \quad \text{with} \quad F(y) = \sqrt{\sum_{i=1}^{D} y_{i}^{2}}.
$$
 (1.43)

We see that

$$
\frac{\partial F}{\partial t} = 0, \quad \frac{\partial F}{\partial y_i} = \frac{y_i}{F}, \quad \frac{\partial^2 F}{\partial y_i \partial y_j} = \frac{\delta_{ij}}{F} - \frac{y_i y_j}{F^3}, \quad 1 \le i, j \le D.
$$

From [\(1.18\)](#page-6-2), we have  $d\langle B_i^{x_i}, B_j^{x_j}\rangle_t = \langle dB_i^{x_i}, dB_j^{x_j}\rangle_t = \delta_{ij}dt$ ,  $1 \le i, j \le D, t \ge 0$ . Then, the second term of  $(1.19)$  for BES<sub>D</sub> becomes

$$
\frac{1}{2} \sum_{1 \le i,j \le D} \left\{ \frac{\delta_{ij}}{F(B^x(t))} - \frac{B_i^{x_i}(t)B_j^{x_j}(t)}{F(B^x(t))^3} \right\} \delta_{ij} dt = \frac{D-1}{2} \frac{1}{F(B^x(t))} dt = \frac{D-1}{2} \frac{dt}{R^x(t)}.
$$

On the other hand, the first term of  $(1.19)$  for  $BES<sub>D</sub>$  is

<span id="page-11-0"></span>
$$
\frac{1}{R^x(t)}\sum_{i=1}^D B_i^{x_i}(t)dB_i^{x_i}(t).
$$
\n(1.44)

It seems to be complicated, but the quadratic variation is calculated as

$$
\left\langle \frac{1}{R^x} \sum_{i=1}^D B_i^{x_i} dB_i^{x_i}, \frac{1}{R^x} \sum_{j=1}^D B_j^{x_j} dB_j^{x_j} \right\rangle_t = \frac{1}{R^x(t)^2} \sum_{i=1}^D \sum_{j=1}^D B_i^{x_i}(t) B_j^{x_j}(t) \langle dB_i^{x_i}, dB_j^{x_j} \rangle_t
$$

$$
= \frac{1}{R^x(t)^2} \sum_{i=1}^D \sum_{j=1}^D B_i^{x_i}(t) B_j^{x_j}(t) \delta_{ij} dt = dt,
$$

where the independence of BMs [\(1.18\)](#page-6-2) and the definition,  $R^x(t)^2 = \sum_{i=1}^D B_i^{x_i}(t)^2$ , have been used. That is, [\(1.44\)](#page-11-0) is equivalent in probability law with an infinitesimal increment of a diffusion process with quadratic variation dt. Then, by introducing a BM,  $B^x(t)$ ,  $t \geq 0$ , which is different from  $B_i^{x_i}(t), t \geq 0, x_i \in \mathbb{R}, i = 1, 2, \ldots, D$  and is started at  $x = |\mathbf{x}| > 0$ , [\(1.44\)](#page-11-0) is identified with  $dB^{x}(t), t \geq 0$ . Hence we have obtained the following equation for  $BES_D$ ,

<span id="page-11-1"></span>
$$
dR^{x}(t) = dB^{x}(t) + \frac{D-1}{2} \frac{dt}{R^{x}(t)}, \quad x > 0, \quad 0 \le t < T^{x}, \tag{1.45}
$$

where  $T^x = \inf\{t > 0 : R^x(t) = 0\}.$ 

The first term of RHS,  $dB^x(t)$ , denotes the infinitesimal increment of BM starting from  $x > 0$ at time  $t = 0$ . This martingale term gives randomness to the motion. On the other hand, if  $D > 1$ , for  $dt > 0$ , the second term in RHS of [\(1.45\)](#page-11-1) is positive definite. It means that there is a drift to increase the value of  $R^x(t)$ . This drift term is increasing in D and decreasing in  $R^x(t)$ . Since as  $R^x(t) \downarrow 0$ , the drift term  $\uparrow \infty$ , it seems that a 'repulsive force' is acting to the D-dimensional BM,  $B^{x}(t), |x| > 0$  to keep the distance from the origin be positive,  $R^{x}(t) = |B^{x}(t)| > 0$  and avoid a collision of the Brownian particle with the origin. A differential equation such as [\(1.45\)](#page-11-1), which involves a random fluctuation term and a drift term is called a stochastic differential equation (SDE).

The following equivalence in probability law is established for arbitrary  $x > 0$ ,

$$
\left(\frac{1}{x}R^x(x^2t)\right)_{t\geq 0} \stackrel{\text{(law)}}{=} (R^1(t))_{t\geq 0}.\tag{1.46}
$$

It inherits [\(1.4\)](#page-3-2) and is called the scaling property of the Bessel process.

The SDE for BES<sub>D</sub> is given by the equation [\(1.20\)](#page-7-2) with  $\sigma(x) := 1$  and  $b(x) = (D-1)/(2x)$ . Then the generator of  $BES<sub>D</sub>$  is obtained as

<span id="page-12-1"></span>
$$
\mathsf{L}^D = \frac{1}{2} \frac{\partial^2}{\partial x^2} + \frac{D - 1}{2x} \frac{\partial}{\partial x}.\tag{1.47}
$$

Let  $p_D(t - s, y|x)$  be the transition probability density of BES<sub>D</sub> from x at times  $s \geq 0$  to y at times  $t \geq s$ . For any  $t \geq s, y \in \mathbb{R}_+$ ,  $p_D(t-s, y|x), x > 0$  solves [\(1.26\)](#page-8-4) with [\(1.47\)](#page-12-1) under the condition  $\lim_{s \uparrow t} p_D(t - s, y|x) = \delta(x - y)$ . In other words,  $p_D(t, y|x)$  solves

$$
\frac{\partial}{\partial t}p_D(t, y|x) = \mathsf{L}_D p_D(t, y|x) \tag{1.48}
$$

under the initial condition  $p_D(0, y|x) = \delta(x-y)$ , which is called the backward Kolmogorov equation for  $BES_D$ .

Let  $I_{\nu}(z)$  be the **modified Bessel function** of the first kind defined by

<span id="page-12-5"></span>
$$
I_{\nu}(z) = \sum_{n=0}^{\infty} \frac{1}{\Gamma(n+1)\Gamma(n+1+\nu)} \left(\frac{z}{2}\right)^{2n+\nu}
$$
\n(1.49)

with the gamma function

$$
\Gamma(z) = \int_0^\infty e^{-u} u^{z-1} du, \quad \text{Re } z > 0.
$$
 (1.50)

The function  $I_{\nu}(z)$  solves the **Bessel differential equation** 

$$
\frac{d^2w}{dz^2} + \frac{1}{z}\frac{dw}{dz} - \left(1 + \frac{\nu^2}{z^2}\right)w = 0.
$$
\n(1.51)

Then we can show that

<span id="page-12-3"></span>
$$
p_D(t, y|x) = \begin{cases} \frac{1}{t} \frac{y^{\nu+1}}{x^{\nu}} e^{-(x^2+y^2)/2t} I_{\nu} \left(\frac{xy}{t}\right), & t > 0, x > 0, y \ge 0, \\ \frac{y^{2\nu+1}}{2^{\nu} t^{\nu+1} \Gamma(\nu+1)} e^{-y^2/2t}, & t > 0, x = 0, y \ge 0 \\ \delta(y-x), & t = 0, x, y \ge 0, \end{cases}
$$
(1.52)

where the index  $\nu$  is specified by the dimension D as

<span id="page-12-2"></span>
$$
\nu = \frac{D-2}{2} \quad \Longleftrightarrow \quad D = 2(\nu + 1). \tag{1.53}
$$

This fact that  $p_D(t, y|x)$  is expressed by using  $I_\nu(z)$  gives the reason why the process  $R^x(t)$  is called the Bessel process.

## <span id="page-12-0"></span>1.4.2 BES<sub>3</sub> and absorbing BM

When  $D = 3$ ,  $\nu = 1/2$  by [\(1.53\)](#page-12-2), and we can use the equality  $I_{1/2}(z) = \sqrt{2/\pi z} \sinh z = (e^z (e^{-z})/\sqrt{2\pi z}$ . Then  $(1.52)$  gives

<span id="page-12-4"></span>
$$
p_3(t, y|x) = \frac{y}{x} \Big\{ p(t, y|x) - p(t, y| - x) \Big\}
$$
 (1.54)



<span id="page-13-1"></span>Figure 1: One realization of a Brownian path from  $x > 0$  to  $y > 0$  is represented by a red curve and denoted by path A, which visits the nonpositive region  $\mathbb{R}_-\cup\{0\}$ , where  $\mathbb{R}_-:=\{x\in\mathbb{R}:x<0\}$ . The first time the path A hits the origin is denoted by  $\tau$ . Path B, which is represented by a black curve, is a mirror image of path A with respect to the origin  $x = 0$ , which is running from  $-x < 0$ to  $-y < 0$ . The path C (blue curve) is then defined as the concatenation of the part of path B from time zero up to time  $\tau$  and the part of path A after  $\tau$  such that it runs from  $-x < 0$  to  $y > 0$ . By the reflection principle of BM [\(1.5\)](#page-3-3) applied at time  $\tau$ , we see that there exists a bijection between path A and path C, which have the same probability weight as Brownian paths. Since the Brownian path A contributes to  $p(t, y|x)$  and the Brownian path C contributes to  $p(t, y|-x)$ , such a path from  $x > 0$  to  $y > 0$  that visits  $\mathbb{R}_-$  is cancelled in  $q_{\text{abs}}(t, y|x)$  given by [\(1.55\)](#page-13-0). In Fig[.1.4.2](#page-13-1) (b), a path from  $x > 0$  to  $y > 0$  which stays in the positive region  $\mathbb{R}_+$  is considered (path A). In this case there is no such a path that perfectly cancels the contribution of path A to [\(1.55\)](#page-13-0) as path B in the case (a). In summary,  $q_{\text{abs}}(t, y|x) = p(t, y|x) - p(t, y|x)$  gives the total weight of Brownian paths which do not hit the origin.

for  $t > 0, x > 0, y \ge 0$ , where  $p(t, y|x)$  is the transition probability density of BM started at x given by  $(1.1)$ . If we put

<span id="page-13-0"></span>
$$
q_{\rm abs}(t, y|x) = p(t, y|x) - p(t, y| - x),
$$
\n(1.55)

we see that  $q_{\text{abs}}(t, 0|x) = 0$  for any  $x > 0$ , since the transition probability density of BM,  $p(t, y|x)$ , is an even function of  $y - x$ .

We consider the situation where an absorbing wall is put at the origin and, if the Brownian particle starting from  $x > 0$  arrives at the origin, it is absorbed there and the motion is stopped. Such a process is called the **absorbing Brownian motion in**  $\mathbb{R}_+$ . Its transition probability density is given by  $q_{\text{abs}}$ .

By absorption, the total mass of paths from  $x > 0$  to  $y > 0$  is then reduced, if we compare the original BM and the absorbing Brownian motion in  $\mathbb{R}_+$ . The factor  $y/x$  appearing in the transition probability density  $(1.54)$  of BES<sub>3</sub> is for renormalization so that

$$
\int_{\mathbb{R}_+} p_3(t, y|x) dy = 1, \quad \forall t > 0, \quad \forall x > 0.
$$
\n(1.56)

We regard this renormalization procedure from  $q_{\text{abs}}$  to  $p_3$  as a transformation. Since

$$
h_1(x) = x \tag{1.57}
$$

is a one-dimensional harmonic function in a rather trivial sense

<span id="page-14-2"></span>
$$
\Delta^{(1)}h_1(x) := \frac{d^2x}{dx^2} = 0,
$$

we say that the  $BES<sub>3</sub>$  is a **harmonic transformation** (h-transformation) of the one-dimensional absorbing BM in the sense of Doob. This implies the following equivalence.

> $BES<sub>3</sub> \leftrightarrow$  one-dimensional Brownian motion conditioned to stay positive  $\iff h$ -transform of absorbing BM in  $\mathbb{R}_+$

Let  $\mathbb{E}_{\text{BES}_3}^x$  denote the expectation with respect to BES<sub>3</sub>,  $(R(t))_{t\geq0}$ , started at  $x \in \mathbb{R}_+$ . For an independent BM,  $(B(t))_{t>0}$  started at the same point  $x \in \mathbb{R}_+$ , let  $\tau = \inf\{t > 0 : B(t) = 0\}$ . Then the above equivalence is written as follows; for any  $\mathcal{F}_t$ -measurable bounded function  $F, t \geq 0$ ,

$$
\mathbb{E}_{\text{BES}_3}^x[F(R(t))] = \mathbb{E}^x \left[ F(B(t)) \mathbf{1}_{(\tau > t)} \frac{B(t)}{x} \right]
$$
  
= 
$$
\mathbb{E}^x \left[ F(B(t)) \mathbf{1}_{(\tau > t)} \frac{h_1(B(t))}{h_1(x)} \right], \quad t \ge 0,
$$
 (1.58)

where  $E^x$  is an expectation with respect to BM started at  $x \in \mathbb{R}_+$ . If F is an even function;  $F(-x) = F(x)$ , then the above gives

<span id="page-14-3"></span>
$$
\mathbb{E}_{\text{BES}_3}^x[F(R(t))] = \mathbb{E}^x\left[F(B(t))\frac{h_1(B(t))}{h_1(x)}\right], \quad t \ge 0,
$$
\n(1.59)

since by the reflection principle of BM [\(1.5\)](#page-3-3), all contribution from paths  $\{\omega : \tau \leq t\}$  should be canceled out.

Here we emphasize the obvious fact that  $p_3(t,0|x) = 0$ ,  $\forall x > 0$ . It implies that BES<sub>3</sub> does not visit the origin. When  $D = 3$ , the outward drift is strong enough to avoid any visit to the origin. Moreover, we can prove that for any  $x > 0$ ,  $R^x(t) \to \infty$  as  $t \to \infty$  with probability 1 and we say the process is transient (see Theorem [1.2](#page-14-1) (ii) below).

#### <span id="page-14-0"></span>1.4.3 Critical dimension  $D_c = 2$

Originally, the Bessel process was defined by  $(1.42)$  for  $D \in \mathbb{N}$ . We find that, however, the modified Bessel function [\(1.49\)](#page-12-5) is an analytic function of  $\nu$  for all values of  $\nu$ . So we will be able to define the Bessel process for any value of dimension  $D \geq 1$  as a diffusion process in  $\mathbb{R}_+$  such that the transition probability density function is given by [\(1.52\)](#page-12-3), where the index  $\nu \ge -1/2$  is determined by [\(1.53\)](#page-12-2) for each value of  $D \geq 1$ .

For BES<sub>D</sub> starting from  $x > 0$ , denote its first visiting time at the origin by

$$
T^x = \inf\{t > 0 : R^x(t) = 0\}.
$$
\n(1.60)

<span id="page-14-1"></span>The following theorem is proved (see, for instance, [\[31\]](#page-88-1)).

**Theorem 1.2** (i)  $D \ge 2 \implies T^x = \infty,^{\forall} x > 0$ , with probability 1.

- (ii)  $D > 2 \implies \lim_{t \to \infty}$  $R^{x}(t) = \infty, \forall x > 0$ , with probability 1, *i.e.* the process is transient.
- (iii)  $D = 2 \implies \inf_{t>0} R^x(t) = 0, \forall x > 0$ , with probability 1D

That is,  $BES_{(2)}$  starting from  $x > 0$  does not visit the origin, but it can visit any neighborhood of the origin.

(iv)  $1 \leq D < 2 \implies T^x < \infty, \forall x > 0$ , with probability 1, *i.e.* the process is recurrent.

# <span id="page-15-0"></span>1.4.4 Bessel flow and another critical dimension  $\overline{D}_c = 3/2$

Theorem [1.2](#page-14-1) states that we have a **critical dimension**,

$$
D_{\rm c}=2,
$$

for competition between the two effects acting the Bessel process, the 'random force' (the martingale term) and the 'entropy force' (the outward drift term) in [\(1.45\)](#page-11-1): when  $D > D<sub>c</sub>$ , the latter dominates the former and the process becomes transient, and when  $D < D<sub>c</sub>$ , the former is relevant and recurrence to the origin of the process is realized frequently.

Here we show that there is another critical dimension [\[47\]](#page-89-0),

$$
\overline{D}_{\rm c} = \frac{3}{2}.
$$

In order to characterize the transition at  $\overline{D}_{c}$ , we have to investigate the dependence of the behavior of  $R^x(t)$  on its initial value,  $x > 0$ . We call the one-parameter family  $\{R^x(t) : t \ge 0\}_{x>0}$  the **Bessel** flow for each fixed  $D > 1$ .

For  $0 < x < y$ , we trace the motions of two BES<sub>D</sub>'s starting from x and y by solving [\(1.45\)](#page-11-1) using the *common* BM,  $B(t), t \geq 0$ ,

$$
R^{x}(t) = x + B(t) + \frac{D-1}{2} \int_{0}^{t} \frac{ds}{R^{x}(s)},
$$
  
\n
$$
R^{y}(t) = y + B(t) + \frac{D-1}{2} \int_{0}^{t} \frac{ds}{R^{y}(s)}, \quad 0 \le t < T^{x}.
$$
\n(1.61)

We will see that

$$
x < y \qquad \Longrightarrow \qquad R^x(t) < R^y(t), \quad 0 \le t < T^x \quad \text{with probability 1}
$$

$$
\qquad \Longrightarrow \qquad T^x \le T^y \quad \text{with probability 1}.
$$

The interesting fact is that in the intermediate fractional dimensions,  $\overline{D}_{c} < D < D_{c}$ , it is possible to have a situation where  $T^x = T^y$  even for  $x < y$ .

**Theorem 1.3** For  $0 < x < y < \infty$ ,

- (i)  $1 \le D \le 3/2 \implies T^x < T^y$  with probability 1.
- (ii)  $3/2 < D < 2 \implies P[T^x = T^y] > 0.$



Figure 2: In the intermediate fractional dimensions,  $3/2 < D < 2$ , there is a positive probability that two Bessel processes starting from different initial positions,  $0 < x < y < \infty$ , return to the origin simultaneously,  $T^x = T^y$ .

### <span id="page-16-1"></span><span id="page-16-0"></span>1.5 Hermitian-matrix-valued BM and the Dyson model with parameter  $\beta$  (DYS<sub>β</sub>)

#### 1.5.1 Multivariate extensions of Bessel processes

Here we consider the stochastic motion of two particles  $(X_1(t), X_2(t))$  in one dimension R satisfying the following SDEs,

<span id="page-16-2"></span>
$$
dX_1(t) = dB_1(t) + \frac{\beta}{2} \frac{dt}{X_1(t) - X_2(t)},
$$
  
\n
$$
dX_2(t) = dB_2(t) + \frac{\beta}{2} \frac{dt}{X_2(t) - X_1(t)},
$$
\n(1.62)

with the initial condition  $x_1 = X_1(0) < x_2 = X_2(0)$  for  $0 \le t < \inf\{t > 0 : X_1(t) = X_2(t)\},$ where  $B_1(t)$  and  $B_2(t)$ ,  $t \geq 0$  are independent BMs and  $\beta > 0$  is the 'coupling constant' of the two particles. The second terms in [\(1.62\)](#page-16-2) represent the repulsive force acting between two particles, which is proportional to the inverse of the distance between them,  $X_2(t) - X_1(t)$ . Since it is a central force  $(i.e.,$  depending only on distance, and thus symmetric for two particles), the 'center of mass'  $X_c(t) := (X_2(t) + X_1(t))/2$  is just a time change of BM; we can calculate the quadratic variation as  $d\langle X_c, X_c\rangle_t = \langle dX_c, dX_c\rangle_t = \langle (dB_1 + dB_2)/2, (dB_1 + dB_2)/2\rangle_t = dt/2$ , since  $d\langle B_1, B_1\rangle_t = \langle dB_1, dB_1\rangle_t = dt, d\langle B_2, B_2\rangle_t = \langle dB_2, dB_2\rangle_t = dt$  and  $d\langle B_1, B_2\rangle_t = \langle dB_1, dB_2\rangle_t = 0.$ Then

$$
(X_c(t))_{t\geq 0} \stackrel{\text{(law)}}{=} \left(\frac{1}{\sqrt{2}}B(t)\right)_{t\geq 0} \stackrel{\text{(law)}}{=} (B(t/2))_{t\geq 0},
$$

where  $B(t)$  is a BM independent from  $B_1(t)$  and  $B_2(t)$ . On the other hand, if we define the relative coordinate by  $X_{\rm r}(t) := (X_2(t) - X_1(t))/\sqrt{2}$ , it satisfies the SDE

$$
dX_{\rm r}(t) = d\widetilde{B}(t) + \frac{\beta}{2} \frac{dt}{X_{\rm r}(t)}
$$
\n(1.63)

for  $0 < t < \inf\{t > 0 : X_{r}(t) = 0\}$ , where  $\widetilde{B}(t), t \ge 0$  is a BM independent from  $B_1(t), B_2(t), B(t),$  $t \geq 0$ . It is nothing but the SDE for BES<sub>D</sub> with  $D = \beta + 1$ :

<span id="page-17-1"></span>
$$
(X_{r}(t))_{t\geq 0} \stackrel{\text{(law)}}{=} (R(t))_{t\geq 0}, \quad D = \beta + 1. \tag{1.64}
$$

Now we consider the following N-particle systems of **interacting Brownian motions** in  $\mathbb{R}$  as a solution  $\mathbf{X}(t) = (X_1(t), X_2(t), \dots, X_N(t))$  of the following system of SDEs: with  $\beta > 0$  and the condition  $x_1 < x_2 < \cdots < x_N$  for initial positions  $x_i = X_i(0), i = 1, \ldots, N$ ,

<span id="page-17-2"></span>
$$
dX_i(t) = dB_i(t) + \frac{\beta}{2} \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{dt}{X_i(t) - X_j(t)}, \quad t \in [0, T^x), \quad i = 1, ..., N,
$$
 (1.65)

where  ${B_i(t)}_{i=1}^N$ ,  $t \ge 0$  are independent BMs and

$$
T_{ij}^{\mathbf{x}} = \inf \{ t > 0 : X_i(t) = X_j(t) \}, \quad 1 \le i < j \le N, T^{\mathbf{x}} = \min_{1 \le i < j \le N} T_{ij}^{\mathbf{x}}.
$$

It is usually called Dyson's Brownian motion model with parameter  $\beta$ . But in this lecture we simply call it the **Dyson model with parameter**  $\beta$  and denoted by  $DYS_{\beta}$ .

As shown above, in the case where  $N = 2$ ,  $DYS_\beta$  is a composition of a BM (the center of mass) and a  $BES_{(\beta+1)}$  (the relative coordinate). In this sense,  $DYS_{\beta}$  can be regarded as a multivariate (multidimensional) extension of  $BES_{(\beta+1)}, \ \beta > 0.$ 

We can prove that, for any  $x \in \mathbb{R}^N$  with  $x_1 < x_2 < \cdots < x_N$ ,  $T^x < \infty$  if  $\beta < 1$ , and  $T^x = \infty$ if  $\beta \geq 1$  [\[65,](#page-90-0) [23\]](#page-88-2). The critical value  $\beta_c = 1$  seems to correspond to  $D_c = \beta_c + 1 = 2$  of BES<sub>D</sub>. Together with [\(1.64\)](#page-17-1), this observation suggests the following relation between parameters,

$$
D = \beta + 1 \quad \Longleftrightarrow \quad \beta = D - 1. \tag{1.66}
$$

In particular, we study the special case of Dyson's BM model with parameter  $\beta = 2$ , DYS<sub>2</sub> As shown above, the case where  $\beta = 2$  corresponds to a BES<sub>D</sub> with  $D = 3$ . In Section [1.4,](#page-10-0) we have shown that  $BES<sub>3</sub>$  has two aspects: [Aspect 1] as a radial coordinate of three-dimensional Brownian motion, which was used to define the Bessel process in Section [1.4.1,](#page-10-1) and [Aspect 2] as a one-dimensional Brownian motion conditioned to stay positive as explained in Section [1.4.2.](#page-12-0) We show that the Dyson model inherits these two aspects from  $BES<sub>3</sub>$  [\[31\]](#page-88-1).

#### <span id="page-17-0"></span> $1.5.2$  DYS<sub>2</sub> realized as eigenvalue process of Hermitian-matrix-valued BM

Dyson [\[19\]](#page-87-0) introduced the processes [\(1.65\)](#page-17-2) with  $\beta = 1, 2$ , and 4 as the eigenvalue processes of matrix-valued stochastic processes in order to realize the point processes in equilibrium called the Gaussian orthogonal ensemble (GOE), the Gaussian unitary ensemble (GUE), and the Gaussian symplectic ensemble  $(GSE)$ , which are extensively studied in **random matrix theory** [\[53,](#page-89-1) [21,](#page-88-3) [2,](#page-86-1) [1\]](#page-86-2).

For  $\beta = 2$  with given  $N \in \mathbb{N}$ , we prepare N-tuples of BMs  $\{B_{ii}^{x_i}(t)\}_{i=1}^N$ ,  $t \geq 0$ , each of which starts from  $x_i \in \mathbb{R}$ , and  $N(N-1)/2$ -tuples of pairs of BMs  $\{B_{ij}(t), \widetilde{B}_{ij}(t)\}_{1 \leq i < j \leq N}$ ,  $t \geq 0$ , starting from the origin. Here, there is a total of  $N + 2 \times N(N-1)/2 = N^2$  BMs, each of them independent from the rest. Then consider an  $N \times N$  **Hermitian-matrix-valued Brownian motion**,

<span id="page-18-0"></span>
$$
\mathbf{B}^{x}(t) = \begin{pmatrix} B_{11}^{x_1}(t) & \frac{B_{12}(t) + \sqrt{-1}\tilde{B}_{12}(t)}{\sqrt{2}} & \cdots & \frac{B_{1N}(t) + \sqrt{-1}\tilde{B}_{1N}(t)}{\sqrt{2}} \\ \frac{B_{12}(t) - \sqrt{-1}\tilde{B}_{12}(t)}{\sqrt{2}} & B_{22}^{x_2}(t) & \cdots & \frac{B_{2N}(t) + \sqrt{-1}\tilde{B}_{2N}(t)}{\sqrt{2}} \\ \cdots & \cdots & \cdots & \cdots \\ \frac{B_{1N}(t) - \sqrt{-1}\tilde{B}_{1N}(t)}{\sqrt{2}} & \frac{B_{2N}(t) - \sqrt{-1}\tilde{B}_{2N}(t)}{\sqrt{2}} & \cdots & B_{NN}^{x_N}(t) \end{pmatrix} .
$$
\n(1.67)

Remember that when we introduced  $BES<sub>D</sub>$  in Section [1.4,](#page-10-0) we considered the D-dimensional vector-valued Brownian motion in  $\mathbb{R}^N$ , [\(1.41\)](#page-10-3), by preparing D-tuples of independent BMs for its elements. Here we consider the space of  $N \times N$  Hermitian matrices denoted by  $\mathcal{H}(N)$ . Since the dimension of this space is  $\dim \mathcal{H}(N) = N^2$ , we need  $N^2$  independent BMs for elements to describe a Brownian motion in this space  $\mathcal{H}(N)$ . Hence we can regard the process  $\mathsf{B}^{\boldsymbol{x}}(t)$ ,  $t \geq 0$  defined by  $(1.67)$  as a 'Brownian motion in  $\mathcal{H}(N)$ '. By definition, the initial state of this Brownian motion is the diagonal matrix

<span id="page-18-1"></span>
$$
B^x(0) = diag(x_1, x_2, \dots, x_N).
$$
 (1.68)

We assume  $x_1 \leq x_2 \leq \cdots \leq x_N$ .

Corresponding to calculating the absolute value [\(1.42\)](#page-10-2) of  $\mathbf{B}^x(t)$ , by which  $BES_D$  was introduced, here we calculate the eigenvalues of  $B^{x}(t)$ . For any  $t \geq 0$ , there is a family of  $N \times N$  unitary matrices  $\{U(t)\}\$  which diagonalizes  $\mathsf{B}^{\boldsymbol{x}}(t)$ ,

$$
\mathsf{U}^*(t)\mathsf{B}^{\mathbf{x}}(t)\mathsf{U}(t) = \mathrm{diag}(\lambda_1(t),\ldots,\lambda_N(t)) =: \Lambda(t), \quad t \ge 0.
$$

Here for a matrix  $M = (M_{ij})_{1 \le i,j \le N}$ , we define its Hermitian conjugate by  $M^* = (\overline{M_{ji}})_{1 \le i,j \le N}$ , where  $\overline{z}$  denotes the complex conjugate of  $z \in \mathbb{C}$ . Consider a subspace of  $\mathbb{R}^N$  defined by

<span id="page-18-3"></span>
$$
\mathbb{W}_N := \{ \boldsymbol{x} = (x_1, x_2, \dots, x_N) \in \mathbb{R}^N : x_1 < x_2 < \dots < x_N \},\tag{1.69}
$$

which is called the **Weyl chamber** in representation theory. If we impose the condition  $(\lambda_i(t))_{i=1}^N$   $\in$  $W_N$ ,  $U(t)$  is uniquely determined at each time  $t \geq 0$ .

The following theorem is established.

**Theorem 1.4** The **eigenvalue process**  $(\lambda_i(t))_{i=1}^N$ ,  $t \ge 0$  of the Hermitian-matrix-valued Brownian motion [\(1.67\)](#page-18-0) started at [\(1.68\)](#page-18-1) satisfies the SDEs,

<span id="page-18-2"></span>
$$
d\lambda_i(t) = dB_i^{x_i}(t) + \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{dt}{\lambda_i(t) - \lambda_j(t)}, \quad t \ge 0, \quad i = 1, \dots, N,
$$
\n(1.70)

where  $(B_i^{x_i}(t))_{i=1}^N$ ,  $t \geq 0$  are independent BMs different from the  $N^2$ -tuples of BMs used to define  $B^{\mathbf{x}}(t)$  in [\(1.67\)](#page-18-0). That is, this process realizes DYS<sub>2</sub>.

The correspondence between  $BES_3$  and  $DYS_2$  in terms of equivalent processes is summarized as

follows.

[Aspect 1]

radial coordinate of BES<sub>3</sub>  $\qquad \Longleftrightarrow \qquad D=3$  vector-valued Brownian motion eigenvalue process of  $N$ -particle DYS<sub>2</sub>  $\iff$   $N \times N$  Hermitian-matrix-valued Brownian motion

Dyson derived [\(1.70\)](#page-18-2) by applying the perturbation theory in quantum mechanics [\[19\]](#page-87-0). Since  $(\lambda_i(t))_{i=1}^N$ ,  $t \geq 0$  are functionals of  $\{B_{ii}^{x_i}(t), B_{ij}(t), \widetilde{B}_{ij}(t)\}_{1 \leq i < j \leq N}$ ,  $t \geq 0$ , we can use Itô's formula to prove Theorem [1.4.](#page-18-3) A key point to prove the theorem is applying Itô's rule for differentiating the product of **matrix-valued semi-martingales** : If  $X(t) = (X_{ij}(t))$  and  $Y(t) = (Y_{ij}(t))$  are  $N \times N$  matrices with semi-martingale elements, then

$$
d(\mathsf{X}^*(t)\mathsf{Y}(t)) = d\mathsf{X}^*(t)\mathsf{Y}(t) + \mathsf{X}^*(t)d\mathsf{Y}(t) + \langle d\mathsf{X}^*, d\mathsf{Y} \rangle_t, \quad t \ge 0,
$$
\n(1.71)

.

where  $\langle dX^*, dY \rangle_t$  denotes an  $N \times N$  matrix-valued process, whose  $(i, j)$ -th element is given by the finite-variation process  $\sum_{k} \langle d\overline{X_{ki}}, dY_{kj} \rangle_t$ ,  $1 \leq i, j \leq N$ .

# <span id="page-19-1"></span><span id="page-19-0"></span>1.6 Determinantal martingale representation (DMR) of  $DYS_2$

#### 1.6.1 General formula

Denote the N-dimensional Laplacian with respect to the variables  $\mathbf{x} = (x_1, \ldots, x_N)$  by

$$
\Delta^{(N)}:=\sum_{i=1}^N\frac{\partial^2}{\partial x_i^2}
$$

We set

$$
h_N(\boldsymbol{x}) := \prod_{1 \le i < j \le N} (x_j - x_i) = \det_{1 \le i, j \le N} [x_j^{i-1}], \tag{1.72}
$$

where the determinant appearing here is called the **Vandermonde determinant**. As we have seen in [1.4.2,](#page-12-0)  $h_1(x) := x$  is a positive harmonic function in  $\mathbb{R}_+ = (0, \infty)$  which satisfies the boundary condition  $h_1(0) = 0$ . Similarly, we can see that

$$
\Delta^{(N)}h_N(\boldsymbol{x}) = 0. \tag{1.73}
$$

and

$$
h_N(\boldsymbol{x}) > 0 \text{, if } \boldsymbol{x} \in \mathbb{W}_N, \quad \text{and} \quad h_N(\boldsymbol{x}) = 0 \text{, if } \boldsymbol{x} \in \partial \mathbb{W}_N. \tag{1.74}
$$

We have another correspondence between  $BES<sub>3</sub>$  and  $DYS<sub>2</sub>$ .

[Aspect 2]



As [Aspect 2], the Dyson model is constructed as the h-transformation of the absorbing Brownian motion in  $\mathbb{W}_{N}$ . Therefore, at any positive time  $t > 0$  the configuration is an element of  $\mathbb{W}_N$ ,

$$
\mathbf{X}(t) = (X_1(t), X_2(t), \dots, X_N(t)) \in \mathbb{W}_N \quad t > 0,
$$
\n(1.75)

and hence there are no multiple points at which coincidence of particle positions,  $X_i(t) = X_j(t), i \neq j$  $j$ , occurs. We can consider, however, the Dyson model starting from initial configurations with multiple points. In order to describe configurations with multiple points, we represent each particle configuration by a sum of delta measures in the form

<span id="page-20-0"></span>
$$
\xi(\cdot) = \sum_{i \in \mathbb{I}} \delta_{x_i}(\cdot) \tag{1.76}
$$

with a sequence of points in  $\mathbb{R}, x = (x_i)_{i \in \mathbb{I}}$ , where  $\mathbb{I}$  is a countable index set. Here for  $y \in \mathbb{R}, \delta_y(\cdot)$ denotes the delta measure such that  $\delta_y({x}) = 1$  for  $x = y$  and  $\delta_y({x}) = 0$  otherwise. Then, for [\(1.76\)](#page-20-0) and  $A \subset \mathbb{R}$ ,  $\xi(A) = \int_A \xi(dx) = \sum_{i \in \mathbb{I}: x_i \in A} 1 = \sharp \{x_i, x_i \in A\}.$ 

The measures of the form [\(1.76\)](#page-20-0) satisfying the condition  $\xi(K) < \infty$  for any compact subset  $K \subset \mathbb{R}$  are called the **nonnegative integer-valued Radon measures** on  $\mathbb{R}$  and we denote the space they form by  $Conf(\mathbb{R});$ 

$$
\text{Conf}(\mathbb{R}) := \left\{ \xi = \sum_{i \in \mathbb{I}} \delta_{x_i} : x_i \in \mathbb{R}, \xi(K) < \infty \text{ for all bounded set } K \subset \mathbb{R} \right\}. \tag{1.77}
$$

The set of configurations without multiple points is denoted by

$$
Conf_0(\mathbb{R}) := \{ \xi \in Conf(\mathbb{R}) : \xi(\lbrace x \rbrace) \leq 1,^{\forall} x \in \mathbb{R} \}.
$$

There is a trivial correspondence between  $\mathbb{W}_N$  and  $\text{Conf}_0(\mathbb{R})$ . We call  $\boldsymbol{x} \in \mathbb{R}^N$  a labeled configuration and  $\xi \in \text{Conf}(\mathbb{R})$  an unlabeled configuration.

We introduce a sequence of independent BMs,  $\mathbf{B}^{\mathbf{x}}(t) = (B_i^{x_i}(t))_{i \in \mathbb{I}}$ ,  $t \geq 0$ , in  $(\Omega, \mathcal{F}, P^{\mathbf{x}})$  with expectation written as  $E^x$ .

In this section we assume that  $\xi = \sum_{i \in \mathbb{I}} \delta_{x_i} \in \text{Conf}_0(\mathbb{R}), \xi(\mathbb{R}) = N \in \mathbb{N}$  and consider  $DYS_2$  as an  $\text{Conf}_0(\mathbb{R})$ -valued diffusion process,

<span id="page-20-1"></span>
$$
\Xi(t,\cdot) = \sum_{i=1}^{N} \delta_{X_i(t)}(\cdot), \quad t \ge 0,
$$
\n(1.78)

starting from the initial configuration  $\xi = \sum_{i=1}^{N} \delta_{x_i}$ , where  $\boldsymbol{X}(t) = (X_1(t), \cdots, X_N(t))$  is the solution of [\(1.65\)](#page-17-2) with  $\beta = 2$  under the initial configuration  $x = (x_1, \dots, x_N) \in W_N$ . We write the process as  $(\Xi, \mathbb{P}^{\xi})$  and express the expectation with respect to the probability law  $\mathbb{P}^{\xi}$  of the Dyson model by  $\mathbb{E}^{\xi}[\cdot]$ . We introduce a filtration  $\{(\mathcal{F}_{\Xi})_t\}_{t\in[0,\infty)}$  on the space of continuous paths  $C([0,\infty) \to \text{Conf}(\mathbb{R}))$  defined by  $(\mathcal{F}_{\Xi})_t = \sigma(\Xi(s), s \in [0,t])$ , where  $\sigma$  denotes the smallest  $\sigma$ -field.

[Aspect 2] of the Dyson model is expressed by the following equality: for any  $(\mathcal{F}_{\Xi})_t$ -measurable bounded function  $F, 0 \le t \le T < \infty$ ,

<span id="page-20-2"></span>
$$
\mathbb{E}^{\xi}[F(\Xi(\cdot))] = \mathbb{E}^{\boldsymbol{x}}\left[F\left(\sum_{i=1}^{N} \delta_{B_{i}(\cdot)}\right) \mathbf{1}_{(\tau > T)} \frac{h_{N}(\boldsymbol{B}(T))}{h_{N}(\boldsymbol{x})}\right],\tag{1.79}
$$

where

$$
\tau := \inf\{t > 0 : \mathbf{B}^{\mathbf{x}}(t) \notin \mathbb{W}_N\}
$$
\n(1.80)

and we have assumed the relations  $\xi = \sum_{i=1}^{N} \delta_{x_i} \in \text{Conf}_0(\mathbb{R}), \boldsymbol{x} = (x_1, \dots, x_N) \in \mathbb{W}_N$  and [\(1.78\)](#page-20-1).

In the following lemma, we claim that even if we delete the indicator  $\mathbf{1}_{(\tau>T)}$  in RHS of [\(1.79\)](#page-20-2), still the equality holds. It is a multivariate extension of the claim by which we replaced [\(1.58\)](#page-14-2) by [\(1.59\)](#page-14-3) in Section [1.4.2.](#page-12-0)

**Lemma 1.5** Assume that  $\xi = \sum_{i=1}^{N} \delta_{x_i} \in \text{Conf}_0(\mathbb{R}), x = (x_1, \ldots, x_N) \in \mathbb{W}_N$ . For any  $(\mathcal{F}_{\Xi})_t$ measurable bounded function  $F$ ,  $0 \le t \le T < \infty$ ,

$$
\mathbb{E}^{\xi}[F(\Xi(\cdot))] = \mathbb{E}^{\boldsymbol{x}}\left[F\left(\sum_{i=1}^{N} \delta_{B_i(\cdot)}\right) \frac{h_N(\boldsymbol{B}(T))}{h_N(\boldsymbol{x})}\right].
$$
\n(1.81)

In Section [1.3,](#page-8-0) we introduced the fundamental martingale polynomials associated with BM,  ${m_n(t, x)}_{n \in \mathbb{N}_0}$ . Since they are monic polynomials, we can prove the equalities

<span id="page-21-0"></span>
$$
\frac{h_N(\mathbf{y})}{h_N(\mathbf{x})} = \frac{1}{h_N(\mathbf{x})} \det_{1 \le i,j \le N} [y_j^{i-1}]
$$
\n
$$
= \frac{1}{h_N(\mathbf{x})} \det_{1 \le i,j \le N} [m_{i-1}(t, y_j)] \tag{1.82}
$$

for an arbitrary  $t \in [0,\infty)$ . This implies that  $(h_N(\mathbf{B}(t))/h_N(\mathbf{x}))_{t>0}$  is a martingale.

Here we extend the integral transformation defined by [\(1.39\)](#page-10-4) with [\(1.36\)](#page-9-2) in Section [1.3](#page-8-0) to a linear integral transformation of multivariate functions as follows. When  $F^{(i)}(\boldsymbol{x}) = \prod_{j=1}^{N} f_j^{(i)}$  $j^{(i)}(x_j),$  $i = 1, 2$  are given for  $\boldsymbol{x} = (x_1, \dots, x_N) \in \mathbb{R}^N$ , then we define

$$
\mathcal{I}\left[F^{(i)}(\mathbf{W})\left|\{(t_{\ell},x_{\ell})\}_{\ell=1}^{N}\right] := \prod_{j=1}^{N} \mathcal{I}\left[f_{j}^{(i)}(W_{j})\right](t_{j},x_{j})\right], \quad i=1,2,
$$

and

$$
\mathcal{I}\Big[c_1 F^{(1)}(\mathbf{W}) + c_2 F^{(2)}(\mathbf{W}) | \{ (t_\ell, x_\ell) \}_{\ell=1}^N \Big] \n:= c_1 \mathcal{I}\Big[F^{(1)}(\mathbf{W}) | \{ (t_\ell, x_\ell) \}_{\ell=1}^N \Big] + c_2 \mathcal{I}\Big[F^{(2)}(\mathbf{W}) | \{ (t_\ell, x_\ell) \}_{\ell=1}^N \Big],
$$

 $c_1, c_2 \in \mathbb{C}$ , for  $0 \le t_i \le \infty$ ,  $i = 1, \ldots, N$ , where  $\mathbf{W} = (W_1, \ldots, W_N) \in \mathbb{R}^N$ . In particular, if  $t_{\ell} = t, 1 \leq \forall \ell \leq N$ , we write  $\mathcal{I}[\cdot | \{(t_{\ell}, x_{\ell})\}_{\ell=1}^{N}]$  simply as  $\mathcal{I}[\cdot | (t, x)]$  with  $\boldsymbol{x} = (x_1, \ldots, x_N)$ . Then [\(1.82\)](#page-21-0) is further rewritten as

<span id="page-21-1"></span>
$$
\frac{h_N(\boldsymbol{y})}{h_N(\boldsymbol{x})} = \frac{1}{h_N(\boldsymbol{x})} \det_{1 \leq i,j \leq N} \left[ \mathcal{I}[(W_j)^{i-1} | (t, y_j)] \right]
$$
\n
$$
= \mathcal{I} \left[ \frac{1}{h_N(\boldsymbol{x})} \det_{1 \leq i,j \leq N} [(W_j)^{i-1}] \middle| (t, \boldsymbol{y}) \right]
$$
\n
$$
= \mathcal{I} \left[ \frac{h_N(\mathbf{W})}{h_N(\boldsymbol{x})} \middle| (t, \boldsymbol{y}) \right], \tag{1.83}
$$

where the multilinearity of determinants has been used.

Now we use the following determinant identity.

**Lemma 1.6** For  $x = (x_1, ..., x_N) \in W_N$ ,  $z = (z_1, ..., z_N) \in \mathbb{C}^N$ ,

$$
\frac{h_N(z)}{h_N(x)} = \det_{1 \le i,j \le N} \left[ \Phi_{\xi}^{x_i}(z_j) \right],\tag{1.84}
$$

where

<span id="page-22-4"></span>
$$
\Phi_{\xi}^{u}(z) = \prod_{\substack{1 \le k \le N, \\ x_k \ne u}} \frac{z - x_k}{u - x_k}, \quad \text{for } \xi = \sum_{i=1}^{N} \delta_{x_i} \in \text{Conf}_0(\mathbb{R}), z, u \in \mathbb{C}. \tag{1.85}
$$

Then [\(1.83\)](#page-21-1) is written as

$$
\frac{h_N(\boldsymbol{y})}{h_N(\boldsymbol{x})} = \mathcal{I}\left[\det_{1 \le i,j \le N}[\Phi_{\xi}^{x_i}(W_j)]\middle|(t,\boldsymbol{y})\right] \n= \det_{1 \le i,j \le N}[\mathcal{M}_{\xi}^{x_i}(t,y_j)],
$$
\n(1.86)

where

<span id="page-22-0"></span> $\mathcal{M}_{\xi}^{x}(t, y) := \mathcal{I}[\Phi_{\xi}^{x}(W)| (t, y)], \quad x, y \in \mathbb{R}, \quad t \ge 0.$  (1.87)

<span id="page-22-1"></span>**Proposition 1.7** Assume that  $\xi = \sum_{i=1}^{N} \delta_{x_i} \in \text{Conf}_0(\mathbb{R})$ . The following are satisfied by [\(1.87\)](#page-22-0).

- (i)  $(\mathcal{M}_{\xi}^{x_i}(t, B(t)))_{t \geq 0}, i = 1, \ldots N$  are continuous martingales.
- (ii) For any time  $t \geq 0$ ,  $\mathcal{M}_{\xi}^{x_i}(t, y), i = 1, ..., N$  are linearly independent functions of y.
- (iii) For  $1 \leq i, j \leq N$ ,  $\lim_{t \downarrow 0} \mathbb{E}^{x_i}[\mathcal{M}_{\xi}^{x_j}(t, B(t))] = \delta_{ij}$ .

Then for any  $(\mathcal{F}_{\Xi})_t$ -measurable bounded function  $F, 0 \le t \le T < \infty$ , the equality

<span id="page-22-2"></span>
$$
\mathbb{E}^{\xi}[F(\Xi(\cdot))] = \mathbb{E}^{\boldsymbol{x}}\left[F\left(\sum_{i=1}^{N} \delta_{B_{i}(\cdot)}\right) \mathcal{D}_{\xi}(T, \boldsymbol{B}(T))\right]
$$
(1.88)

holds, where

$$
\mathcal{D}_{\xi}(t, \mathbf{y}) = \det_{1 \leq i, j \leq N} [\mathcal{M}_{\xi}^{y_j}(t, y_i)], \quad \mathbf{y} = (y_1, \dots, y_N) \in \mathbb{W}_N, \quad t \geq 0.
$$
 (1.89)

We remark that  $\mathcal{D}_{\xi}(t, \mathbf{B}(t)), t \geq 0$  is indeed a continuous martingale by part (i) and is not identically zero by part (ii) of Proposition [1.7.](#page-22-1) We call  $\mathcal{D}_{\xi}(t, \boldsymbol{B}(t))$ ,  $t \geq 0$  a determinantal martingale and the equality [\(1.88\)](#page-22-2) the determinantal martingale representation (DMR) of  $DYS_2$  [\[30\]](#page-88-0).

For  $n \in \mathbb{N}$ , an index set  $\{1, 2, \ldots, n\}$  is denoted by  $\mathbb{I}_n$ . Fixing  $N \in \mathbb{N}$  with  $N' \in \mathbb{I}_N$ , we write  $\mathbb{J} \subset \mathbb{I}_N, \sharp \mathbb{J} = N'$ , if  $\mathbb{J} = \{j_1, \ldots, j_{N'}\}, 1 \leq j_1 < \cdots < j_{N'} \leq N$ . For  $\boldsymbol{x} = (x_1, \ldots, x_N) \in \mathbb{R}^N$ , put  $\boldsymbol{x}_{\mathbb{J}}=(x_{j_1},\ldots,x_{j_{N'}}).$  In particular, we write  $\boldsymbol{x}_{N'}=\boldsymbol{x}_{\mathbb{I}_{N'}}, 1\leq N'\leq N.$  (By definition  $\boldsymbol{x}_N=\boldsymbol{x}.$ ) A collection of all permutations of elements in  $\mathbb J$  is denoted by  $\mathbb S(\mathbb J)$ . In particular, we write  $\mathbb{S}_{N'} = \mathbb{S}(\mathbb{I}_{N'}), 1 \leq N' \leq N.$ 

<span id="page-22-3"></span>The following shows the reducibility of the determinantal martingale in the sense that, if we observe a symmetric function depending on N' variables,  $N' \leq N$ , then the size of determinantal martingale can be reduced from  $N$  to  $N'$ .

**Lemma 1.8** Assume that  $\xi = \sum_{i=1}^{N} \delta_{x_i} = \sum_{i \in I_N} \delta_{x_i}$  with  $x \in W_N$ . Let  $1 \leq N' \leq N$ . For  $0 < t \leq T < \infty$  and an  $(\mathcal{F}_{\Xi})_t$ -measurable symmetric function  $F_{N'}$  on  $\mathbb{R}^{N'}$ ,

$$
\sum_{\mathbb{J}\subset\mathbb{I}_N,\sharp\mathbb{J}=N'} \mathbb{E}^{\mathbf{x}} \left[ F_{N'}(\mathbf{B}_{\mathbb{J}}(t)) \mathcal{D}_{\xi}(T,\mathbf{B}(T)) \right]
$$
\n
$$
= \int_{\mathbb{W}_{N'}} \xi^{\otimes N'}(d\mathbf{v}) \mathbb{E}^{\mathbf{v}} \left[ F_{N'}(\mathbf{B}_{N'}(t)) \mathcal{D}_{\xi}(T,\mathbf{B}_{N'}(T)) \right]. \tag{1.90}
$$

# <span id="page-23-0"></span>1.6.2 Time-dependent density function  $\rho_{\xi}(t, x)$

The **density function** at a single time for  $(\Xi, \mathbb{P}^{\xi}), \xi \in \text{Conf}_0(\mathbb{R})$  is denoted by  $\rho_{\xi}(t, x)$ . It is defined as a continuous function of  $x \in \mathbb{R}$  for  $0 \le t \le T < \infty$  such that for any **test function**,  $\chi \in C_c(\mathbb{R})$ ,

<span id="page-23-2"></span>
$$
\mathbb{E}^{\xi}\left[\int_{\mathbb{R}}\chi(x)\Xi(t,dx)\right]=\int_{\mathbb{R}}dx\,\chi(x)\rho_{\xi}(t,x). \tag{1.91}
$$

The test function  $\chi$  is symmetrized as  $g(x) = \sum_{n=1}^{N}$  $i=1$  $\chi(x_i)$ , which is applied as F to the DMR [\(1.88\)](#page-22-2), and we obtain the equality

<span id="page-23-1"></span>
$$
\mathbb{E}^{\xi}\left[\sum_{i=1}^{N}\chi(X_i(t))\right] = \mathbb{E}^{\boldsymbol{x}}\left[\sum_{i=1}^{N}\chi(B_i(t))\mathcal{D}_{\xi}(T,\boldsymbol{B}(T))\right], \quad 0 \le t \le T < \infty.
$$
 (1.92)

The LHS of [\(1.92\)](#page-23-1) gives

$$
\mathbb{E}^{\xi}\left[\sum_{i=1}^{N} \chi(X_i(t))\right] = \mathbb{E}^{\xi}\left[\int_{\mathbb{R}} \chi(x)\Xi(t, dx)\right]
$$

by [\(1.78\)](#page-20-1). On the other hand, RHS of [\(1.92\)](#page-23-1) is reduced by Lemma [1.8](#page-22-3) as

$$
\sum_{i=1}^{N} \mathbf{E}^{\mathbf{x}}[\chi(B_i(t))\mathcal{D}_{\xi}(T,\mathbf{B}(T))] = \int_{\mathbb{R}} \xi(dv)\mathbf{E}^v[\chi(B(t))\mathcal{M}_{\xi}^v(t,B(t))]
$$
  

$$
= \int_{\mathbb{R}} \xi(dv)\int_{\mathbb{R}} dx \chi(x)p(t,x|v)\mathcal{M}_{\xi}^v(t,x).
$$

By Fubini's theorem, we can rewrite it as  $\int_{\mathbb{R}} dx \,\chi(x) \mathcal{G}_{\xi}(t, x; t, x)$ , where

<span id="page-23-3"></span>
$$
\mathcal{G}_{\xi}(s,x;t,y) = \int_{\mathbb{R}} \xi(dv)p(s,x|v) \mathcal{M}_{\xi}^{v}(t,y). \tag{1.93}
$$

Then [\(1.91\)](#page-23-2) gives

$$
\rho_{\xi}(t,x) = \mathcal{G}_{\xi}(t,x;t,x), \quad x \in \mathbb{R}, \quad t \ge 0.
$$
\n(1.94)

# <span id="page-24-0"></span>**1.6.3** Two-time correlation function  $\rho_{\xi}(s, x; t, y)$

For  $0\leq t_1< t_2\leq T<\infty,$  set

$$
g_1(\boldsymbol{x}) = \sum_{i=1}^N \chi_1(x_i), \quad g_2(\boldsymbol{x}) = \sum_{i=1}^N \chi_2(x_i),
$$

where  $\chi_m \in C_c(\mathbb{R}), m = 1, 2$ , and put

$$
F(\Xi(\cdot)) = \prod_{m=1}^{2} g_m(\mathbf{X}(t_m)).
$$

If we apply this to DMR, [\(1.88\)](#page-22-2), we obtain the equality

<span id="page-24-1"></span>
$$
\mathbb{E}^{\xi} \left[ \sum_{i=1}^{N} \sum_{j=1}^{N} \chi_{1}(X_{i}(t_{1})) \chi_{2}(X_{j}(t_{2})) \right]
$$
  
= 
$$
\mathbb{E}^{x} \left[ \sum_{i=1}^{N} \sum_{j=1}^{N} \chi_{1}(B_{i}(t_{1})) \chi_{2}(B_{j}(t_{2})) \mathcal{D}_{\xi}(T, \mathbf{B}(T)) \right], 0 \le t \le T < \infty.
$$
 (1.95)

The LHS of [\(1.95\)](#page-24-1) defines the two-time correlation function  $\rho_{\xi}(s,x;t,y)$  as

<span id="page-24-2"></span>
$$
\mathbb{E}^{\xi}\left[\sum_{i=1}^{N}\sum_{j=1}^{N}\chi_{1}(X_{i}(t_{1}))\chi_{2}(X_{j}(t_{2}))\right]=\int_{\mathbb{R}^{2}}dx_{1}dx_{2}\,\chi_{1}(x_{1})\chi_{2}(x_{2})\rho_{\xi}(t_{1},x_{1};t_{2},x_{2}).\tag{1.96}
$$

On the other hand, RHS of [\(1.95\)](#page-24-1) gives

$$
\sum_{i=1}^{N} \sum_{j=1}^{N} \mathbf{E}^{\mathbf{x}}[\chi_1(B_i(t_1))\chi_2(B_j(t_2))\mathcal{D}_{\xi}(T,\mathbf{B}(T))]
$$
\n
$$
= \sum_{\substack{1 \le i,j \le N, \\ i \ne j}} \mathbf{E}^{\mathbf{x}}[\chi_1(B_i(t_1))\chi_2(B_j(t_2))\mathcal{D}_{\xi}(T,\mathbf{B}(T))]
$$
\n
$$
+ \sum_{1 \le i \le N} \mathbf{E}^{\mathbf{x}}[\chi_1(B_i(t_1))\chi_2(B_i(t_2))\mathcal{D}_{\xi}(T,\mathbf{B}(T))].
$$

By the reducibility of DMR given by Lemma [1.8,](#page-22-3) the last expression becomes

$$
\int_{\mathbb{R}^2} \xi^{\otimes 2}(d\mathbf{v})
$$
\n
$$
\times \mathbf{E}^{(v_1, v_2)} \left[ \chi_1(B_1(t_1)) \chi_2(B_2(t_2)) \det \begin{pmatrix} \mathcal{M}_{\xi}^{v_1}(T, B_1(T)) & \mathcal{M}_{\xi}^{v_1}(T, B_2(T)) \\ \mathcal{M}_{\xi}^{v_2}(T, B_1(T)) & \mathcal{M}_{\xi}^{v_2}(T, B_2(T)) \end{pmatrix} \right]
$$
\n
$$
+ \int_{\mathbb{R}} \xi(dv) \mathbf{E}^v[\chi_1(B(t_1)) \chi_2(B(t_2)) \mathcal{M}_{\xi}^v(T, B(T))]
$$

If we use the martingale property (i) of  $\mathcal{M}_{\xi}^{v}$  in Proposition [1.7,](#page-22-1) it is written as

$$
\int_{\mathbb{R}^2} \xi^{\otimes 2}(dv)
$$
\n
$$
\times \mathbf{E}^{(v_1,v_2)} \left[ \chi_1(B_1(t_1)) \chi_2(B_2(t_2)) \det \begin{pmatrix} \mathcal{M}_{\xi}^{v_1}(t_1, B_1(t_1)) & \mathcal{M}_{\xi}^{v_1}(t_2, B_2(t_2)) \\ \mathcal{M}_{\xi}^{v_2}(t_1, B_1(t_1)) & \mathcal{M}_{\xi}^{v_2}(t_2, B_2(t_2)) \end{pmatrix} \right]
$$
\n
$$
+ \int_{\mathbb{R}} \xi(dv) \mathbf{E}^v[\chi_1(B(t_1)) \chi_2(B(t_2)) \mathcal{M}_{\xi}^v(t_2, B(t_2))].
$$

By Fubini's theorem, this is equal to

$$
\int_{\mathbb{R}^2} dx_1 dx_2 \chi_1(x_1) \chi_2(x_2) \det \begin{pmatrix} \mathcal{G}_{\xi}(t_1, x_1; t_1, x_1) & \mathcal{G}_{\xi}(t_1, x_1; t_2, x_2) \\ \mathcal{G}_{\xi}(t_2, x_2; t_1, x_1) & \mathcal{G}_{\xi}(t_2, x_2; t_2, x_2) \end{pmatrix} + \int_{\mathbb{R}^2} dx_1 dx_2 \chi_1(x_1) \chi_2(x_2) \mathcal{G}_{\xi}(t_1, x_1; t_2, x_2) p(t_2 - t_1, x_2 | x_1)
$$
\n
$$
= \int_{\mathbb{R}^2} dx_1 dx_2 \chi_1(x_1) \chi_2(x_2)
$$
\n
$$
\times \det \begin{pmatrix} \mathcal{G}_{\xi}(t_1, x_1; t_1, x_1) & \mathcal{G}_{\xi}(t_1, x_1; t_2, x_2) \\ \mathcal{G}_{\xi}(t_2, x_2; t_1, x_1) - p(t_2 - t_1, x_2 | x_1) & \mathcal{G}_{\xi}(t_2, x_2; t_2, x_2) \end{pmatrix}.
$$

Since this is equal to [\(1.96\)](#page-24-2), the two-time correlation function is determined as

$$
\rho_{\xi}(s,x;t,y) = \det \left( \begin{array}{cc} \mathbb{K}_{\xi}(s,x;s,x) & \mathbb{K}_{\xi}(s,x;t,y) \\ \mathbb{K}_{\xi}(t,y;s,x) & \mathbb{K}_{\xi}(t,y;t,y) \end{array} \right)
$$
(1.97)

for  $0 \leq s < t < \infty$ ,  $x, y \in \mathbb{R}$ , where

<span id="page-25-3"></span>
$$
\mathbb{K}_{\xi}(s,x;t,y) = \mathcal{G}_{\xi}(s,x;t,y) - \mathbf{1}_{(s>t)} p(s-t,x|y). \tag{1.98}
$$

## <span id="page-25-1"></span><span id="page-25-0"></span>1.7 Determinantal stochastic processes (DSPs)

#### 1.7.1 From DMR to DSP

In the previous section, the density function at a single time  $\rho_{\xi}(t,x)$  and the two-time (and twopoint) correlation function  $\rho_{\xi}(s, x; t, y)$  were defined by [\(1.91\)](#page-23-2) and [\(1.96\)](#page-24-2), respectively. In order to give a general definition of spatio-temporal correlation functions here we consider the Laplace transformations of the multitime joint distribution functions of  $(\Xi, \mathbb{P}^{\xi})$ . For any integer  $M \in \mathbb{N}$ , a sequence of times  $\mathbf{t} = (t_1, \ldots, t_M) \in [0, \infty)^M$  with  $0 \le t_1 < \cdots < t_M < \infty$ , and a sequence of functions  $\boldsymbol{f} = (f_{t_1}, \dots, f_{t_M}) \in C_c(\mathbb{R})^M$ , let

<span id="page-25-2"></span>
$$
\Psi_{\xi}^{\mathbf{t}}[\mathbf{f}] \equiv \mathbb{E}^{\xi} \left[ \exp \left\{ \sum_{m=1}^{M} \int_{\mathbb{R}} f_{t_m}(x) \Xi(t_m, dx) \right\} \right]. \tag{1.99}
$$

By [\(1.78\)](#page-20-1), if we set test functions as

$$
\chi_{t_m}(\cdot) = e^{f_{t_m}(\cdot)} - 1, \quad 1 \le m \le M,\tag{1.100}
$$

we can rewrite [\(1.99\)](#page-25-2) in the form

<span id="page-26-2"></span>
$$
\Psi_{\xi}^{\mathbf{t}}[\mathbf{f}] = \mathbb{E}^{\xi} \left[ \prod_{m=1}^{M} \prod_{i=1}^{N} \{ 1 + \chi_{t_m}(X_i(t_m)) \} \right]. \tag{1.101}
$$

We expand this with respect to test functions and define the spatio-temporal correlation functions  $\{\rho_{\xi}\}\$ as coefficients,

$$
\Psi_{\xi}^{\boldsymbol{t}}[\boldsymbol{f}] = \sum_{\substack{0 \le N_m \le N, \\ 1 \le m \le M}} \int_{\prod_{m=1}^M \mathbb{W}_{N_m^{\mathbf{A}}} m} \prod_{m=1}^M dx_{N_m}^{(m)} \prod_{i=1}^{N_m} \chi_{t_m}\left(x_i^{(m)}\right) \rho_{\xi}\left(t_1, x_{N_1}^{(1)}; \ldots; t_M, x_{N_M}^{(M)}\right), \tag{1.102}
$$

where  $\boldsymbol{x}_{N_{m}}^{(m)}$  $\binom{m}{N_m}$  denotes  $\binom{m}{1}$  $x_1^{(m)},\ldots,x_{N_m}^{(m)}$  $\binom{(m)}{N_m}$  and  $d\boldsymbol{x}^{(m)}_{N_m}$  $\mathcal{L}_{N_m}^{(m)} = \prod_{i=1}^{N_m} dx_i^{(m)}, \ 1 \leq m \leq M.$  Here the empty products equal 1 by convention and the term with  $N_m = 0, 1 \leq \forall m \leq M$  is considered to be 1. The previous two examples  $\rho_{\xi}(t, x)$  and  $\rho_{\xi}(s, x; t, y)$  are special cases in which we set  $M = 1, t_1 =$ t,  $N_1 = 1$ ,  $x_1^{(1)} = x$ , and  $M = 2$ ,  $t_1 = s$ ,  $t_2 = t$ ,  $N_1 = N_2 = 1$ ,  $x_1^{(1)} = x$ ,  $x_1^{(2)} = y$ , respectively. The function  $\Psi_{\xi}^{t}[f]$  is a generating function of correlation functions.

Given an integral kernel,  $\mathbf{K}(s, x; t, y)$ ,  $(s, x)$ ,  $(t, y) \in [0, \infty) \times \mathbb{R}$ , and a sequence of functions  $(\chi_{t_1},\ldots,\chi_{t_M})\in \mathrm{C}_c(\mathbb{R})^M$ ,  $M\in\mathbb{N}$ , the **Fredholm determinant** associated with **K** and  $(\chi_{t_m})_{m=1}^M$ is defined as

$$
\text{Det}_{(s,t) \in \{t_1, \ldots, t_M\}^2, \left[\delta_{st}\delta_x(\{y\}) + \mathbf{K}(s, x; t, y)\chi_t(y)\right]}\n= \sum_{\substack{0 \le N_m \le N, \\ 1 \le m \le M}} \int_{\prod_{m=1}^M \mathbb{W}_{N_m}^{\mathbf{A}}} \prod_{m=1}^M dx_{N_m}^{(m)} \prod_{k=1}^{N_m} \chi_{t_m}\left(x_k^{(m)}\right) \det_{\substack{1 \le i \le N_m, 1 \le j \le N_n, \\ 1 \le m, n \le M}} \left[\mathbf{K}(t_m, x_i^{(m)}; t_n, x_j^{(n)})\right].
$$
\n(1.103)

<span id="page-26-1"></span>Definition 1.9 If any moment generating function [\(1.99\)](#page-25-2) is given by a Fredholm determinant, the process  $(\Xi, \mathbb{P}^{\xi})$  is said to be a **determinantal stochastic process** (DSP). In this case all spatio-temporal correlation functions are given by determinants as

<span id="page-26-3"></span>
$$
\rho_{\xi}\left(t_{1}, \boldsymbol{x}_{N_{1}}^{(1)}; \ldots; t_{M}, \boldsymbol{x}_{N_{M}}^{(M)}\right) = \det_{\substack{1 \leq i \leq N_{m}, 1 \leq j \leq N_{n}, \\1 \leq m, n \leq M}} \left[\mathbb{K}_{\xi}(t_{m}, x_{i}^{(m)}; t_{n}, x_{j}^{(n)})\right],\tag{1.104}
$$

 $0 \le t_1 < \cdots < t_M < \infty, \ 1 \le N_m \le N, \ \pmb{x}_{N_m}^{(m)}$  $\mathbb{R}_{N_m}^{(m)} \in \mathbb{R}^{N_m}, 1 \leq m \leq M \in \mathbb{N}$ . Here the integral kernel,  $\mathbb{K}_{\xi}: ([0,\infty) \times \mathbb{R})^2 \mapsto \mathbb{R}$ , is a function of the initial configuration  $\xi$  and is called the **correlation** kernel.

**Remark 1.10** If the process  $(\Xi, \mathbb{P}^{\xi})$  is DSP, then, at each time  $0 \le t < \infty$ , all spatial correlation functions are given by determinants as

<span id="page-26-0"></span>
$$
\rho_{\xi}(\boldsymbol{x}_{N'}) = \det_{1 \le i,j \le N'} [K(x_i, x_j)], \quad 1 \le N' \le N,
$$
\n(1.105)

with  $K(x, y) = \mathbb{K}_{\xi}(t, x; t, y)$ . In general a random integer-valued Radon measure in Conf(R) (resp.  $Conf_0(\mathbb{R})$  is called a **point process** (resp. **simple point process**). A simple point process is said to be a determinantal point process  $(DPP)$  (or Fermion point process) with kernel K, if its spatial correlation functions exist and are given in the form [\(1.105\)](#page-26-0). When K is symmetric, *i.e.*,  $K(x, y) = K(y, x), x, y \in \mathbb{R}$ , Soshnikov [\[81\]](#page-91-0) and Shirai and Takahashi [\[76,](#page-91-1) [77,](#page-91-2) [78\]](#page-91-3) gave sufficient conditions for K to be a correlation kernel of a determinantal point process. See also [\[35\]](#page-88-4). The notion of DSP given by Definition [1.9](#page-26-1) is a dynamical extension of the determinantal point process [\[10,](#page-87-1) [38\]](#page-89-2).

Let  $\Xi$  and  $\widetilde{\Xi}$  be determinantal processes with correlation kernels K and K, respectively. If there is a function  $G(s, x)$ , which is continuous with respect to  $x \in \mathbb{R}$  for any fixed  $s \in [0, \infty)$ , such that

<span id="page-27-0"></span>
$$
\mathbb{K}(s,x;t,y) = \frac{G(s,x)}{G(t,y)} \widetilde{\mathbb{K}}(s,x;t,y), \quad (s,x),(t,y) \in [0,\infty) \times \mathbb{R}, \tag{1.106}
$$

<span id="page-27-3"></span>then we say that  $\widetilde{\Xi}$  is a **gauge transform** of  $\Xi$ .

Lemma 1.11 Under [\(1.106\)](#page-27-0),

$$
(\Xi(t))_{t\geq 0} \stackrel{\text{(law)}}{=} (\widetilde{\Xi}(t))_{t\geq 0}.
$$

In other words, DSP is gauge invariant.

By Proposition [1.7,](#page-22-1) [\(1.101\)](#page-26-2) has the DMR

$$
\Psi_{\xi}^{t}[f] = \mathbf{E}^{x} \left[ \prod_{m=1}^{N} \prod_{i=1}^{N} \{1 + \chi_{t_m}(B_i(t_m))\} \mathcal{D}_{\xi}(T, \mathbf{B}(T)) \right],
$$
\n(1.107)

where  $T \ge t_M$  and  $\xi = \sum_{i=1}^N \delta_{x_i}$ . The following equality is established [\[30\]](#page-88-0).

**Lemma 1.12** Let  $\boldsymbol{x} \in \mathbb{W}_{N}^{A}$  and  $\xi = \sum_{i=1}^{N} \delta_{x_i}$ . Then for any  $M \in \mathbb{N}$ ,  $0 \le t_1 < \cdots < t_M \le T < \infty$ ,  $\chi_{t_m} \in C_c(\mathbb{R}), 1 \leq m \leq M$ , the equality

$$
\mathbf{E}^{\mathbf{x}} \left[ \prod_{m=1}^{M} \prod_{i=1}^{N} \{1 + \chi_{t_m}(B_i(t_m))\} \mathcal{D}_{\xi}(T, \mathbf{B}(T)) \right] \n= \text{Det} \text{Det} \text{Set} \left[ \delta_{st} \delta_x(\{y\}) + \mathbb{K}_{\xi}(s, x; t, y) \chi_t(y) \right] \n(x, y) \in \mathbb{R}^2
$$
\n(1.108)

holds, where  $\mathbb{K}_{\xi}$  is given by [\(1.98\)](#page-25-3) with [\(1.93\)](#page-23-3).

Now we arrive at the following full characterization of  $\text{DYS}_2$ ,  $(\Xi, \mathbb{P}^{\xi})$ .

Theorem 1.13 For any finite and fixed initial configuration without multiple points, that is, for  $\xi \in \text{Conf}_0(\mathbb{R}), \xi(\mathbb{R}) = N \in \mathbb{N}, \text{DYS}_2$  is determinantal. Its correlation kernel is given by

<span id="page-27-1"></span>
$$
\mathbb{K}_{\xi}(s,x;t,y) = \mathcal{G}_{\xi}(s,x;t,y) - \mathbf{1}_{(s>t)}p(s-t,x|y), \quad (s,x),(t,y) \in [0,\infty) \times \mathbb{R} \tag{1.109}
$$

with

<span id="page-27-2"></span>
$$
\mathcal{G}_{\xi}(s,x;t,y) = \int_{\mathbb{R}} \xi(dv) p(s,x|v) \mathcal{M}_{\xi}^{v}(t,y). \tag{1.110}
$$

#### <span id="page-28-0"></span>1.7.2 Generalization for initial configuration with multiple points

For general  $\xi = \sum_{i=1}^{N} \delta_{x_i} \in \text{Conf}(\mathbb{R})$  with  $\xi(\mathbb{R}) = N < \infty$ , define supp $\xi = \{x \in \mathbb{R} : \xi(x) > 0\}$  and let  $\xi_*(\cdot) = \sum_{v \in \text{supp } \xi} \overline{\delta_v(\cdot)}$ . For  $s \in [0, \infty)$ ,  $v, x \in \mathbb{R}, z, \zeta \in \mathbb{C}$ , let

<span id="page-28-7"></span>
$$
\phi_{\xi}^{v}((s,x);z,\zeta) = \frac{p(s,x|\zeta)}{p(s,x|v)} \frac{1}{z-\zeta} \prod_{i=1}^{N} \frac{z-x_i}{\zeta-x_i},
$$
\n(1.111)

and

<span id="page-28-2"></span>
$$
\Phi_{\xi}^{v}((s,x);z) = \frac{1}{2\pi\sqrt{-1}} \oint_{C(\delta_{v})} d\zeta \, \phi_{\xi}^{v}((s,x);z,\zeta)
$$
  
\n
$$
= \text{Res}\left[\phi_{\xi}^{v}((s,x);z,\zeta); \zeta=v\right],
$$
\n(1.112)

where  $C(\delta_v)$  is a closed contour on the complex plane C encircling the point v once in the positive direction. This function [\(1.112\)](#page-28-2) is entire with respect to  $z \in \mathbb{C}$  parameterized by  $(s, x) \in [0, \infty) \times \mathbb{R}$ in addition to  $v \in \mathbb{C}, \xi \in \text{Conf}(\mathbb{R})$ . Remark that the polynomial function  $\Phi_{\xi}^{u}(z)$  defined by  $(1.85)$ is parameterized only by  $u \in \mathbb{C}$  and  $\xi \in \text{Conf}_0(\mathbb{R})$ . Here we start from this entire function and consider its  $I$ -transformation,

<span id="page-28-4"></span>
$$
\mathcal{M}_{\xi}^{v}((s,x)|(t,y)) = \mathcal{I}\left[\Phi_{\xi}^{v}((s,x);W)\Big|(t,y)\right], \quad (s,x),(t,y) \in [0,\infty) \times \mathbb{R},\tag{1.113}
$$

which provides a martingale, if we put  $y = B(t)$ ,  $t \ge 0$ . Then it is easy to see that [\(1.109\)](#page-27-1) with [\(1.110\)](#page-27-2) is rewritten as

<span id="page-28-3"></span>
$$
\mathbb{K}_{\xi}(s,x;t,y) = \int_{S} \xi_{*}(dv) p(s,x|v) \mathcal{M}_{\xi}^{v}((s,x)|(t,y)) - \mathbf{1}(s>t) p(s-t,x|y), \tag{1.114}
$$

 $(s, x), (t, y) \in [0, \infty) \times \mathbb{R}$ .

We note that the kernel [\(1.114\)](#page-28-3) with [\(1.113\)](#page-28-4) is bounded and integrable also for  $\xi \in \text{Conf}(\mathbb{R}) \setminus \mathbb{R}$ Conf<sub>0</sub>( $\mathbb{R}$ ). Therefore, the spatio-temporal correlations are given by [\(1.104\)](#page-26-3) for any  $0 \le t_1 < \cdots <$  $t_M < \infty, M \in \mathbb{N}$  and the finite-dimensional distributions are determined.

**Proposition 1.14** Also for  $\xi \in \text{Conf}(\mathbb{R}) \setminus \text{Conf}_0(\mathbb{R})$ , the DSPs with the correlation kernels [\(1.114\)](#page-28-3) are well-defined.

The complete proof of this proposition was given in Section 4.1 of [\[39\]](#page-89-3). See also [\[60,](#page-90-1) [61,](#page-90-2) [83,](#page-91-4) [62\]](#page-90-3). The above extension will provide the **entrance laws** for the processes  $(\Xi(t), t > 0, \mathbb{P}^{\xi})$  in the sense of Section XII.4 in [\[64\]](#page-90-4).

#### <span id="page-28-1"></span>1.7.3 DSP with extended Hermitian kernel

In order to give an example of Proposition [1.14,](#page-28-5) here we study the extreme case where all  $N$  points are concentrated on an origin,

<span id="page-28-6"></span><span id="page-28-5"></span>
$$
\xi = N\delta_0 \iff \xi_* = \delta_0 \quad \text{with} \quad \xi(\{0\}) = N. \tag{1.115}
$$

For [\(1.115\)](#page-28-6), [\(1.111\)](#page-28-7) and [\(1.112\)](#page-28-2) become

$$
\begin{array}{rcl}\n\phi_{N\delta_0}^0((s,x);z,\zeta) & = & \frac{p(s,x|\zeta)}{p(s,x|0)} \frac{1}{z-\zeta} \left(\frac{z}{\zeta}\right)^N \\
& = & \frac{p(s,x|\zeta)}{p(s,x|0)} \sum_{\ell=0}^\infty \frac{z^{N-\ell-1}}{\zeta^{N-\ell}},\n\end{array}
$$

and

<span id="page-29-0"></span>
$$
\Phi_{N\delta_0}^0((s,x);z) = \frac{1}{p(s,x|0)} \sum_{\ell=0}^{\infty} z^{N-\ell-1} \frac{1}{2\pi\sqrt{-1}} \oint_{C(\delta_0)} d\zeta \frac{p(s,x|\zeta)}{\zeta^{N-\ell}} \n= \frac{1}{p(s,x|0)} \sum_{\ell=0}^{N-1} z^{N-\ell-1} \frac{1}{2\pi\sqrt{-1}} \oint_{C(\delta_0)} d\zeta \frac{p(s,x|\zeta)}{\zeta^{N-\ell}},
$$
\n(1.116)

since the integrands are holomorphic when  $\ell \geq N$ .

For BM with the transition probability density [\(1.1\)](#page-3-0), [\(1.116\)](#page-29-0) gives

$$
\Phi_{N\delta_0}^0((s,x);z) = \sum_{\ell=0}^{N-1} z^{N-\ell-1} \frac{1}{2\pi\sqrt{-1}} \oint_{C(\delta_0)} d\zeta \frac{e^{x\zeta/s - \zeta^2/2s}}{\zeta^{N-\ell}}
$$
  
\n
$$
= \sum_{\ell=0}^{N-1} \left(\frac{z}{\sqrt{2s}}\right)^{N-\ell-1} \frac{1}{2\pi\sqrt{-1}} \oint_{C(\delta_0)} d\eta \frac{e^{2(x/\sqrt{2s})\eta - \eta^2}}{\eta^{N-\ell}}
$$
  
\n
$$
= \sum_{\ell=0}^{N-1} \left(\frac{z}{\sqrt{2s}}\right)^{N-\ell-1} \frac{1}{(N-\ell-1)!} H_{N-\ell-1} \left(\frac{x}{\sqrt{2s}}\right),
$$

where we have used the contour integral representation of the Hermite polynomials [\(1.34\)](#page-9-3). Thus its integral transformation is calculated as

$$
\mathcal{I}\left[\Phi_{N\delta_{0}}^{0}((s,x);W)\right](t,y)=\sum_{\ell=0}^{N-1}\frac{1}{(N-\ell-1)!}H_{N-\ell-1}\left(\frac{x}{\sqrt{2s}}\right)\frac{1}{(2s)^{(N-\ell-1)/2}}\mathcal{I}[W^{N-\ell-1}](t,y)=\sum_{\ell=0}^{N-1}\frac{1}{(N-\ell-1)!}H_{N-\ell-1}\left(\frac{x}{\sqrt{2s}}\right)\frac{1}{(2s)^{(N-\ell-1)/2}}m_{N-\ell-1}(t,y)=\sum_{\ell=0}^{N-1}\frac{1}{(N-\ell-1)!2^{N-\ell-1}}\left(\frac{t}{s}\right)^{(N-\ell-1)/2}H_{N-\ell-1}\left(\frac{x}{\sqrt{2s}}\right)H_{N-\ell-1}\left(\frac{y}{\sqrt{2t}}\right),
$$

where we have used Lemma [1.1](#page-9-4) and [\(1.40\)](#page-10-5) in Section [1.3.](#page-8-0) Then we obtain the following,

<span id="page-29-1"></span>
$$
\mathcal{M}_{N\delta_0}^0((s,x)|(t,B(t))) = \sum_{n=0}^{N-1} \frac{1}{n!2^n} m_n(s,x) m_n(t,B(t))
$$
  
=  $\sqrt{\pi}e^{x^2/4s + B(t)^2/4t} \sum_{n=0}^{N-1} \left(\frac{t}{s}\right)^{n/2} \varphi_n\left(\frac{x}{\sqrt{2s}}\right) \varphi_n\left(\frac{B(t)}{\sqrt{2t}}\right),$  (1.117)

where

$$
\varphi_n(x) = \frac{1}{\sqrt{\sqrt{\pi}2^n n!}} H_n(x) e^{-x^2/2}, \quad x \in \mathbb{R}, \quad n \in \mathbb{N}_0,
$$
\n(1.118)

are the Hermite orthonormal functions on R,

$$
\int_{\mathbb{R}} dx \,\varphi_n(x)\varphi_m(x) = \delta_{nm}, \quad n, m \in \mathbb{N}_0.
$$
\n(1.119)

The following expression for the transition probability density [\(1.1\)](#page-3-0) of BM is known as Mehler's formula,

<span id="page-30-0"></span>
$$
p(s-t,x|y) = \frac{e^{-x^2/4s}}{e^{-y^2/4t}} \frac{1}{\sqrt{2s}} \sum_{n=0}^{\infty} \left(\frac{t}{s}\right)^{n/2} \varphi_n\left(\frac{x}{\sqrt{2s}}\right) \varphi_n\left(\frac{y}{\sqrt{2t}}\right). \tag{1.120}
$$

Since  $m_n$ ,  $n \in \mathbb{N}_0$  are the fundamental martingale polynomials associated with BM, the process [\(1.117\)](#page-29-1) is a continuous martingale. Then we see that

$$
\mathbb{E}\left[\mathcal{M}_{N\delta_0}^0((s,x)|(t,\boldsymbol{B}(t)))\right] = \mathbb{E}\left[\mathcal{M}_{N\delta_0}^0((s,x)|(0,\boldsymbol{B}(0)))\right] = 1
$$

for  $(s, x) \in [0, \infty) \times \mathbb{R}, 0 \leq t < \infty$ .

By the formula [\(1.114\)](#page-28-3), we obtain the correlation kernels as

<span id="page-30-1"></span>
$$
\mathbb{K}_{N\delta_0}(s,x;t,y) = p(s,x|0)\mathcal{M}_{N\delta_0}^0((s,x)|(t,y)) - \mathbf{1}(s>t)p(s-t,x|y) \n= \frac{e^{-x^2/4s}}{e^{-y^2/4t}} \mathbf{K}_{\text{Hermite}}^{(N)}(s,x;t,y)
$$
\n(1.121)

with

$$
\mathbf{K}_{\text{Hermite}}^{(N)}(s, x; t, y) = \frac{1}{\sqrt{2s}} \sum_{n=0}^{N-1} \left(\frac{t}{s}\right)^{n/2} \varphi_n \left(\frac{x}{\sqrt{2s}}\right) \varphi_n \left(\frac{y}{\sqrt{2t}}\right)
$$

$$
- \mathbf{1}(s > t) \frac{1}{\sqrt{2s}} \sum_{n=0}^{\infty} \left(\frac{t}{s}\right)^{n/2} \varphi_n \left(\frac{x}{\sqrt{2s}}\right) \varphi_n \left(\frac{y}{\sqrt{2t}}\right).
$$

$$
= \begin{cases} \frac{1}{\sqrt{2s}} \sum_{n=0}^{N-1} \left(\frac{t}{s}\right)^{n/2} \varphi_n \left(\frac{x}{\sqrt{2s}}\right) \varphi_n \left(\frac{y}{\sqrt{2t}}\right) & \text{for } s \le t, \\ -\frac{1}{\sqrt{2s}} \sum_{n=N}^{\infty} \left(\frac{t}{s}\right)^{n/2} \varphi_n \left(\frac{x}{\sqrt{2s}}\right) \varphi_n \left(\frac{y}{\sqrt{2t}}\right) & \text{for } s > t, \end{cases} \tag{1.122}
$$

where Mehler's formula [\(1.120\)](#page-30-0) was used. By the gauge invariance (Lemma [1.11\)](#page-27-3), the factor  $e^{-x^2/4s}/e^{-y^2/4t}$  in [\(1.121\)](#page-30-1) is irrelevant for DSPs. The kernel  $\mathbf{K}_{\text{Hermite}}^{(N)}$  is known as the extended Hermite kernel (see, for instance, Exercise 11.6.3 in [\[21\]](#page-88-3)).

The equal-time correlation kernel

$$
K_{\text{Hermite}}^{(N,t)}(x,y) \equiv \mathbf{K}_{\text{Hermite}}^{(N)}(t,x;t,y)
$$

$$
= \frac{1}{\sqrt{2t}} \sum_{n=0}^{N-1} \varphi_n \left(\frac{x}{\sqrt{2t}}\right) \varphi_n \left(\frac{y}{\sqrt{2t}}\right)
$$

has the following expression,

$$
K_{\text{Hermite}}^{(N,t)}(x,y) = \sqrt{\frac{N}{2}} \frac{\varphi_N(x/\sqrt{2t})\varphi_{N-1}(y/\sqrt{2t}) - \varphi_{N-1}(x/\sqrt{2t})\varphi_N(y/\sqrt{2t})}{x-y},
$$
  
\nif  $x \neq y$ , (1.123)  
\n
$$
K_{\text{Hermite}}^{(N,t)}(x,x) = \frac{1}{\sqrt{2t}} \left[ N \left\{ \varphi_N\left(\frac{x}{\sqrt{2t}}\right) \right\}^2 - \sqrt{N(N+1)}\varphi_{N-1}\left(\frac{x}{\sqrt{2t}}\right) \varphi_{N+1}\left(\frac{x}{\sqrt{2t}}\right) \right].
$$

This spatial correlation kernel is a special case of the Christoffel–Dorboux kernel (see, for instance, Chapter 9 in [\[21\]](#page-88-3) and Chapter 3 in [\[2\]](#page-86-1)). It is called the Hermite kernel and defines the determinantal point process  $[81, 77]$  $[81, 77]$  on  $\mathbb R$  such that a spatial correlation function is given by

$$
\rho_{\text{Hermite}}^{(N,t)}(\boldsymbol{x}_{N'}) = \det_{1 \le i,j \le N'} \left[ \mathbf{K}_{\text{Hermite}}^{(N,t)}(x_i, x_j) \right] \tag{1.125}
$$

(1.124)

for any  $1 \leq N' \leq N$  and  $\boldsymbol{x}_{N'} = (x_1, \ldots, x_{N'}) \in \mathbb{R}^{N'}$ ,  $t > 0$ . We write the probability measure of this determinantal point process as  $P_{Hermite}^{(N,t)}$ .

#### <span id="page-31-1"></span><span id="page-31-0"></span>1.8 Exercises 1

#### 1.8.1 Exercise 1.1

For  $t > 0$ ,  $x \in \mathbb{R}$ , and  $\alpha \in \mathbb{C}$ , let

$$
G_{\alpha}(t,x) = e^{\alpha x - \alpha^2 t/2}.
$$
\n(1.126)

Let  $(B(t))_{t\geq0}$  be the 1 dimensional standard Brownian motion (BM) in  $(\Omega, \mathcal{P}, \mathcal{F})$  with the filtration  $\mathcal{F}_t = \sigma(B(s) : 0 \le s \le t), t \ge 0.$ 

(1) Prove the equality

$$
\frac{e^{\alpha B(t)}}{\mathcal{E}[e^{\alpha B(t)}]} = G_{\alpha}(t, B(t)), \quad t \ge 0.
$$
\n(1.127)

(2) Prove that  $(G_{\alpha}(t, B(t)))_{t>0}$  is  $\mathcal{F}_t$ -martingale;

$$
E[G_{\alpha}(t, B(t))|\mathcal{F}_s] = G_{\alpha}(s, B(s)), \quad 0 < s < t. \tag{1.128}
$$

(3) Let  $\widehat{G}_w(t, x)$  be the Fourier transformation of  $G_\alpha(t, x)$ ;

$$
\widehat{G}_w(t,x) := \int_{-\infty}^{\infty} \frac{e^{-\sqrt{-1}\alpha w}}{2\pi} G_{\alpha}(t,x) d\alpha.
$$
 (1.129)

Show that

$$
\widehat{G}_w(t,x) = \frac{e^{-(\sqrt{-1}x+w)^2/2t}}{\sqrt{2\pi t}}.
$$
\n(1.130)

(4) Define the integral transformation by

<span id="page-32-2"></span>
$$
\mathcal{I}[f(W)|(t,x)] = \int_{-\infty}^{\infty} f(\sqrt{-1}w)\widehat{G}_w(t,x)dw.
$$
 (1.131)

For  $n \in \mathbb{N}_0 := \{0, 1, \dots\}$ , let

$$
m_n(t, x) = \mathcal{I}[W^n|(t, x)]
$$
  
= 
$$
\int_{-\infty}^{\infty} (\sqrt{-1}w)^n \widehat{G}_w(t, x) dw, \quad t \ge 0.
$$
 (1.132)

Express  $m_n(t, x)$ ,  $n \in \mathbb{N}_0$  using the Hermite polynomials,

<span id="page-32-1"></span>
$$
H_n(x) := (-1)^n e^{x^2} \frac{d^n e^{-x^2}}{dx^n}
$$
  
= 
$$
\sum_{k=0}^{[n/2]} (-1)^k \frac{n!}{k!(n-2k)!} (2x)^{n-2k}.
$$
 (1.133)

- (5) Prove the following for  $\{m_n(t,x)\}_{n\in\mathbb{N}_0}$ .
	- (i) They are monic polynomials of degrees  $n \in \mathbb{N}_0$  with time-dependent coefficients:

<span id="page-32-3"></span>
$$
m_n(t,x) = x^n + \sum_{k=0}^{n-1} c_n^{(k)}(t) x^k, \quad t \ge 0.
$$

(ii) For  $0 \le k \le n-1$ ,  $c_n^{(k)}(0) = 0$ . That is,

$$
m_n(0, x) = x^n, \quad n \in \mathbb{N}_0.
$$

(iii) If we set  $x = B(t)$ , they provide martingales:

$$
E[m_n(t, B(t)) | \mathcal{F}_s] = m_n(s, B(s)), \quad 0 \le s \le t, \quad n \in \mathbb{N}_0.
$$
 (1.134)

# <span id="page-32-0"></span>1.8.2 Exercise 1.2

Let

$$
h_N(\boldsymbol{x}) := \prod_{1 \le i < j \le N} (x_j - x_i) = \det_{1 \le i, j \le N} [x_j^{i-1}]. \tag{1.135}
$$

(1) Using the basic properties of determinant, verify the equalities,

$$
\frac{h_N(\boldsymbol{y})}{h_N(\boldsymbol{x})} = \frac{1}{h_N(\boldsymbol{x})} \det_{1 \le i,j \le N} [y_j^{i-1}]
$$
\n
$$
= \frac{1}{h_N(\boldsymbol{x})} \det_{1 \le i,j \le N} [m_{i-1}(t, y_j)], \qquad (1.136)
$$

where  ${m_n(t, x)}_{n \in \mathbb{N}_0}$  are given by [\(1.132\)](#page-32-1) in Exercise 1.1.

(2) We extend the integral transformation  $\mathcal I$  defined by [\(1.131\)](#page-32-2) in Exercise 1.1 to a linear integral transformation of multivariate functions as follows: When  $F^{(i)}(\boldsymbol{x}) = \prod_{j=1}^{N} f_j^{(i)}$  $j^{(i)}(x_j), i=1,2$ are given for  $\boldsymbol{x} = (x_1, \ldots, x_N) \in \mathbb{R}^N$ , then we define

$$
\mathcal{I}\left[F^{(i)}(\mathbf{W})\left|\{(t_{\ell},x_{\ell})\}_{\ell=1}^{N}\right] := \prod_{j=1}^{N} \mathcal{I}\left[f_{j}^{(i)}(W_{j})\right](t_{j},x_{j})\right], \quad i=1,2,
$$
\n(1.137)

and

$$
\mathcal{I}\Big[c_{1}F^{(1)}(\mathbf{W})+c_{2}F^{(2)}(\mathbf{W})\left|\{(t_{\ell},x_{\ell})\}_{\ell=1}^{N}\right] \n:=c_{1}\mathcal{I}\Big[F^{(1)}(\mathbf{W})\left|\{(t_{\ell},x_{\ell})\}_{\ell=1}^{N}\right]+c_{2}\mathcal{I}\Big[F^{(2)}(\mathbf{W})\left|\{(t_{\ell},x_{\ell})\}_{\ell=1}^{N}\right],
$$
\n(1.138)

 $c_1, c_2 \in \mathbb{C}$ , for  $0 < t_i < \infty$ ,  $i = 1, \ldots, N$ , where  $\mathbf{W} = (W_1, \ldots, W_N) \in \mathbb{R}^N$ . In particular, if  $t_{\ell} = t, 1 \leq \forall \ell \leq N$ , we write  $\mathcal{I}[\cdot | \{(t_{\ell}, x_{\ell})\}_{\ell=1}^{N}]$  simply as  $\mathcal{I}[\cdot | (t, x)]$  with  $\boldsymbol{x} = (x_1, \ldots, x_N)$ . Using multilinearity of determinants, verify the following equalities,

$$
\frac{h_N(\boldsymbol{y})}{h_N(\boldsymbol{x})} = \frac{1}{h_N(\boldsymbol{x})} \det_{1 \le i,j \le N} \left[ \mathcal{I}[(W_j)^{i-1} | (t, y_j)] \right]
$$
\n
$$
= \mathcal{I} \left[ \frac{1}{h_N(\boldsymbol{x})} \det_{1 \le i,j \le N} [(W_j)^{i-1}] \middle| (t, \boldsymbol{y}) \right]
$$
\n
$$
= \mathcal{I} \left[ \frac{h_N(\mathbf{W})}{h_N(\boldsymbol{x})} \middle| (t, \boldsymbol{y}) \right].
$$
\n(1.139)

(3) Prove the equality,

<span id="page-33-1"></span>
$$
\frac{h_N(z)}{h_N(x)} = \det_{1 \le i,j \le N} \left[ \Phi_{\xi}^{x_i}(z_j) \right],\tag{1.140}
$$

where

<span id="page-33-0"></span>
$$
\Phi_{\xi}^{u}(z) = \prod_{\substack{1 \le k \le N, \\ x_k \ne u}} \frac{z - x_k}{u - x_k},\tag{1.141}
$$

 $\xi = \sum_{i=1}^{N} \delta_{x_i} \in \text{Conf}_0(\mathbb{R}), z, u \in \mathbb{C}$ , following the steps given below.

(i) Let

$$
H(\boldsymbol{x}, \boldsymbol{z}) = \det_{1 \leq i, j \leq N} \left[ \prod_{1 \leq k \leq N, k \neq i} (z_j - x_k) \right]. \tag{1.142}
$$

Show that (a) it is a polynomial function of  $x_i, z_i, 1 \leq i \leq N$  with degree  $N(N-1)$ , (b)  $H(x,x) = h(x)^2$ , and (c)  $H(x,z) = 0$ , if  $x_i = x_j$  or  $z_i = z_j$  for some i, j with  $1 \leq i < j \leq N$ . It is implied by (a)–(c) that  $H(\boldsymbol{x}, \boldsymbol{z}) = h_N(\boldsymbol{x})h_N(\boldsymbol{z})$ .

(ii) By the definition [\(1.141\)](#page-33-0), hence the RHS of [\(1.140\)](#page-33-1) is equal to  $H(x, z)/h(x)^2$  for  $\xi \in \text{Conf}_0(\mathbb{R}).$ 

#### <span id="page-34-0"></span>1.8.3 Exercise 1.3

For a pair of space-time coordinates  $(s, x), (t, y) \in [0, \infty) \times \mathbb{R}$ , consider the function,

$$
\mathcal{M}_N((s,x)|(t,y)) = \sum_{n=0}^{N-1} \frac{1}{n!2^n} m_n(s,x) m_n(t,y), \qquad (1.143)
$$

where  ${m_n(t, x)}_{n \in \mathbb{N}_0}$  are given by [\(1.132\)](#page-32-1) in Exercise 1.1.

## (1) The Hermite orthonormal functions are defined by

$$
\varphi_n(x) = \frac{1}{\sqrt{\sqrt{\pi} 2^n n!}} H_n(x) e^{-x^2/2}, \quad x \in \mathbb{R}, \quad n \in \mathbb{N}_0,
$$
\n(1.144)

where  ${H_n(x)}_{n \in \mathbb{N}_0}$  are the Hermite polynomials defined by [\(1.133\)](#page-32-3). Prove the following expression of  $\mathcal{M}_N((s,x)|(t,y))$ ,

$$
\mathcal{M}_N((s,x)|(t,y)) = \sqrt{\pi}e^{x^2/4s+y^2/4t} \sum_{n=0}^{N-1} \left(\frac{t}{s}\right)^{n/2} \varphi_n\left(\frac{x}{\sqrt{2s}}\right) \varphi_n\left(\frac{y}{\sqrt{2t}}\right). \tag{1.145}
$$

(2) Let  $(B(t))_{t\geq0}$  be BM adapted to the filtration  $\mathcal{F}_t, t\geq0$ . Then prove that, for fixed  $(s, x) \in$  $[0,\infty) \times \mathbb{R}, (\mathcal{M}_N((s,x)|(t,B(t))))_{t\geq 0}$  is an  $\mathcal{F}_t$ -martingale, and

 $E[\mathcal{M}_N((s,x)|(t, \mathcal{B}(t)))] = E[\mathcal{M}_N((s,x)|(0, \mathcal{B}(0))))] = 1.$  (1.146)

(3) The following identity is known as Mehler's formula,

$$
\sum_{n=0}^{\infty} \frac{H_n(x)H_n(y)}{2^n n!} a^n = \frac{e^{[2xya - (x^2+y^2)a^2]/(1-a^2)}}{\sqrt{1-a^2}}, \quad |a| < 1. \tag{1.147}
$$

Prove the following limit,

$$
\lim_{N \to \infty} \mathcal{M}_N((s, x)|(t, y)) = \frac{p(s - t, x|y)}{p(s, x|0)}, \quad 0 < t < s, \quad x, y \in \mathbb{R}, \tag{1.148}
$$

where  $p(t, y|x)$  denotes the transition probability density of BM,

$$
p(t, y|x) = \frac{e^{-(y-x)^2/2t}}{\sqrt{2\pi t}}, \quad t > 0, \quad x, y \in \mathbb{R}.
$$
 (1.149)

(4) Define the spatio-temporal kernel as follows,

$$
K(s, x; t, y) := p(s, x|0) \mathcal{M}_N((s, x)|(t, y)) - \mathbf{1}(s > t) p(s - t, x|y), \tag{1.150}
$$

 $(s, x), (t, y) \in [0, \infty) \times \mathbb{R}$ , where  $\mathbf{1}(\omega)$  is an indicator;  $\mathbf{1}(\omega) = 1$ , if  $\omega$  is satisfied, and  $\mathbf{1}(\omega) = 0$ , otherwise. Then prove the equality

$$
\mathbb{K}_{N\delta_0}(s,x;t,y) = \frac{e^{-x^2/4s}}{e^{-y^2/4t}} \mathbf{K}_{\text{Hermite}}^{(N)}(s,x;t,y),\tag{1.151}
$$

where

$$
\mathbf{K}_{\text{Hermite}}^{(N)}(s, x; t, y)
$$
\n
$$
= \begin{cases}\n\frac{1}{\sqrt{2s}} \sum_{n=0}^{N-1} \left(\frac{t}{s}\right)^{n/2} \varphi_n\left(\frac{x}{\sqrt{2s}}\right) \varphi_n\left(\frac{y}{\sqrt{2t}}\right) & \text{for } s \le t, \\
-\frac{1}{\sqrt{2s}} \sum_{n=N}^{\infty} \left(\frac{t}{s}\right)^{n/2} \varphi_n\left(\frac{x}{\sqrt{2s}}\right) \varphi_n\left(\frac{y}{\sqrt{2t}}\right) & \text{for } s > t.\n\end{cases} \tag{1.152}
$$

The spatio-temporal kernel  $\mathbf{K}_{\text{Hermite}}^{(N)}(s, x; t, y)$  is known as the extended Hermite kernel giving the correlation kernel for  $DYS_2$  starting from the configuration such that all N particles are at the origin,  $N\delta_0 \in \text{Conf}(\mathbb{R}) \setminus \text{Conf}_0(\mathbb{R})$ .

# <span id="page-35-1"></span><span id="page-35-0"></span>2 Gaussian Analytic Functions and Their Zero-Point Processes

## <span id="page-35-2"></span>2.1 Reproducing kernel Hilbert spaces

## 2.1.1 Bergman spaces and Hardy spaces

A Hilbert function space of holomorphic functions is a Hilbert space  $H$  of functions on a domain D in C equipped with the inner product  $\langle \cdot, \cdot \rangle_{\mathcal{H}},$ 

$$
\langle af + bg, h \rangle_{\mathcal{H}} = a \langle f, h \rangle_{\mathcal{H}} + b \langle g, h \rangle_{\mathcal{H}},
$$
  
\n
$$
\langle h, af + bg \rangle_{\mathcal{H}} = \overline{a} \langle h, f \rangle_{\mathcal{H}} + \overline{b} \langle h, g \rangle_{\mathcal{H}}, \qquad f, g, h \in \mathcal{H}, \quad a, b \in \mathbb{C}
$$
  
\n
$$
\langle g, f \rangle_{\mathcal{H}} = \overline{\langle f, g \rangle_{\mathcal{H}}} = \langle \overline{f}, \overline{g} \rangle_{\mathcal{H}},
$$

and the norm  $||f||_{\mathcal{H}} := \sqrt{\langle f, f \rangle_{\mathcal{H}}}$ . The evaluation at each point of D is a continuous functional on H. Therefore, for each point  $w \in D$ , there is an element of H, which is called the **reproducing kernel at**  $w$  and denote by  $k_w$ , with the property

$$
\langle f, k_w \rangle_{\mathcal{H}} = f(w), \quad \forall f \in \mathcal{H}.
$$

Because  $k_w \in \mathcal{H}$ , it is itself a function on D,

$$
k_w(z) = \langle k_w, k_z \rangle_{\mathcal{H}}.
$$

We write

$$
k_{\mathcal{H}}(z,w) := k_{w}(z) = \langle k_{w}, k_{z} \rangle_{\mathcal{H}}
$$
\n(2.1)

and call it the reproducing kernel for  $H$ . By definition, it is Hermitian;

$$
\overline{k_{\mathcal{H}}(z,w)} = k_{\mathcal{H}}(w,z), \quad z, w \in D.
$$

If H is a holomorphic Hilbert function space, then  $k<sub>H</sub>$  is holomorphic in the first variable and anti-holomorphic in the second. We see that  $k<sub>H</sub>(z, w)$  is a **positive semi-definite kernel**: for any  $n \in \mathbb{N} := \{1, 2, \dots\}$ , for any points  $z_i \in D$  and  $\xi_i \in \mathbb{C}, i = 1, 2, \dots, n$ ,

$$
\sum_{i=1}^{n} \sum_{j=1}^{n} k_{\mathcal{H}}(z_i, z_j) \xi_i \overline{\xi_j} = \sum_{i=1}^{n} \sum_{j=1}^{n} \langle k_{z_j}, k_{z_i} \rangle_{\mathcal{H}} \xi_i \overline{\xi_j} = \sum_{i=1}^{n} \sum_{j=1}^{n} \langle k_{\mathcal{H}}(\cdot, z_j), k_{\mathcal{H}}(\cdot, z_i) \rangle_{\mathcal{H}} \xi_i \overline{\xi_j}
$$

$$
= \left\langle \sum_{i=1}^{n} k_{\mathcal{H}}(z_i, \cdot) \xi_i, \sum_{j=1}^{n} k_{\mathcal{H}}(z_j, \cdot) \xi_j \right\rangle_{\mathcal{H}} = \left\| \sum_{i=1}^{n} \xi_i k_{\mathcal{H}}(z_i, \cdot) \right\|_{\mathcal{H}}^2 \ge 0. \tag{2.2}
$$
Let  $\{e_n : n \in \mathcal{I}\}\)$  be any **complete orthonormal systme** (CONS) for  $\mathcal{H}$ , where  $\mathcal{I}$  is an index set,

$$
\langle e_n, e_m \rangle_{\mathcal{H}} = \delta_{nm}, n, m \in \mathcal{I},
$$
  

$$
f \in \mathcal{H} \iff f = \sum_{n \in \mathcal{I}} c_n e_n \text{ with } (c_n)_{n \in \mathcal{I}} \in \ell^2(\mathcal{I}).
$$

Then one can prove that the reproducing kernel for  $\mathcal H$  is written in the form

<span id="page-36-0"></span>
$$
k_{\mathcal{H}}(z,w) = \sum_{n \in \mathcal{I}} e_n(z) \overline{e_n(w)}.
$$
\n(2.3)

Actually, this gives

$$
\langle f(\cdot), k_{\mathcal{H}}(\cdot, w) \rangle_{\mathcal{H}} = \left\langle \sum_{m \in \mathcal{I}} c_m e_m(\cdot), \sum_{n \in \mathcal{I}} e_n(\cdot) \overline{e_n(w)} \right\rangle_{\mathcal{H}} = \sum_{m \in \mathcal{I}} \sum_{n \in \mathcal{I}} c_m \langle e_m, e_n \rangle_{\mathcal{H}} e_n(w)
$$

$$
= \sum_{n \in \mathcal{I}} c_n e_n(w) = f(w), \ \forall f \in \mathcal{H}, w \in D.
$$

We note that the positive definiteness of the kernel  $(2.2)$  is equivalent with the situation such that, for any points  $z_i \in D$ ,  $i \in \mathbb{N}$ , the matrix  $(k_{\mathcal{H}}(z_i, z_k))_{1 \leq i,j \leq n}$  has a nonnegative determinant,

$$
\det_{1\leq i,j\leq n}[k_{\mathcal{H}}(z_i,z_j)] \geq 0 \text{ for any } n \in \mathbb{N}.
$$

First we show two examples of holomorphic Hilbert function spaces, the Bergman space and the **Hardy space**, for a unit disk  $\mathbb{D} := \{z \in \mathbb{C} : |z| < 1\}$  and the domains which are conformally transformed from D.

**Example 2.1** The **Bergman space** on  $\mathbb{D}$ , denoted by  $L^2(\mathbb{D})$ , is the Hilbert space of holomorphic functions on  $\mathbb D$  which are square-integrable with respect to the Lebesgue measure on  $\mathbb C$ . The inner product for  $L^2(\mathbb{D})$  is given by

$$
\langle f, g \rangle_{L^2_B(\mathbb{D})} := \frac{1}{\pi} \int_{\mathbb{D}} f(z) \overline{g(z)} m(dz) = \sum_{n=0}^{\infty} \frac{\widehat{f}(n) \overline{\widehat{g}(n)}}{n+1},
$$

where the *n*th Taylor coefficient of f at 0 is denoted by  $\widehat{f}(n)$ ;  $f(z) = \sum_{n=0}^{\infty} \widehat{f}(n)z^n$ . Let

<span id="page-36-1"></span>
$$
\widetilde{e_n}(z) := \sqrt{n+1}z^n, n \in \mathbb{N}_0.
$$

Then  $\{\tilde{e}_n(z)\}_{n\in\mathbb{N}_0}$  form a CONS for  $L^2(\mathbb{D})$  and the reproducing kernel [\(2.3\)](#page-36-0) is given by

$$
K_{\mathbb{D}}(z, w) := k_{L_{\mathbb{B}}^2(\mathbb{D})}(z, w)
$$
  
= 
$$
\sum_{n \in \mathbb{N}_0} (n+1)(z\overline{w})^n = \frac{1}{(1 - z\overline{w})^2}, \quad z, w \in \mathbb{D}.
$$
 (2.4)

This kernel is called the **Bergman kernel** of  $D$ .

**Example 2.2** The Hardy space on  $\mathbb{D}$ ,  $H^2(\mathbb{D})$ , consists of holomorphic functions on  $\mathbb{D}$  such that the Taylor coefficients form a square-summable series;

$$
||f||_{H^2(\mathbb{D})}^2 := \sum_{n \in \mathbb{N}_0} |\widehat{f}(n)|^2 < \infty, \quad f \in H^2(\mathbb{D}).
$$

For every  $f \in H^2(\mathbb{D})$ , the non-tangential limit  $\lim_{r \uparrow 1} f(re^{\sqrt{-1}\phi})$  exists a.e. by Fatou's theorem and we write it as  $f(e^{\sqrt{-1}\phi})$ . It is known that  $f(e^{\sqrt{-1}\phi}) \in L^2(\partial \mathbb{D})$ . Then one can prove that the inner product of  $H^2(\mathbb{D})$  is given by the following three different ways,

<span id="page-37-0"></span>
$$
\langle f, g \rangle_{H^2(\mathbb{D})} = \begin{cases} \sum_{n \in \mathbb{N}_0} \widehat{f}(n) \overline{\widehat{g}(n)}, \\ \lim_{r \uparrow 1} \frac{1}{2\pi} \int_0^{2\pi} f(re^{\sqrt{-1}\phi}) \overline{g(re^{\sqrt{-1}\phi})} d\phi, & f, g \in H^2(\mathbb{D}), \\ \frac{1}{2\pi} \int_0^{2\pi} f(e^{\sqrt{-1}\phi}) \overline{g(e^{\sqrt{-1}\phi})} d\phi, \end{cases}
$$
(2.5)

with  $||f||_{H^2(\mathbb{D})}^2 = \langle f, f \rangle_{H^2(\mathbb{D})}$ . Let  $\sigma$  be the measure on the boundary of D which is the usual arc length measure. Then the last expression of the inner product [\(2.5\)](#page-37-0) is written as  $\langle f, g \rangle_{H^2(\mathbb{D})}$  =  $(1/2\pi)\int_{\gamma_1} f(z)\overline{g(z)}\sigma(dz)$ , where  $\gamma_1$  is a unit circle  $\{e^{\sqrt{-1}\phi} : \phi \in [0, 2\pi)\}\$ giving the boundary of D. If we set  $e_n(z) := e_n^{(0,0)}(z) = z^n, n \in \mathbb{N}_0$ , then  $\{e_n(z)\}_{n \in \mathbb{N}_0}$  form CONS for  $H^2(\mathbb{D})$ . The reproducing kernel  $(2.3)$  is given by

<span id="page-37-1"></span>
$$
S_{\mathbb{D}}(z, w) := k_{H^2(\mathbb{D})}(z, w)
$$
  
= 
$$
\sum_{n \in \mathbb{N}_0} (z\overline{w})^n = \frac{1}{1 - z\overline{w}}, \quad z, w \in \mathbb{D},
$$
 (2.6)

which is called the **Szegő** kernel of  $D$ .

Let  $f : D \to \tilde{D}$  be a **conformal transformation** (the angle-preserving one-to-one map) between two bounded domains  $D, \widetilde{D} \subset \mathbb{C}$  with  $\mathcal{C}^{\infty}$  smooth boundary (analytic boundary). We find an argument in Chapter 12 of  $[8]$  concluding that the derivative of the transformation f denoted by f' has a single valued square root on D. We let  $\sqrt{f'(z)}$  denote one of the square roots of f'. The Szegő kernel and the Bergman kernel are then transformed by  $f$  as [\[8,](#page-87-0) Chapters 12, 16]

<span id="page-37-2"></span>
$$
S_D(z, w) = \sqrt{f'(z)} \sqrt{f'(w)} S_{\tilde{D}}(f(z), f(w)),
$$
  
\n
$$
K_D(z, w) = |f'(z)||f'(w)| K_{\tilde{D}}(f(z), f(w)), \quad z, w \in D.
$$
\n(2.7)

Consider the special case in which  $D \subseteq \mathbb{C}$  is a simply connected domain with analytic boundary and  $\tilde{D} = \mathbb{D}$ . For each  $\alpha \in D$ , **Riemann's mapping theorem** gives a unique conformal transformation;

<span id="page-37-3"></span>
$$
h_{\alpha}: D \to \mathbb{D} \quad \text{conformal such that} \quad h_{\alpha}(\alpha) = 0, \ h'_{\alpha}(\alpha) > 0. \tag{2.8}
$$

Such  $h_{\alpha}$  is called the **Riemann mapping function**. By [\(2.6\)](#page-37-1), the first equation in [\(2.7\)](#page-37-2) gives

$$
S_D(z, w) = \sqrt{h'_{\alpha}(z)} \overline{\sqrt{h'_{\alpha}(w)}} S_{\mathbb{D}}(h_{\alpha}(z), h_{\alpha}(w))
$$
  
=  $\sqrt{h'_{\alpha}(z)} \overline{\sqrt{h'_{\alpha}(w)}} \frac{1}{1 - h_{\alpha}(z) \overline{h_{\alpha}(w)}}, \quad z, w \in D.$ 

Since  $h_{\alpha}(\alpha) = 0$ , if we put  $z = \alpha$ , or  $w = \alpha$ , or  $z = w = \alpha$  in this formula, then we readily obtain

$$
S_D(\alpha, w) = \sqrt{h'_\alpha(\alpha)} \overline{\sqrt{h'_\alpha(w)}} \frac{1}{1 - 0 \times \overline{h_\alpha(w)}} = \sqrt{h'_\alpha(\alpha)} \overline{\sqrt{h'_\alpha(w)}},
$$
  
\n
$$
S_D(z, \alpha) = \sqrt{h'_\alpha(z)} \overline{\sqrt{h'_\alpha(\alpha)}}, \quad S_D(\alpha, \alpha) = \sqrt{h'_\alpha(\alpha)} \overline{\sqrt{h'_\alpha(\alpha)}}
$$
  
\n
$$
\implies \frac{S_D(z, \alpha) S_D(\alpha, w)}{S_D(\alpha, \alpha)} = \frac{\sqrt{h'_\alpha(\alpha)} \overline{\sqrt{h'_\alpha(w)}} \overline{\sqrt{h'_\alpha(z)} \overline{\sqrt{h'_\alpha(\alpha)}}}}{\sqrt{h'_\alpha(\alpha)} \overline{\sqrt{h'_\alpha(\alpha)}}} = \sqrt{h'_\alpha(z)} \overline{\sqrt{h'_\alpha(w)}}
$$

Hence the above formula is written as [\[7\]](#page-87-1),

<span id="page-38-0"></span>
$$
S_D(z, w) = \frac{S_D(z, \alpha) \overline{S_D(w, \alpha)}}{S_D(\alpha, \alpha)} \frac{1}{1 - h_{\alpha}(z) \overline{h_{\alpha}(w)}}, \quad z, w, \alpha \in D.
$$
 (2.9)

Similarly, we have

$$
K_D(z, w) = \frac{S_D(z, \alpha)^2 \overline{S_D(w, \alpha)^2}}{S_D(\alpha, \alpha)^2} \frac{1}{(1 - h_\alpha(z) \overline{h_\alpha(w)})^2}, \quad z, w, \alpha \in D.
$$
 (2.10)

Hence the following relationship is established,

<span id="page-38-5"></span>
$$
S_D(z, w)^2 = K_D(z, w), \quad z, w \in D.
$$
\n(2.11)

Here we notice that  $(2.9)$  is rewritten as

<span id="page-38-4"></span>
$$
S_D(z, w) - \frac{S_D(z, \alpha)S_D(\alpha, w)}{S_D(\alpha, \alpha)} = S_D(z, w)h_\alpha(z)\overline{h_\alpha(w)}, \quad z, w, \alpha \in D. \tag{2.12}
$$

LHS is regarded as the reproducing kernel for the Hilbert subspace

$$
H_a^2(D) := \{ f \in H^2(D) : f(\alpha) = 0 \},\tag{2.13}
$$

and we will denote such a **conditional kernel** with a deterministic zero at  $\alpha$  by  $S_{D}^{\alpha}(z, w)$ . Hence we obtain the formula

<span id="page-38-2"></span>
$$
S_D^{\alpha}(z, w) = S_D(z, w) h_{\alpha}(z) \overline{h_{\alpha}(w)}, \quad z, w, \alpha \in D.
$$
 (2.14)

Let  $q \in (0,1)$  be a fixed number and we consider the **annulus** 

$$
\mathbb{A}_q:=\{z\in\mathbb{C}:q<|z|<1\}.
$$

<span id="page-38-3"></span>As the third example of the Hilbert spaces of holomorphic functions, we consider the Hardy space on  $\mathbb{A}_q$ .

**Example 2.3** The Hardy space for  $\mathbb{A}_q$ ,  $H^2(\mathbb{A}_q)$ , is the Hilbert space of holomorphic functions on  $\mathbb{A}_q$  equipped with the inner product

<span id="page-38-1"></span>
$$
\langle f, g \rangle_{H^2(\mathbb{A}_q)} = \frac{1}{2\pi} \int_0^{2\pi} f(e^{\sqrt{-1}\phi}) \overline{g(e^{\sqrt{-1}\phi})} d\phi + \frac{1}{2\pi} \int_0^{2\pi} f(qe^{\sqrt{-1}\phi}) \overline{g(qe^{\sqrt{-1}\phi})} q d\phi, \ f, g \in H^2(\mathbb{A}_q). \tag{2.15}
$$

A CONS of  $H^2(\mathbb{A}_q)$  is given by  $\{e_n^{(q,q)}\}_{n\in\mathbb{Z}}$  with

$$
e_n^{(q,q)}(z) = \frac{z^n}{\sqrt{1+q^{2n+1}}}, \quad z \in \mathbb{A}_q, \quad n \in \mathbb{Z},
$$

and the reproducing kernel is given by [\[52\]](#page-89-0)

$$
S_{\mathbb{A}_q}(z,w) = \sum_{n \in \mathbb{Z}} e_n^{(q,q)}(z) \overline{e_n^{(q,q)}(w)} = \sum_{n=-\infty}^{\infty} \frac{(z\overline{w})^n}{1 + q^{2n+1}}.
$$
 (2.16)

This infinite series converges absolutely for  $z, w \in A_q$  and is called the **Szegő kernel of**  $A_q$ .

The nonexistence of zero in  $\mathbb{D}$  of  $S_{\mathbb{D}}(\cdot,\alpha), \alpha \in \mathbb{D}$  and the uniqueness of zero in  $\mathbb{A}_q$  of  $S_{\mathbb{A}_q}(\cdot,\alpha)$ ,  $\alpha \in \mathbb{A}_q$  are concluded from a general consideration (see, for instance, [\[8,](#page-87-0) Section 27]). Define

<span id="page-39-1"></span>
$$
\widehat{\alpha} := -\frac{q}{\overline{\alpha}}, \quad \alpha \in \mathbb{A}_q. \tag{2.17}
$$

The fact

<span id="page-39-0"></span>
$$
S_{\mathbb{A}_q}(\widehat{\alpha}, \alpha) = S_{\mathbb{A}_q}(\alpha, \widehat{\alpha}) = 0, \quad \alpha \in \mathbb{A}_q
$$
\n(2.18)

was proved as Theorem 1 in [\[82\]](#page-91-0) by direct calculation, for which a simpler proof can be given using theta functions (see Lemma [2.5](#page-43-0) below).

Again we consider the conformal transformation  $f: D \to \widetilde{D}$  between two bounded domains  $D, \widetilde{D} \subsetneq \mathbb{C}$  with analytic boundaries. Here we consider the case in which D is a 2-connected domain and  $D = \mathbb{D}$ . For each  $\alpha \in D$ , there exists a unique function  $f_a$  giving a branched 2 to 1 covering map of D to D, in which a unique point  $\hat{\alpha} = \hat{\alpha}(\alpha) \neq \alpha$ ,  $\hat{\alpha} \in D$  exists such that

$$
f_{\alpha}(\alpha) = f_{\alpha}(\widehat{\alpha}) = 0.
$$

The map  $f_{\alpha}, \alpha \in D$  is called the **Ahlfors mapping function** [\[8,](#page-87-0) Chapter 13]. By the fact [\(2.18\)](#page-39-0) with [\(2.17\)](#page-39-1), for an arbitrary, but fixed  $\alpha \in D$ , we can verify that the following set forms a CONS for  $H^2(\mathbb{A}_q)$  which is different from  $\{e_n^{(q,q)}(z)\}_{n\in\mathbb{Z}}$  given by  $(2.15)$ ,

$$
\widehat{e}_{jn}(z) := \frac{S_{\mathbb{A}_q}(z, \alpha_j)}{\sqrt{S_{\mathbb{A}_q}(\alpha_j, \alpha_j)}} f_\alpha(z)^n, \quad j = 0, 1, \quad n \in \mathbb{N}_0,
$$
\n(2.19)

where

<span id="page-39-2"></span>
$$
\alpha_0 := \alpha, \quad \alpha_1 := \widehat{\alpha}.
$$

Hence the Szegő kernel of  $\mathbb{A}_q$  can be expressed as

$$
S_{\mathbb{A}_q}(z,w) = \sum_{j=0}^{1} \sum_{n=0}^{\infty} \widehat{e}_{jn}(z) \overline{\widehat{e}_{jn}(w)}
$$
  
= 
$$
\left(\frac{S_{\mathbb{A}_q}(z,\alpha) \overline{S_{\mathbb{A}_q}(w,\alpha)}}{S_{\mathbb{A}_q}(\alpha,\alpha)} + \frac{S_{\mathbb{A}_q}(z,\widehat{\alpha}) \overline{S_{\mathbb{A}_q}(w,\widehat{\alpha})}}{S_{\mathbb{A}_q}(\widehat{\alpha},\widehat{\alpha})}\right) \sum_{n=0}^{\infty} f_{\alpha}(z)^n \overline{f_{\alpha}(z)^n}
$$
  
= 
$$
\left(\frac{S_{\mathbb{A}_q}(z,\alpha) \overline{S_{\mathbb{A}_q}(w,\alpha)}}{S_{\mathbb{A}_q}(\alpha,\alpha)} + \frac{S_{\mathbb{A}_q}(z,\widehat{\alpha}) \overline{S_{\mathbb{A}_q}(w,\widehat{\alpha})}}{S_{\mathbb{A}_q}(\widehat{\alpha},\widehat{\alpha})}\right) \frac{1}{1 - f_{\alpha}(z) \overline{f_{\alpha}(w)}}, z, w \in \mathbb{A}_q.
$$
(2.20)

Then, if we consider the Hilbert subspace of  $\mathbb{A}_q$  such that

$$
H_{\alpha,\widehat{\alpha}}^2(\mathbb{A}_q) := \{ f \in H^2(\mathbb{A}_q) : f(\alpha) = f(\widehat{\alpha}) = 0 \},
$$

then its reproducing kernel should be given by

$$
S_{\mathbb{A}_q}^{\alpha,\widehat{\alpha}}(z,w) := S_{\mathbb{A}_q}(z,w) - \left( \frac{S_{\mathbb{A}_q}(z,\alpha)S_{\mathbb{A}_q}(\alpha,w)}{S_{\mathbb{A}_q}(\alpha,\alpha)} + \frac{S_{\mathbb{A}_q}(z,\widehat{\alpha})S_{\mathbb{A}_q}(\widehat{\alpha},w)}{S_{\mathbb{A}_q}(\widehat{\alpha},\widehat{\alpha})} \right).
$$
(2.21)

It is easy to verify that [\(2.20\)](#page-39-2) gives [\[8\]](#page-87-0)

<span id="page-40-0"></span>
$$
S_{\mathbb{A}_q}^{\alpha,\widehat{\alpha}}(z,w) = S_{\mathbb{A}_q}(z,w) f_{\alpha}(z) \overline{f_{\alpha}(w)}, \quad z, w \in \mathbb{A}_q.
$$
 (2.22)

Now we ask the following question.

(Q) For  $\alpha \in A_{q}$ , consider the conditional Hilbert space

$$
H_{\alpha}^{2}(\mathbb{A}_{q}) := \{ f \in H^{2}(\mathbb{A}_{q}) : f(\alpha) = 0 \},
$$

whose reproducing kernel is given by

<span id="page-40-6"></span>
$$
S_{\mathbb{A}_q}^{\alpha}(z,w) := S_{\mathbb{A}_q}(z,w) - \frac{S_{\mathbb{A}_q}(z,\alpha)S_{\mathbb{A}_q}(\alpha,w)}{S_{\mathbb{A}_q}(\alpha,\alpha)}, \quad z, w \in \mathbb{A}_q.
$$
 (2.23)

Is it possible to factorize this **conditional kernel** in the form as  $(2.14)$  and  $(2.22)$ ?

## <span id="page-40-5"></span>2.1.2 Theta function

Assume that  $p \in \mathbb{C}$  is a fixed number such that  $0 < |p| < 1$ . We use the following standard notation,

$$
(a; p)_n := \prod_{i=0}^{n-1} (1 - ap^i), \qquad (a; p)_\infty := \prod_{i=0}^\infty (1 - ap^i),
$$
  

$$
(a_1, \dots, a_k; p)_\infty := (a_1; p)_\infty \cdots (a_k; p)_\infty.
$$
 (2.24)

The **theta function** with argument  $z$  and **nome**  $p$  is defined by

<span id="page-40-4"></span><span id="page-40-3"></span>
$$
\theta(z;p) := (z, p/z; p)_{\infty}.
$$
\n(2.25)

We often use the shorthand notation  $\theta(z_1,\ldots,z_n;p) := \prod_{i=1}^n \theta(z_i;p)$ .

As a function of z, the theta function  $\theta(z; p)$  is holomorphic in  $\mathbb{C}^\times$  and has single zeros precisely at  $p^i$ ,  $i \in \mathbb{Z}$ , that is,

$$
\{z \in \mathbb{C}^{\times} : \theta(z; p) = 0\} = \{p^i : i \in \mathbb{Z}\}.
$$
\n(2.26)

We will use the inversion formula

<span id="page-40-1"></span>
$$
\theta(1/z;p) = -\frac{1}{z}\theta(z;p)
$$
\n(2.27)

and the quasi-periodicity property

<span id="page-40-2"></span>
$$
\theta(pz;p) = -\frac{1}{z}\theta(z;p)
$$
\n(2.28)

of the theta function. By comparing [\(2.27\)](#page-40-1) and [\(2.28\)](#page-40-2) and performing the transformation  $z \mapsto 1/z$ , we immediately see the periodicity property,

$$
\theta(p/z; p) = \theta(z; p). \tag{2.29}
$$

By Jacobi's triple product identity (see, for instance, [\[22,](#page-88-0) Section 1.6]), we have the Laurent expansion

$$
\theta(z;p) = \frac{1}{(p;p)_{\infty}} \sum_{n \in \mathbb{Z}} (-1)^n p^{\binom{n}{2}} z^n.
$$

One can show that [\[59,](#page-90-0) Chapter 20]

$$
\lim_{p \to 0} \theta(z; p) = 1 - z,\tag{2.30}
$$

$$
\theta'(1;p) := \frac{\partial \theta(z;p)}{\partial z}\Big|_{z=1} = -(p;p)^2_{\infty}.
$$
\n(2.31)

The theta function satisfies the following Weierstrass' addition formula [\[41\]](#page-89-1),

<span id="page-41-1"></span>
$$
\theta(xy, x/y, uv, u/v; p) - \theta(xv, x/v, uy, u/y; p) = \frac{u}{y}\theta(yv, y/v, xu, x/u; p).
$$
 (2.32)

When p is real and  $p \in (0,1)$ , we see that

<span id="page-41-2"></span><span id="page-41-0"></span>
$$
\overline{\theta(z;p)} = \theta(\overline{z};p). \tag{2.33}
$$

In this case the definition [\(2.25\)](#page-40-3) with [\(2.24\)](#page-40-4) implies that

$$
\begin{aligned}\n\theta(x;p) > 0, & x \in (p^{2i+1}, p^{2i}) \\
\theta(x;p) &= 0, & x = p^i \\
\theta(x;p) < 0, & x \in (p^{2i}, p^{2i-1}) \\
\theta(x;p) > 0, & x \in (-\infty, 0).\n\end{aligned}\n\quad i \in \mathbb{Z},
$$
\n(2.34)

Moreover, we can prove the following: In the interval  $x \in (-\infty,0)$ ,  $\theta(x) := \theta(x;p)$  is strictly convex with

$$
\min_{x \in (-\infty, 0)} \theta(x) = \theta(-\sqrt{p}) = \prod_{n=1}^{\infty} (1 + p^{n-1/2})^2 > 0,
$$
\n(2.35)

and  $\lim_{x\downarrow-\infty}\theta(x) = \lim_{x\uparrow0}\theta(x) = +\infty$ , and in the interval  $x \in (p,1)$ ,  $\theta(x)$  is strictly concave with

<span id="page-41-3"></span>
$$
\max_{x \in (p,1)} \theta(x) = \theta(\sqrt{p}) = \prod_{n=1}^{\infty} (1 - p^{n-1/2})^2,
$$
\n(2.36)

 $\theta(x) \sim (p; p)_{\infty}^2 (x-p)/p$  as  $x \downarrow p$ , and  $\theta(x) \sim (p; p)_{\infty}^2 (1-x)$  as  $x \uparrow 1$ , where [\(2.28\)](#page-40-2) and [\(2.31\)](#page-41-0) were used.

#### <span id="page-42-2"></span>2.1.3 Weighted Szegő kernels of annulus and conformal transformations

We introduce a positive parameter  $r > 0$ . Consider the Hilbert space of holomorphic functions on  $A_q$  equipped with the inner product

$$
\langle f, g \rangle_{H_r^2(\mathbb{A}_q)} = \frac{1}{2\pi} \int_{\gamma_1 \cup \gamma_q} f(z) \overline{g(z)} \sigma_r(dz), \quad f, g \in H_r^2(\mathbb{A}_q)
$$

with

$$
\sigma_r(dz) = \begin{cases} d\phi, & \text{if } z \in \gamma_1 := \{e^{\sqrt{-1}\phi} : \phi \in [0, 2\pi)\}, \\ r d\phi, & \text{if } z \in \gamma_q := \{qe^{\sqrt{-1}\phi} : \phi \in [0, 2\pi)\}, \end{cases}
$$

which we write as  $H_r^2(\mathbb{A}_q)$ . A CONS of  $H_r^2(\mathbb{A}_q)$  is given by  $\{e_n^{(q,r)}\}_{n\in\mathbb{Z}}$  with

$$
e_n^{(q,r)}(z) = \frac{z^n}{\sqrt{1 + r q^{2n}}}, \quad z \in \mathbb{A}_q, \quad n \in \mathbb{Z},
$$

and the reproducing kernel is given by [\[52\]](#page-89-0)

<span id="page-42-0"></span>
$$
S_{\mathbb{A}_q}(z, w; r) = \sum_{n \in \mathbb{Z}} e_n^{(q, r)}(z) \overline{e_n^{(q, r)}(w)} = \sum_{n = -\infty}^{\infty} \frac{(z\overline{w})^n}{1 + r q^{2n}}.
$$
 (2.37)

This infinite series converges absolutely for  $z, w \in A_q$ .

When  $r = q$ , this Hilbert function space is reduced to the Hardy space for  $\mathbb{A}_q$ ,  $H^2(\mathbb{A}_q)$ , introduced in Example [2.3](#page-38-3) in Section [2.1.1,](#page-35-1) associated with the reproducing kernel

$$
S_{\mathbb{A}_q}(\cdot,\cdot):=S_{\mathbb{A}_q}(\cdot,\cdot;q).
$$

The kernel [\(2.37\)](#page-42-0) with a parameter  $r > 0$  is called the **weighted Szegő kernel** of  $\mathbb{A}_q$  and  $H_r^2(\mathbb{A}_q)$ is the **reproducing kernel Hilbert space** (RKHS) with respect to  $S_{\mathbb{A}_q}(\cdot,\cdot;r)$  [\[52\]](#page-89-0). We call  $r$  the weight parameter [\[36\]](#page-88-1). Notice that  $(2.37)$  implies that  $S_{A_q}(z, z; r)$  is a monotonically decreasing function of the weight parameter  $r \in (0,\infty)$  for each fixed  $z \in \mathbb{A}_q$ .

From now on, we put

<span id="page-42-3"></span>
$$
p = q^2 \tag{2.38}
$$

for the nome p of the theta function [\(2.25\)](#page-40-3), and assume that  $\theta(z) = \theta(z; q)$ , where  $q \in (0, 1)$  specifies the radius of the inner circle of the annulus  $\mathbb{A}_q$ .

## Proposition 2.4 (McCullough and Shen [\[52\]](#page-89-0)) For  $r > 0$

<span id="page-42-1"></span>
$$
S_{\mathbb{A}_q}(z, w; r) = \frac{q_0^2 \theta(-rz\overline{w})}{\theta(-r, z\overline{w})}, \quad z, w \in \mathbb{A}_q,
$$
\n(2.39)

where

$$
q_0 := \prod_{n \in \mathbb{N}} (1 - q^{2n}) = (q^2; q^2)_{\infty}.
$$
\n(2.40)

In particular,

$$
S_{\mathbb{A}_q}(z,w) = S_{\mathbb{A}_q}(z,w;q) = \frac{q_0^2 \theta(-qz\overline{w})}{\theta(-q,z\overline{w})}, \quad z, w \in \mathbb{A}_q.
$$
 (2.41)

Proof. The following function has been studied in [\[52,](#page-89-0) [14,](#page-87-2) [84\]](#page-91-1),

<span id="page-43-1"></span>
$$
f^{JK}(z, a; q) := \sum_{n \in \mathbb{Z}} \frac{z^n}{1 - aq^{2n}},
$$
\n(2.42)

with  $q^2 < |z| < 1$ ,  $a \notin \{q^{2i} : i \in \mathbb{Z}\}$ , which is called the **Jordan–Kronecker function** [\[84,](#page-91-1) p.59]. We readily find that

<span id="page-43-2"></span>
$$
S_{\mathbb{A}_q}(z, w; r) = f^{\mathbb{J}\mathbb{K}}(z\overline{w}, -r; q), \quad z, w \in \mathbb{A}_q.
$$
 (2.43)

The **bilateral basic hypergeometric series** in base  $p$  with one numerator parameter  $a$  and one denominator parameter  $b$  is defined by  $[22]$ 

$$
{}_1\psi_1(a;b;p,z) = {}_1\psi_1\left[\begin{array}{c} a \\ b \end{array};p,z\right] := \sum_{n \in \mathbb{Z}} \frac{(a;p)_n}{(b;p)_n} z^n, \quad |b/a| < |z| < 1.
$$

The Jordan–Kronecker function [\(2.42\)](#page-43-1) is a special case of the  $_1\psi_1$  function [\[14,](#page-87-2) [84\]](#page-91-1);

$$
f^{JK}(z,a;q) = \frac{1}{1-a} \psi_1(a;aq^2;q^2,z).
$$

The following equality is known as **Ramanujan's**  $_1\psi_1$  summation formula [\[14,](#page-87-2) [22,](#page-88-0) [84\]](#page-91-1),

$$
\sum_{n\in\mathbb{Z}}\frac{(a;p)_n}{(b;p)_n}z^n=\frac{(az,p/(az),p,b/a;p)_\infty}{(z,b/(az),b,p/a;p)_\infty},\quad |b/a|<|z|<1.
$$

Combining the above two equalities with an appropriate change of variables, we obtain [\[14,](#page-87-2) [84\]](#page-91-1)

<span id="page-43-3"></span>
$$
f^{\text{JK}}(z,a;q) = \frac{(az,q^2/(az),q^2,q^2;q^2)_{\infty}}{(z,q^2/z,a,q^2/a;q^2)_{\infty}} = \frac{q_0^2 \theta(za;q^2)}{\theta(z,a;q^2)}.
$$
\n(2.44)

By  $(2.43)$ ,  $(2.44)$  proves the equality  $(2.39)$ . The proof is complete.

Here we notice that the following symmetries of  $f^{JK}$  are readily verified by [\(2.44\)](#page-43-3) using [\(2.27\)](#page-40-1) and [\(2.28\)](#page-40-2) [\[14,](#page-87-2) [84\]](#page-91-1).

$$
f^{JK}(z,a) = f^{JK}(a,z),
$$
\n(2.45)

<span id="page-43-5"></span><span id="page-43-4"></span>
$$
f^{\text{JK}}(z,a) = -f^{\text{JK}}(z^{-1}, a^{-1}),\tag{2.46}
$$

<span id="page-43-0"></span>
$$
f^{\text{JK}}(z,a) = z f^{\text{JK}}(z,aq^2) = af^{\text{JK}}(zq^2,a). \tag{2.47}
$$

Since  $\theta$  is holomorphic in  $\mathbb{C}^{\times}$ , the expression [\(2.39\)](#page-42-1) implies that  $S_{\mathbb{A}_q}(z,\alpha;r)$  is meromorphic in  $\mathbb{C}^{\times}$ . By the basic properties of  $\theta$  given in Section [2.1.2,](#page-40-5) the following is easily verified.

**Lemma 2.5** Assume that  $\alpha \in \mathbb{A}_q$ . Then  $S_{\mathbb{A}_q}(z,\alpha;r)$  has zeros at  $z = -q^{2i}/(\overline{\alpha}r)$ ,  $i \in \mathbb{Z}$  in  $\mathbb{C}^{\times}$ . In particular,  $S_{\mathbb{A}_q}(z,\alpha)$  has a unique zero in  $\mathbb{A}_q$  at  $z=\widehat{\alpha}$  given by [\(2.17\)](#page-39-1).

We have found the answer to the question  $(Q)$  addressed in Section [2.1.1](#page-35-1) in the paper by McCullough and Shen [\[52\]](#page-89-0).

Proposition 2.6 (Mccullough and Shen [\[52\]](#page-89-0)) The equality

<span id="page-44-1"></span>
$$
S_{\mathbb{A}_q}^{\alpha}(z, w; r) = S_{\mathbb{A}_q}(z, w; r|\alpha|^2) h_{\alpha}^q(z) \overline{h_{\alpha}^q(w)}, \quad z, w, \alpha \in \mathbb{A}_q,
$$
\n(2.48)

holds with

<span id="page-44-0"></span>
$$
h_{\alpha}^{q}(z) := z \frac{\theta(\alpha/z)}{\theta(\overline{\alpha}z)} = -\alpha \frac{\theta(z/\alpha)}{\theta(z\overline{\alpha})}, \quad z, \alpha \in \mathbb{A}_{q}.
$$
 (2.49)

The special case with  $r = q$  answers the question (Q) as

$$
S_{\mathbb{A}_q}^{\alpha}(z,w) = S_{\mathbb{A}_q}(z,w;q|\alpha|^2)h_{\alpha}^q(z)\overline{h_{\alpha}^q(w)}, \quad z, w, \alpha \in \mathbb{A}_q. \tag{2.50}
$$

*Proof.* We put  $(2.23)$  with  $(2.39)$  and  $(2.49)$  to  $(2.48)$ , then the equality is expressed by theta functions. After multiplying both sides by the common denominator, we see that the equality [\(2.48\)](#page-44-1) is equivalent to the following,

<span id="page-44-2"></span>
$$
\theta(-rz\overline{w}, -r|\alpha|^2, \overline{\alpha}z, \alpha \overline{w}) - \theta(-r\overline{\alpha}z, -r\alpha \overline{w}, z\overline{w}, |\alpha|^2) \n= z\overline{w}\theta(-rz\overline{w}|\alpha|^2, \alpha z^{-1}, \overline{\alpha}\,\overline{w}^{-1}, -r).
$$
\n(2.51)

Now we change the variables from  $\{z, \overline{w}, \alpha, r\}$  to  $\{x, y, u, v\}$  as  $\overline{\alpha}z = x/y$ ,  $\alpha \overline{w} = u/v$ ,  $z\overline{w} = x/v$ ,  $|\alpha|^2 = u/y$ , and  $r = -yv$ . Then LHS of [\(2.51\)](#page-44-2) becomes  $\theta(xy, x/y, uv, u/v) - \theta(xv, x/v, uy, u/y)$ , and RHS side becomes  $(x/v)\theta(yv,(y/v)^{-1}, xu,(x/u)^{-1})$  which is equal to  $(u/y)\theta(yv,y/v,xu,x/u)$ by [\(2.27\)](#page-40-1). Hence Weierstrass' addition formula [\(2.32\)](#page-41-1) proves the equality [\(2.51\)](#page-44-2). The proof is complete.

Following the formula [\(2.12\)](#page-38-4) for  $S^{\alpha}_{\mathbb{A}_q}$ , conditional kernels  $S^{\alpha_1,...,\alpha_n}_{\mathbb{A}_q}$  $\mathbb{A}_q^{\alpha_1,\dots,\alpha_n}$  are inductively defined as

$$
S_{\mathbb{A}_q}^{\alpha_1, ..., \alpha_n}(z, w) = (S_{\mathbb{A}_q}^{\alpha_1, ..., \alpha_{n-1}})^{\alpha_n}(z, w), \quad z, w, \alpha_1, ..., \alpha_n \in D, \quad n = 2, 3, .... \tag{2.52}
$$

The kernels  $S^{\alpha_1,...,\alpha_n}_{\mathbb{A}_n}$  $\mathbb{A}_q^{a_1,...,a_n}$ ,  $n = 2, 3,...$ , will construct Hilbert subspaces of the Hardy space for  $H^2(\mathbb{A}_q)$ ,

$$
H^{2}_{\alpha_1,...,\alpha_n}(\mathbb{A}_q) := \{ f \in \mathcal{H}_k : f(\alpha_1) = \cdots = f(\alpha_n) = 0 \}.
$$

For  $n \in \mathbb{N}, \alpha_1, \ldots, \alpha_n \in \mathbb{A}_q$ , define

<span id="page-44-6"></span>
$$
\gamma_{\{\alpha_{\ell}\}_{\ell=1}^n}^q(z) := \prod_{\ell=1}^n h_{\alpha_{\ell}}^q(z), \quad z \in \mathbb{A}_q. \tag{2.53}
$$

Then the Mccullough and Shen formula [\(2.48\)](#page-44-1) [\[52\]](#page-89-0) is generalized as

<span id="page-44-3"></span>
$$
S_{\mathbb{A}_q}^{\alpha_1,\dots,\alpha_n}(z,w;r) = S_{\mathbb{A}_q}(z,w;r) \prod_{\ell=1}^n |\alpha_\ell|^2 \gamma_{\{\alpha_\ell\}_{\ell=1}^n}^q(z) \overline{\gamma_{\{\alpha_\ell\}_{\ell=1}^n}^q(w)}, \quad z, w \in \mathbb{A}_q, \tag{2.54}
$$

for  $n \in \mathbb{N}, \alpha_1, \ldots, \alpha_n \in \mathbb{A}_q$ . As a special case of [\(2.54\)](#page-44-3), we have

<span id="page-44-4"></span>
$$
S_{\mathbb{A}_q}^{\alpha,\widehat{\alpha}}(z,w) = S_{\mathbb{A}_q}(z,w) f_{\alpha}^q(z) \overline{f_{\alpha}^q(w)}, \quad z, w \in \mathbb{A}_q
$$
\n(2.55)

with

<span id="page-44-5"></span>
$$
f_{\alpha}^{q}(z) := \frac{1}{z} h_{\alpha}^{q}(z) h_{\hat{\alpha}}^{q}(z)
$$
  
= 
$$
z \frac{\theta(-qz\overline{\alpha}, \alpha/z)}{\theta(-qz/\alpha, \overline{\alpha}z)} = -\alpha \frac{\theta(-qz\overline{\alpha}, z/\alpha)}{\theta(-qz/\alpha, z\overline{\alpha})}.
$$
 (2.56)

The equation [\(2.55\)](#page-44-4) is nothing but [\(2.22\)](#page-40-0), and  $f_{\alpha}^{q}$  given by [\(2.56\)](#page-44-5) is identified with the **Ahlfors map** from  $\mathbb{A}_q$  to  $\mathbb{D}$ .

We can prove the following.

**Lemma 2.7** For  $\alpha \in \mathbb{A}_q$ ,

<span id="page-45-2"></span><span id="page-45-1"></span>
$$
(i) \quad h^q_\alpha(\alpha) = 0,\tag{2.57}
$$

(ii) 
$$
0 < |h^q_\alpha(z)| < 1 \quad \forall z \in \mathbb{A}_q \setminus \{\alpha\},\tag{2.58}
$$

(iii) 
$$
|h_{\alpha}^{q}(z)| = \begin{cases} 1, & \text{if } z \in \gamma_1 := \{z \in \mathbb{C} : |z| = 1\}, \\ |\alpha|, & \text{if } z \in \gamma_q := \{z \in \mathbb{C} : |z| = q\}, \end{cases}
$$
 (2.59)

(iv) 
$$
h_{\alpha}^{q'}(\alpha) = -\frac{\theta'(1)}{\theta(|\alpha|^2)} = \frac{q_0^2}{\theta(|\alpha|^2)} > 0,
$$
 (2.60)

<span id="page-45-3"></span><span id="page-45-0"></span>(v) 
$$
\lim_{q \to 0} h_{\alpha}^{q}(z) = \frac{z - \alpha}{1 - \overline{\alpha}z} = z \frac{1 - \alpha/z}{1 - \overline{\alpha}z} =: h_{\alpha}(z).
$$
 (2.61)

Hence,  $h^q_\alpha$  is identified with a **conformal map** from  $\mathbb{A}_q$  to the unit disk with a **circular slit** in it, in which  $\alpha \in A_q$  is sent to the origin [\[52\]](#page-89-0). Notice that  $h_\alpha$  is the **Möbius transformation**  $\mathbb{D} \to \mathbb{D}$ sending  $\alpha$  to the origin.



Figure 3: Conformal map  $h^q_\alpha : \mathbb{A}_q \to \mathbb{D} \setminus \{ \text{a circular slit} \}$  is illustrated for  $q = 1/3$  and  $\alpha = 2/3$ . The point  $\alpha = 2/3$  in  $\mathbb{A}_{1/3}$  is mapped to the origin. The outer boundary  $\gamma_1$  of  $\mathbb{A}_{1/3}$  (denoted by a red circle) is mapped to a unit circle (a red circle) making the boundary of D. The inner boundary  $\gamma_{1/3}$  of  $\mathbb{A}_{1/3}$  (a green circle) is mapped to a circular slit (denoted by a green arc) which is a part of the circle with radius  $\alpha = 2/3$ , where the map is two-to-one except the two points on  $\gamma_{1/3}$  mapped to the two edges of the circular slit.

## 2.2 I.i.d.Gaussian power series, Laurent series, and their zero point processes

For a domain  $D \subset \mathbb{C}$ , let X be a random variable on a probability space which takes values in the space of analytic functions on D. If  $(X(z_1),...,X(z_n))$  follows a centered (mean zero) complex Gaussian distribution for every  $n \in \mathbb{N}$  and every  $z_1, \ldots, z_n \in D$ , X is said to be a **Gaussian** analytic function (GAF) [\[26\]](#page-88-2). The zero set of X is regarded as a point process on D denoted by a nonnegative-integer-valued Radon measure

$$
\mathcal{Z}_X = \sum_{z \in D: X(z) = 0} \delta_z,
$$

and it is simply called the **zero point process** of the GAF. Zero-point processes of GAFs have been extensively studied in quantum and statistical physics as solvable models of quantum chaotic systems and interacting particle systems. Many important characterizations of their probability laws have been reported in probability theory.

A typical example of GAF is provided by the i.i.d. Gaussian power series defined on the unit disk D: Let  $\mathbb{N}_0 := \{0, 1, 2, \dots\}$  and  $\{\zeta_n\}_{n \in \mathbb{N}_0}$  be i.i.d. standard complex Gaussian random variables with density

$$
p(z) = \frac{1}{\pi}e^{-|z|^2} = \frac{1}{\sqrt{\pi}}e^{-x^2} \cdot \frac{1}{\sqrt{\pi}}e^{-y^2}, \quad z = x + \sqrt{-1}y, \quad x, y \in \mathbb{R},
$$

and consider a random power series,

$$
X_{\mathbb{D}}(z) = \sum_{n=0}^{\infty} \zeta_n z^n,
$$
\n(2.62)

which defines an analytic function on  $\mathbb D$  a.s. This gives a GAF on  $\mathbb D$  with a **covariance kernel** 

<span id="page-46-1"></span>
$$
\mathbf{E}[X_{\mathbb{D}}(z)\overline{X_{\mathbb{D}}(w)}] = \frac{1}{1 - z\overline{w}} =: S_{\mathbb{D}}(z, w), \quad z, w \in \mathbb{D}.
$$
 (2.63)

This kernel is identified with the **Szegő kernel** of  $D(2.6)$  $D(2.6)$  of the Hardy space  $H^2(D)$  introduced in Section [2.1.1\)](#page-35-1). Peres and Virág [\[63\]](#page-90-1) proved that  $\mathcal{Z}_{X_{\mathbb{D}}}$  is a **determinantal point process** (DPP) such that the **correlation kernel** is given by  $S_{\mathbb{D}}(z,w)^2 = (1 - z\overline{w})^{-2}$ ,  $z, w \in \mathbb{D}$  with respect to the reference measure  $\lambda = m/\pi$ . Here m represents the Lebesgue measure on C;  $m(dz) := dxdy$ ,  $z = x + \sqrt{-1}y \in \mathbb{C}$ . This correlation kernel is identified with the reproducing kernel of the Bergman space on  $\mathbb{D}$  [\(2.4\)](#page-36-1), which is called **Bergman kernel** of  $\mathbb{D}$  in Section [2.1.1.](#page-35-1) Thus the study of Peres and Virág on  $X_{\mathbb{D}}$  and  $\mathcal{Z}_{X_{\mathbb{D}}}$  is associated with the following relationship between kernels on  $\mathbb{D}$  [\[63\]](#page-90-1),

<span id="page-46-2"></span>
$$
\mathbf{E}[X_{\mathbb{D}}(z)\overline{X_{\mathbb{D}}(w)}]^2 = S_{\mathbb{D}}(z,w)^2 = K_{\mathbb{D}}(z,w), \quad z, w \in \mathbb{D}.
$$
 (2.64)

Associated with the RKHS  $H_r^2(\mathbb{A}_q)$  studied in Section [2.1.3](#page-42-2) we consider the **Gaussian Laurent** series

<span id="page-46-0"></span>
$$
X_{\mathbb{A}_q}^r(z) := \sum_{n \in \mathbb{Z}} \zeta_n e_n^{(q,r)}(z) = \sum_{n = -\infty}^{\infty} \zeta_n \frac{z^n}{\sqrt{1 + r q^{2n}}},\tag{2.65}
$$

where  $\{\zeta_n\}_{n\in\mathbb{Z}}$  are i.i.d. standard complex Gaussian random variables with density  $e^{-|z|^2}/\pi$ . Since  $\lim_{n\to\infty} |\zeta_n|^{1/n} = 1$  a.s., we apply the Cauchy–Hadamard criterion to the positive and negative powers of  $X_{\mathbb{A}_q}^r(z)$  separately to conclude that this random Laurent series converges a.s. whenever  $z \in A_q$ . Moreover, since the distribution  $\zeta_n$  is symmetric, both of  $\gamma_1$  and  $\gamma_q$  are natural boundaries. Hence  $X_{\mathbb{A}_q}^r$  provides a GAF on  $\mathbb{A}_q$  whose covariance kernel is given by the **weighted Szegő kernel** of  $\mathbb{A}_q$ ,

$$
\mathbf{E}[X_{\mathbb{A}_q}^r(z)\overline{X_{\mathbb{A}_q}^r(w)}] = S_{\mathbb{A}_q}(z, w; r), \quad z, w \in \mathbb{A}_q,
$$

and the zero point process is denoted by  $\mathcal{Z}_{X_{\mathbb{A}_q}^r} := \sum_{z \in \mathbb{A}_q : X_{\mathbb{A}_q}^r} (z) = 0 \delta_z$ . In particular, we write  $X_{\mathbb{A}_q}(z):=X^q_{\mathbb{A}}$  $\mathcal{A}_{\mathcal{A}_q}(z), z \in \mathbb{A}_q$  and  $\mathcal{Z}_{X_{\mathbb{A}_q}} := \mathcal{Z}_{X_{\mathbb{A}_q}^q}$  as mentioned above.

We recall **Schottky's theorem**: The group of conformal transformations from  $A_q$  to itself is generated by the rotations and the  $q$ -inversions

<span id="page-47-2"></span>
$$
T_q(z) := \frac{q}{z}.
$$

The invariance of the present GAF and its zero point process under rotation is obvious. Using the properties of  $S_{A_q}$ , we can prove the following.

# **Proposition 2.8** (i) The GAF  $X_{\mathbb{A}_q}^r$  given by [\(2.65\)](#page-46-0) has the  $(q, r)$ -inversion symmetry in the sense that

$$
\left\{ (T_q'(z))^{1/2} X_{\mathbb{A}_q}^r(T_q(z)) \right\} \stackrel{\text{d}}{=} \left\{ \sqrt{\frac{q}{r}} X_{\mathbb{A}_q}^{q^2/r}(z) \right\}, \quad z \in \mathbb{A}_q,
$$

where  $T'_q(z) := \frac{dT_q}{dz}(z) = -q/z^2$ .

- (ii) For  $\mathcal{Z}_{X_{\mathbb{A}_q}} = \sum_i \delta_{Z_i}$ , let  $T_q^* \mathcal{Z}_{X_{\mathbb{A}_q}^r} := \sum_i \delta_{T_q^{-1}(Z_i)}$ . Then  $T_q^* \mathcal{Z}_{X_{\mathbb{A}_q}^r} \stackrel{\rm d}{=} \mathcal{Z}_{X_{\mathbb{A}_q}^{q^2/r}}$
- (iii) In particular, when  $r = q$ , the GAF  $X_{A_q}$  is invariant under conformal transformations which preserve  $\mathbb{A}_q$ , and so is its zero point process  $\mathcal{Z}_{X_{\mathbb{A}_q}}$ .

Proof. The equivalence in distribution asserted by (i) is implied by proving the following equality of the covariance kernel given by the weighted Szegő kernel of  $\mathbb{A}_q$ ,

<span id="page-47-0"></span>
$$
\sqrt{T_q'(z)}\sqrt{T_q'(w)}S_{\mathbb{A}_q}(T_q(z), T_q(w); r) = \frac{q}{r}S_{\mathbb{A}_q}(z, w; q^2/r).
$$
\n(2.66)

<span id="page-47-1"></span>.

If we use the expression [\(2.43\)](#page-43-2) of  $S_{A_q}(z, w; r)$ , the symmetries of the Jordan–Kronecker function [\(2.46\)](#page-43-4) and [\(2.47\)](#page-43-5) prove [\(2.66\)](#page-47-0). By the definition of zero point processes, (ii) is concluded from (i). When  $r = q$ , [\(2.66\)](#page-47-0) becomes

$$
\sqrt{T'_q(z)}\sqrt{T'_q(w)}S_{\mathbb{A}_q}(T_q(z),T_q(w);q)=\sqrt{T'_q(z)}\sqrt{T'_q(w)}S_{\mathbb{A}_q}(T_q(z),T_q(w))=S_{\mathbb{A}_q}(z,w),
$$

which implies the invariance of the GAF  $X_{A_q}$  under conformal transformations preserving  $A_q$  by Schottkey's theorem and (iii) is proved.  $\blacksquare$ 

We can give probabilistic interpretations of the above facts as follows.

**Proposition 2.9** For any  $\alpha_1, \ldots, \alpha_n \in \mathbb{A}_q$ ,  $n \in \mathbb{N}$ , the following hold.

(i) The following equality is established,

$$
\{X_{\mathbb{A}_q}^r(z) : z \in \mathbb{A}_q\} \text{ given } \{X_{\mathbb{A}_q}^r(\alpha_1) = \dots = X_{\mathbb{A}_q}^r(\alpha_n) = 0\}
$$
  

$$
\stackrel{\text{d}}{=} \left\{\gamma_{\{\alpha_\ell\}_{\ell=1}^n}^q(z) X_{\mathbb{A}_q}^r \prod_{\ell=1}^n |\alpha_\ell|^2(z) : z \in \mathbb{A}_q\right\}.
$$

<span id="page-47-3"></span>(ii) Let  $\mathcal{Z}_{X_{A_q}}^{\alpha_1,\dots,\alpha_n}$  denote the zero point process of the GAF  $X_{A_q}^r(z)$  given  $\{X_{A_q}^r(\alpha_1) = \cdots =$  $X_{\mathbb{A}_q}^r(\alpha_n) = 0$ . Then,

$$
\mathcal{Z}_{X_{\mathbb{A}_q}^r}^{\alpha_1,\dots,\alpha_n} \stackrel{\text{d}}{=} \mathcal{Z}_{X_{\mathbb{A}_q}^r \prod_{\ell=1}^n |\alpha_\ell|^2} + \sum_{i=1}^n \delta_{\alpha_i}.
$$

**Remark 2.10** For the GAF on  $\mathbb D$  studied by Peres and Virág [\[63\]](#page-90-1),  $\{X_{\mathbb D}(z) : z \in \mathbb D\}$  given  $\{X_{\mathbb D}(\alpha) =$ 0} is equal in law to  $\{h_{\alpha}(z)X_{\mathbb{D}}(z): z \in \mathbb{D}\}\,$ ,  $\forall \alpha \in \mathbb{D}$ , where  $h_{\alpha}$  is given by [\(2.61\)](#page-45-0), and then, in the notation used in Proposition [2.9,](#page-47-1)

$$
\mathcal{Z}_{X_{\mathbb{D}}}^{\alpha} \stackrel{\mathrm{d}}{=} \mathcal{Z}_{X_{\mathbb{D}}} + \delta_{\alpha}, \quad \forall \alpha \in \mathbb{D}.
$$

Hence, no new GAF nor new zero point process appear by conditioning of zeros. For the present GAF on  $\mathbb{A}_q$ , however, conditioning of zeros induces new GAFs and new zero point processes as shown by Proposition [2.9.](#page-47-1) Actually, by [\(2.37\)](#page-42-0) the covariance of the induced GAF  $X_{\mathbb{A}}^r \prod_{\ell=1}^n |\alpha_\ell|^2$  $\mathbb{A}_q^{I}$   $\mathbb{I}_{\ell=1}^{|\alpha_{\ell}|}$  is expressed by

$$
S_{\mathbb{A}_q}\left(z,w;r\prod_{\ell=1}^n|\alpha_\ell|^2\right)=\sum_{n=-\infty}^\infty\frac{(z\overline{w})^n}{1+r\prod_{\ell=1}^n|\alpha_\ell|^2q^{2n}}.
$$

Since  $q < |\alpha_{\ell}| < 1$ , as increasing the number of conditioning zeros, the variance of induced GAF monotonically increases, in which the increment is a decreasing function of  $|\alpha_{\ell}| \in (q, 1)$ .

## 2.3 Permanental-determinantal point processes

#### 2.3.1 Correlation functions of point processes

A point process is formulated as follows [\[35\]](#page-88-3). Let S be a base space, which is locally compact Hausdorff space with a countable base, and a **reference measure**  $\lambda$  is given on S. The configuration space of a point process on  $S$  is given by

$$
Conf(S) = \Big\{\xi = \sum_{i} \delta_{x_i} : x_i \in S, \, \xi(\Lambda) < \infty \text{ for all bounded set } \Lambda \subset S\Big\}.
$$

We say  $\xi_n, n \in \mathbb{N} := \{1, 2, \dots\}$  converges to  $\xi$  in the **vague topology**, if

$$
\int_{S} f(x)\xi_n(dx) \to \int_{S} f(x)\xi(dx), \quad \forall f \in \mathcal{C}_c(S),
$$

where  $\mathcal{C}_c(S)$  is the set of all continuous real-valued functions with compact support. A point process on S is a Conf(S)-valued random variable  $\Xi = \Xi(\cdot)$ . If  $\Xi({x}) \in \{0,1\}$  for any point  $z \in S$ , then the point process is said to be **simple**. Assume that  $\Lambda_i$ ,  $i = 1, \ldots, m$ ,  $m \in \mathbb{N}$  are disjoint bounded sets in S and  $k_i \in \mathbb{N}_0, i = 1, \ldots, m$  satisfy  $\sum_{i=1}^m k_i = n \in \mathbb{N}_0$ . A symmetric measure  $\lambda^n$  on  $S^n$  is called the  $n$ -th **correlation measure**, if it satisfies

$$
\mathbf{E}\left[\prod_{i=1}^m\frac{\Xi(\Lambda_i)!}{(\Xi(\Lambda_i)-k_i)!}\right]=\lambda^n(\Lambda_1^{k_1}\times\cdots\times\Lambda_m^{k_m}),
$$

where when  $\Xi(\Lambda_i) - k_i < 0$ , we interpret  $\Xi(\Lambda_i)! / (\Xi(\Lambda_i) - k_i)! = 0$ . If  $\lambda^n$  is absolutely continuous with respect to the *n*-product measure  $\lambda^{\otimes n}$ , the Radon–Nikodym derivative  $\rho^{n}(x_1,\ldots,x_n)$  is called the *n*-point correlation function with respect to the reference measure  $\lambda$ ; that is,

$$
\lambda^{n}(dx_1\cdots dx_n)=\rho^{n}(x_1,\ldots,x_n)\lambda^{\otimes n}(dx_1\cdots dx_n).
$$

## 2.3.2 Permanental-determinantal point processes (PDPPs) with hierarchical structures

We introduce the following notation. For an  $n \times n$  matrix  $M = (m_{ij})_{1 \leq i,j \leq n}$ ,

<span id="page-49-3"></span><span id="page-49-0"></span>
$$
perfect M = perfect[m_{ij}] := per M det M,
$$
\n(2.67)

that is, perdet M denotes per M multiplied by det M. Note that perdet is a special case of hy**perdeterminants** introduced by Gegenbauer following Cayley [\[36\]](#page-88-1). If  $M$  is a positive semidefinite Hermitian matrix, then per  $M \geq \det M \geq 0$ , and hence perdet  $M \geq 0$  by the definition [\(2.67\)](#page-49-0). The following is the main theorem in Section [2.](#page-35-2)

**Theorem 2.11** The zero point process  $\mathcal{Z}_{X_{\mathbb{A}_q}^r}$  on  $\mathbb{A}_q$  is a **permanental-determinantal point process** (PDPP) in the sense that it has correlation functions  $\{\rho_{A_q}^n\}_{n\in\mathbb{N}}$  given by

<span id="page-49-1"></span>
$$
\rho_{\mathbb{A}_q}^n(z_1,\ldots,z_n;r) = \frac{\theta(-r)}{\theta(-r\prod_{k=1}^n |z_k|^4)} \operatorname{perdet}_{1 \le i,j \le n} \left[ S_{\mathbb{A}_q}\left(z_i,z_j;r\prod_{\ell=1}^n |z_\ell|^2\right) \right]
$$
(2.68)

for every  $n \in \mathbb{N}$  and  $z_1, \ldots, z_n \in \mathbb{A}_q$  with respect to  $m/\pi$ .

The **density of zeros on**  $\mathbb{A}_q$  with respect to  $m/\pi$  is given by

<span id="page-49-5"></span>
$$
\rho_{\mathbb{A}_q}^1(z;r) = \frac{\theta(-r)}{\theta(-r|z|^4)} S_{\mathbb{A}_q}(z,z;r|z|^2)^2 = \frac{q_0^4 \theta(-r,-r|z|^4)}{\theta(-r|z|^2,|z|^2)^2}, \quad z \in \mathbb{A}_q,
$$
\n(2.69)

which is always positive. Since  $\rho_{\mathbb{A}_q}^1(z; r)$  depends only on the modulus of the coordinate  $|z| \in (q, 1)$ , the PDPP is rotationally invariant. As shown by  $(2.34)$ – $(2.36)$  in Section [2.1.2,](#page-40-5) in the interval  $x \in (-\infty, 0), \theta(x)$  is positive and strictly convex with  $\lim_{x \downarrow -\infty} \theta(x) = \lim_{x \uparrow 0} \theta(x) = +\infty$ , while in the interval  $x \in (q^2, 1)$ ,  $\theta(x)$  is positive and strictly concave with  $\theta(x) \sim q_0^2(x - q^2)/q^2$  as  $x \downarrow q^2$ and  $\theta(x) \sim q_0^2(1-x)$  as  $x \uparrow 1$ . Therefore, the density shows divergence both at the inner and outer boundaries as

<span id="page-49-4"></span>
$$
\rho_{\mathbb{A}_q}^1(z;r) \sim \begin{cases} \frac{q^2}{(|z|^2 - q^2)^2}, & |z| \downarrow q, \\ \frac{1}{(1 - |z|^2)^2}, & |z| \uparrow 1, \end{cases}
$$
\n(2.70)

which is independent of r and implies  $\mathbf{E}[\mathcal{Z}_{X_{\mathbb{A}_q}^r}(\mathbb{A}_q)] = \infty$ . See Fig[.4.](#page-50-0)

If M is a 2×2 matrix, we see that perdet  $M = \det(M \circ M)$ , where  $M \circ M$  denotes the Hadamard product of M, i.e., entrywise multiplication,  $(M \circ M)_{ij} = M_{ij}M_{ij}$ . Then the two-point correlation is expressed by a single determinant as

<span id="page-49-6"></span>
$$
\rho_{\mathbb{A}_q}^2(z_1, z_2; r) = \frac{\theta(-r)}{\theta(-r|z_1|^4|z_2|^4)} \det_{1 \le i, j \le 2} \left[ S_{\mathbb{A}_q}(z_i, z_j; r|z_1|^2|z_2|^2)^2 \right], \quad z_1, z_2 \in \mathbb{A}_q. \tag{2.71}
$$

<span id="page-49-2"></span>Remark 2.12 (i) The PDPP with correlation functions [\(2.68\)](#page-49-1) turns out to be a simple point process, i.e., there is no multiple point a.s., due to the existence of two-point correlation function with respect to the Lebesgue measure  $m/\pi$  [\[28,](#page-88-4) Lemma 2.7]. (ii) Using the explicit expression [\(2.68\)](#page-49-1) together with the Frobenius determinantal formula [\(2.90\)](#page-56-0) given below, we can verify that for every  $n \in \mathbb{N}$ , the *n*-point correlation  $\rho_{\mathbb{A}_q}^n(z_1,\ldots,z_n) > 0$  if all coordinates  $z_1,\ldots,z_n \in \mathbb{A}_q$ are different from each other, and that  $\rho_{\mathbb{A}_q}^n(z_1,\ldots,z_n) = 0$  if some of  $z_1,\ldots,z_n$  coincide; e.g.,  $z_i = z_i, i \neq j$ , by the determinantal factor in perdet [\(2.67\)](#page-49-0).



<span id="page-50-0"></span>Figure 4: Dependence on the radial coordinate  $r = |z| \in (q, 1)$  of the density function  $\rho_{\mathbb{A}_q}^1(z; r)$  in the case  $q = r = 1/3$ .

Remark 2.13 The determinantal point processes (DPPs) and the permanental point processes (PPPs) have the n-correlation functions of the forms

$$
\rho_{\text{DPP}}^n(z_1,\ldots,z_n) = \det_{1 \le i,j \le n} [K(z_i,z_j)], \quad \rho_{\text{PPP}}^n(z_1,\ldots,z_n) = \operatorname{per}_{1 \le i,j \le n} [K(z_i,z_j)],
$$

respectively (cf.  $[77, 26]$  $[77, 26]$ ). Due to **Hadamard's inequality** for the determinant  $[51, Section II.4]$ and Lieb's inequality for the permanent [\[49\]](#page-89-3), we have

$$
\rho_{\text{DPP}}^2(z_1, z_2) \leq \rho_{\text{DPP}}^1(z_1)\rho_{\text{DPP}}^1(z_2), \quad \rho_{\text{PPP}}^2(z_1, z_2) \geq \rho_{\text{PPP}}^1(z_1)\rho_{\text{PPP}}^1(z_2).
$$

These correlation inequalities suggest a repulsive nature (negative correlation) for DPPs and an attractive nature (positive correlation) for PPPs. Since perdet is considered to have intermediate nature between determinant and permanent, PDPPs are expected to exhibit both repulsive and attractive characters, depending on the position of two points  $z_1$  and  $z_2$ . For example, Remark [2.12](#page-49-2) (ii) shows the repulsive nature inherited from the DPP side. The two-sidedness of the present PDPP will be summarized in Theorem [2.22](#page-57-0) given below.

**Remark 2.14** The  $(q, r)$ -inversionb symmetry given by Proposition [2.8](#page-47-2) (ii) is rephrased using correlation functions as

<span id="page-50-1"></span>
$$
\rho_{\mathbb{A}_q}^n(T_q(z_1),\ldots,T_q(z_n);r)\prod_{\ell=1}^n|T'_q(z_\ell)|^2=\rho_{\mathbb{A}_q}^n(z_1,\ldots,z_n;q^2/r)
$$
\n(2.72)

for any  $n \in \mathbb{N}$  and  $z_1, \ldots, z_n \in \mathbb{A}_q$ , where  $T_q(z) = q/z$  and  $|T'_q(z)|^2 = q^2/|z|^4$ . In the correlation functions  $\{\rho_{A_q}^n\}_{n\in\mathbb{N}}$  given by Theorem [2.11,](#page-49-3) we see an **inductive structure** such that the functional form of the permanental-determinantal correlation kernel  $S_{\mathbb{A}_q}(\cdot,\cdot;r\prod_{\ell=1}^n|z_\ell|^2)$  is depending on the points  $\{z_1, \ldots, z_n\}$ , which we intend to measure by  $\rho_{\mathbb{A}_q}^n$ , via the weight parameter  $r \prod_{\ell=1}^n |z_\ell|^2$ . This is due to the inductive structure of the induced GAFs generated in conditioning of zeros as mentioned in Remark [2.10.](#page-47-3) In addition, the reference measure  $m/\pi$  is also weighted by  $\theta(-r)/\theta(-r\prod_{k=1}^n |z_k|^4)$ . Such a **hierarchical structure** of correlation functions and reference measures is necessary to realize the  $(q, r)$ -inversion symmetry  $(2.72)$  and the invariance under conformal transformations preserving  $\mathbb{A}_q$  when  $r = q$ .

## 2.3.3 Simpler but still non-trivial PDPP in  $q \to 0$  limit

The above GAF and the PDPP induce the following limiting cases. With fixed  $r > 0$  we take the limit  $q \to 0$ . By the reason explained in Remark [2.16](#page-52-0) below, in this limiting procedure, we should consider the point processes  $\{\mathcal{Z}_{X_{\mathbb{A}_q}^r}: q > 0\}$  to be defined on the **punctured unit disk** 

$$
\mathbb{D}^{\times} := \{ z \in \mathbb{C} : 0 < |z| < 1 \}
$$

instead of  $\mathbb{D}$ . Although the limit point process is given on  $\mathbb{D}^{\times}$  by definition, it can be naturally viewed as a point process defined on  $\mathbb{D}$ , which we will introduce below. Let  $H_r^2(\mathbb{D})$  be the **Hardy** space on  $\mathbb D$  with the weight parameter  $r > 0$ , whose inner product is given by

$$
\langle f, g \rangle_{H_r^2(\mathbb{D})} = \frac{1}{2\pi} \int_0^{2\pi} f(e^{\sqrt{-1}\phi}) \overline{g(e^{\sqrt{-1}\phi})} d\phi + rf(0)g(0), \quad f, g \in H_r^2(\mathbb{D}).
$$

The reproducing kernel of  $H_r^2(\mathbb{D})$  is given by

$$
S_{\mathbb{D}}(z, w; r) = \sum_{n=0}^{\infty} e_n^{(0,r)}(z) \overline{e_n^{(0,r)}(w)} = \frac{1}{1+r} + \sum_{n=1}^{\infty} (z\overline{w})^n
$$

$$
= \frac{1+r z \overline{w}}{(1+r)(1-z\overline{w})}, \quad z, w \in \mathbb{D}.
$$
 (2.73)

The GAF associated with  $H_r^2(\mathbb{D})$  is then defined by

<span id="page-51-2"></span><span id="page-51-0"></span>
$$
X_{\mathbb{D}}^r(z) = \frac{\zeta_0}{\sqrt{1+r}} + \sum_{n=1}^{\infty} \zeta_n z^n, \quad z \in \mathbb{D}
$$
\n(2.74)

so that the covariance kernel is given by  $\mathbf{E}[X_{\mathbb{D}}^r(z)\overline{X_{\mathbb{D}}^r(w)}] = S_{\mathbb{D}}(z,w;r)$ ,  $z, w \in \mathbb{D}$ . For the conditional GAF given a zero at  $\alpha \in \mathbb{D}$ , the covariance kernel is given by

$$
S_{\mathbb{D}}^{\alpha}(z,w;r) = S_{\mathbb{D}}(z,w;r|\alpha|^2)h_{\alpha}(z)\overline{h_{\alpha}(w)}, \quad z, w, \alpha \in \mathbb{D},
$$

where the replacement of the weight parameter r by  $r|\alpha|^2$  should be done, even though the factor  $h_{\alpha}(z)$  is simply given by the Möbius transformation [\(2.61\)](#page-45-0).

For the zero point process Theorem [2.11](#page-49-3) is reduced to the following by the formula

$$
\lim_{q \to 0} \theta(z; q^2) = 1 - z.
$$

**Corollary 2.15** Assume that  $r > 0$ . Then  $\mathcal{Z}_{X_{\mathbb{D}}^r}$  is a PDPP on  $\mathbb{D}$  with the correlation functions

<span id="page-51-3"></span>
$$
\rho_{\mathbb{D}}^n(z_1,\ldots,z_n;r) = \frac{1+r}{1+r\prod_{k=1}^n |z_k|^4} \operatorname{perdet}_{1\leq i,j\leq n} \left[ S_{\mathbb{D}}\left(z_i,z_j;r\prod_{\ell=1}^n |z_\ell|^2\right) \right]
$$
(2.75)

for every  $n \in \mathbb{N}$  and  $z_1, \ldots, z_n \in \mathbb{D}$  with respect to  $m/\pi$ . In particular, the density of zeros on  $\mathbb{D}$  is given by

<span id="page-51-1"></span>
$$
\rho_{\mathbb{D}}^1(z;r) = \frac{(1+r)(1+r|z|^4)}{(1+r|z|^2)^2(1-|z|^2)^2}, \quad z \in \mathbb{D}.
$$
\n(2.76)

As r increases the first term in  $(2.74)$ , which gives the value of the GAF at the origin, decreases and hence the variance at the origin,  $\mathbf{E}[|X_{\mathbb{D}}^r(0)|^2] = S_{\mathbb{D}}(0,0;r) = (1+r)^{-1}$  decreases monotonically. As a result the density of zeros in the vicinity of the origin increases as  $r$  increases. Actually we see that  $\rho_{\mathbb{D}}^1(0;r) = 1+r$ .

<span id="page-52-0"></span>**Remark 2.16** The asymptotics [\(2.70\)](#page-49-4) show that the density of zeros of  $\mathcal{Z}_{X^r_{A_q}}$  diverges at the inner boundary  $\gamma_q = \{z : |z| = q\}$  for each  $q > 0$  while the density of  $\mathcal{Z}_{X_{\mathbb{D}}}$  is finite at the origin as in [\(2.76\)](#page-51-1). Therefore infinitely many zeros near the inner boundary  $\gamma_q$  seem to vanish in the limit as  $q \to 0$ . This is the reason why we regard the base space of  $\{\mathcal{Z}_{X_{\mathbb{A}_q}^r} : q > 0\}$  and the limit point process  $\mathcal{Z}_{X^r_{\mathbb{D}}}$  as  $\mathbb{D}^\times$  instead of  $\mathbb D$  as mentioned before. Indeed, in the vague topology, with which we equip a configuration space, we cannot see configurations outside each compact set, hence infinitely many zeros are not observed on each compact set in  $\mathbb{D}^{\times}$  (not  $\mathbb{D}$ ) for any sufficiently small  $q > 0$ depending on the compact set that we take.

#### 2.3.4 Reduction to GAF and DPP of Peres and Virág

If we take the further limit  $r \to 0$  in [\(2.73\)](#page-51-2), we obtain the Szegő kernel of  $\mathbb D$  given by [\(2.63\)](#page-46-1). Since the matrix  $(S_{\mathbb{D}}(z_i, z_j)^{-1})_{1 \leq i,j \leq n} = (1 - z_i \overline{z_j})_{1 \leq i,j \leq n}$  has rank 2, the following equality called Borchardt's identity holds (see Theorem 3.2 in [\[58\]](#page-90-2), Theorem 5.1.5 in [\[26\]](#page-88-2)),

$$
\operatorname{perfect}_{1 \le i,j \le n} \left[ (1 - z_i \overline{z_j})^{-1} \right] = \det_{1 \le i,j \le n} \left[ (1 - z_i \overline{z_j})^{-2} \right]. \tag{2.77}
$$

By the relation [\(2.64\)](#page-46-2), the  $r \to 0$  limit of  $\mathcal{Z}_{X_{\mathbb{D}}}$  is identified with the DPP on  $\mathbb{D}, \mathcal{Z}_{X_{\mathbb{D}}}$ , studied by Peres and Virág [\[63\]](#page-90-1), whose correlation functions are given by

$$
\rho_{\mathbb{D},\mathrm{PV}}^n(z_1,\ldots,z_n)=\det_{1\leq i,j\leq n}[K_{\mathbb{D}}(z_i,z_j)],\quad n\in\mathbb{N},\quad z_1,\ldots,z_n\in\mathbb{D},
$$

with respect to  $m/\pi$ 

Here we give the result by Peres and Virág as a theorem.

**Theorem 2.17 (Peres and Virág [\[63\]](#page-90-1))**  $\mathcal{Z}_{X_{\mathbb{D}}}$  is a DPP on  $\mathbb{D}$  such that the correlation kernel with respect to  $m/\pi$  is given by the Bergman kernel  $K_{\mathbb{D}}$  of  $\mathbb D$  given by [\(2.4\)](#page-36-1).

Remark 2.18 We see from [\(2.73\)](#page-51-2) that

$$
\lim_{r \to \infty} S_{\mathbb{D}}(z, w; r) = (1 - z\overline{w})^{-1} - 1, \quad z, w \in \mathbb{D},
$$

which can be identified with the conditional kernel given a zero at the origin;  $S_{\mathbb{D}}^0(z, w) = S_{\mathbb{D}}(z, w) S_{\mathbb{D}}(z,0)S_{\mathbb{D}}(0,w)/S_{\mathbb{D}}(0,0)$  for  $S_{\mathbb{D}}(z,0) \equiv 1$ . In this limit we can use Borchardt's identity again, since the rank of the matrix  $(S_{\mathbb{D}}(z_i, z_j; \infty)^{-1})_{1 \leq i,j \leq n} = (z_i^{-1}z_j^{-1} - 1)_{1 \leq i,j \leq n}$  is two. Then, thanks to the proper limit of the prefactor of perdet in [\(2.75\)](#page-51-3) when  $z_k \in \mathbb{D}^\times$  for all  $k = 1, 2, \ldots, n$ ;

$$
\lim_{r \to \infty} \frac{1+r}{1+r \prod_{k=1}^{n} |z_k|^4} = \prod_{k=1}^{n} |z_k|^{-4},
$$

we can verify that

$$
\lim_{r \to \infty} \rho_{\mathbb{D}}^n(z_1, \dots, z_n; r) = \rho_{\mathbb{D},\text{PV}}^n(z_1, \dots, z_n)
$$

for every  $n \in \mathbb{N}$ , and every  $z_1, \ldots, z_n \in \mathbb{D}^{\times}$ . On the other hand, taking  $(2.74)$  into account, we have

$$
X_{\mathbb{D}}^{\infty}(z) = z \sum_{n=1}^{\infty} \zeta_n z^{n-1} \stackrel{\mathrm{d}}{=} z X_{\mathbb{D}}(z),
$$

from which, we can see that

$$
\lim_{r \to \infty} \mathcal{Z}_{X_{\mathbb{D}}^r} =: \mathcal{Z}_{X_{\mathbb{D}}^{\infty}} \stackrel{\mathrm{d}}{=} \mathcal{Z}_{X_{\mathbb{D}}} + \delta_0,
$$

that is, the DPP of Peres and Virág with a deterministic zero added at the origin. This is consistent with the fact that  $\rho_{\mathbb{D}}^1(0; r) = 1 + r$  diverges as  $r \to \infty$ . Since

$$
\lim_{r \to 0} \mathcal{Z}_{X^r_{\mathbb{D}}} =: \mathcal{Z}_{X^0_{\mathbb{D}}} \stackrel{\text{d}}{=} \mathcal{Z}_{X_{\mathbb{D}}}
$$

as mentioned above, the one-parameter family of PDPPs  $\{\mathcal{Z}_{X_{\mathbb{D}}^{r}}: r \in (0, \infty)\}\)$  can be regarded as an **interpolation** between the DPP of Peres and Virág and that DPP with a deterministic zero added at the origin.

The distribution of  $\mathcal{Z}_{X_{\mathbb{D}}}$  is invariant under Möbius transformations that preserve  $\mathbb{D}$  [\[63\]](#page-90-1). This invariance is a special case of the following, which can be proved using the conformal transformations of the Szegő kernel and the Bergman kernel given by  $(2.7)$  [\[63,](#page-90-1) [26\]](#page-88-2).

**Proposition 2.19 (Peres and Virág [\[63\]](#page-90-1))** Let  $\widetilde{D} \subsetneq \mathbb{C}$  be a simply connected domain wit analytic boundary. Then there is a GAF  $X_{\tilde{D}}$  with covariance kernel  $\mathbf{E}[X_{\tilde{D}}(z)\overline{X_{\tilde{D}}(w)}] = S_{\tilde{D}}(z, w),$  $z, w \in D$ , where  $S_{\tilde{D}}$  denotes the Szegő kernel of D. The zero point process  $\mathcal{Z}_{X_{\tilde{D}}}$  is the DPP such that the correlation kernel is given by the Bergman kernel  $K_{\tilde{D}}$  of  $\tilde{D}$ . This DPP is conformally invariant in the following sense. If  $D \subsetneq \mathbb{C}$  is another simply connected domain with analytic boundary, and

$$
f:D\to D
$$

is a conformal transformation, then for  $\mathcal{Z}_{X_D} = \sum_i \delta_{Z_i}$ ,

$$
\sum_i \delta_{f^{-1}(Z_i)} =: f^* \mathcal{Z}_{X_{\tilde{D}}} \stackrel{\text{d}}{=} \mathcal{Z}_{X_D}.
$$

That is,  $f^* \mathcal{Z}_{X_{\tilde{D}}}$  is a DPP such that the correlation kernel is equal to the Bergman kernel  $K_D$  of D. We can say that the zero point process of Peres and Virág  $[63]$  is a **conformal invariant DPP**.

## 2.4 Proof of Theorem [2.11](#page-49-3)

## 2.4.1 Edelman-Kostlan formulas and second log-derivatives of Szegő kernels

For the DPP of Peres and Virág the **Edelman–Kostlan formula** [\[20\]](#page-88-5) gives the density of  $\mathcal{Z}_{X_{\text{D}}}$ with respect to  $m/\pi$  as

$$
\rho_{\mathbb{D},\text{PV}}^1(z) = \frac{1}{4} \Delta \log S_{\mathbb{D}}(z,z), \quad z \in \mathbb{D},
$$

where  $\Delta := 4\partial_z\partial_{\overline{z}}$ . Moreover, we have the equality

<span id="page-53-0"></span>
$$
K_{\mathbb{D}}(z,w) = \partial_z \partial_{\overline{w}} \log S_{\mathbb{D}}(z,w) = S_{\mathbb{D}}(z,w)^2, \quad z, w \in \mathbb{D}.
$$
 (2.78)

As explained above [\(2.11\)](#page-38-5), this gives an example of the general formula

<span id="page-54-3"></span><span id="page-54-1"></span>
$$
S_D(z, w)^2 = K_D(z, w), \quad z, w \in D,
$$
\n(2.79)

which is establised for the kernels on any simply connected domain  $D \subsetneq \mathbb{C}$ .

Here we have reported our work to generalize the above to a family of GAFs and their zero point processes on the annulus  $\mathbb{A}_q$ . We can prove the following [\[36\]](#page-88-1).

**Proposition 2.20** For  $r > 0$ , the following equality holds,

<span id="page-54-2"></span>
$$
\partial_z \partial_{\overline{w}} \log S_{\mathbb{A}_q}(z, w; r) = \frac{\theta(-r)}{\theta(-r(z\overline{w})^2)} S_{\mathbb{A}_q}(z, w; rz\overline{w})^2, \quad z, w \in \mathbb{A}_q. \tag{2.80}
$$

In particular,

<span id="page-54-0"></span>
$$
\Delta \log S_{\mathbb{A}_q}(z, z; r) = 4 \frac{\theta(-r)}{\theta(-r|z|^4)} S_{\mathbb{A}_q}(z, z; r|z|^2)^2, \quad z \in \mathbb{A}_q.
$$
 (2.81)

By comparing the expression [\(2.69\)](#page-49-5) for the density obtained from Theorem [2.11](#page-49-3) with [\(2.81\)](#page-54-0) in Proposition [2.20,](#page-54-1) we can recover the Edelman–Kostlan formula as follows,

$$
\rho_{\mathbb{A}_q}^1(z;r) = \frac{\theta(-r)}{\theta(-r|z|^4)} S_{\mathbb{A}_q}(z,z;r|z|^2)^2 = \frac{1}{4} \Delta \log S_{\mathbb{A}_q}(z,z;r), \quad z \in \mathbb{A}_q.
$$

However, [\(2.78\)](#page-53-0) does not hold for the weighted Szegő kernel for  $H_r^2(\mathbb{A}_q)$ . As shown by [\(2.80\)](#page-54-2), the second log-derivative of  $S_{\mathbb{A}_q}(z, w; r)$  cannot be expressed by  $S_{\mathbb{A}_q}(z, w; r)$  itself but a new function  $S_{\mathbb{A}_q}(z, w; r z \overline{w})$  should be introduced.

In addition the proportionality between the square of the Szegő kernel and the Bergman kernel [\(2.79\)](#page-54-3), which holds as [\(2.79\)](#page-54-3) on  $\mathbb{D}$ , is no longer valid for the point processes on  $\mathbb{A}_q$  The following are verified [\[36\]](#page-88-1). A CONS for the **Bergman space on**  $\mathbb{A}_q$  is given by  $\{\tilde{e}_n^{(q)}(z)\}_{n\in\mathbb{Z}}$  where we set

<span id="page-54-4"></span>
$$
\widetilde{e}_n^{(q)}(z) = \begin{cases} \sqrt{\frac{n+1}{1-q^{2(n+1)}}} z^n, & n \in \mathbb{Z} \setminus \{-1\}, \\ \sqrt{\frac{1}{-2\log q}} z^{-1}, & n = -1. \end{cases}
$$

The Bergman kernel of  $\mathbb{A}_q$  is then given by

$$
K_{\mathbb{A}_q}(z, w) := k_{L^2_{\mathbb{B}}(\mathbb{A}_q)}(z, w) = \sum_{n \in \mathbb{Z}} \tilde{e}_n^{(q)}(z) \overline{\tilde{e}_n^{(q)}(w)}
$$
  
= 
$$
-\frac{1}{2 \log q} \frac{1}{z \overline{w}} + \frac{1}{z \overline{w}} \sum_{n \in \mathbb{Z} \setminus \{0\}} \frac{n}{1 - q^{2n}} (z \overline{w})^n, \quad z, w \in \mathbb{A}_q.
$$
 (2.82)

The Weierstrass  $\wp$ -function is defined by

$$
\wp(\phi) = -\frac{1}{12} + 2\sum_{n=1}^{\infty} \frac{q^{2n}}{(1-q^{2n})^2} - \sum_{n=-\infty}^{\infty} \frac{e^{\sqrt{-1}\phi} q^{2n}}{(1 - e^{\sqrt{-1}\phi} q^{2n})^2}.
$$
\n(2.83)

Then  $(2.82)$  is expressed as as [\[9\]](#page-87-3)

$$
K_{\mathbb{A}_q}(z,w) = -\frac{1}{2\log q} \frac{1}{z\overline{w}} - \frac{1}{z\overline{w}} \left( \wp(\phi_{z\overline{w}}) + \frac{P(q)}{12} \right), \quad z, w \in \mathbb{A}_q,
$$
\n
$$
(2.84)
$$

where we have used the notation,

$$
z = e^{\sqrt{-1}\phi_z} \iff \phi_z = -\sqrt{-1}\log z,\tag{2.85}
$$

and

$$
P(q) = 1 - 24 \sum_{n=1}^{\infty} \frac{q^{2n}}{(1 - q^{2n})^2}.
$$

We can prove the equality [\[36\]](#page-88-1)

$$
S_{\mathbb{A}_q}(z,w)^2 = K_{\mathbb{A}_q}(z,w) + \frac{a(q)}{z\overline{w}}, \quad z, w \in \mathbb{A}_q,
$$
\n
$$
(2.86)
$$

where

$$
a(q) = -2\sum_{n \in \mathbb{N}} \frac{(-1)^n n q^n}{1 - q^{2n}} + \frac{1}{2 \log q}.
$$
 (2.87)

## 2.4.2 Hammersley formula and Shirai's Proposition

We recall a general formula for correlation functions of zero point process of a GAF, which is called the Hammersley formula in [\[63\]](#page-90-1), but here we use a slightly different expression given by Shirai (Proposition 6.1 of [\[75\]](#page-91-3)). Let  $\partial_z \partial_{\overline{w}} := \frac{\partial^2}{\partial z \partial_{\overline{w}}}$  $\frac{\partial^2}{\partial z \partial \overline{w}}$ .

<span id="page-55-0"></span>**Proposition 2.21 (Shirai [\[75\]](#page-91-3))** The correlation functions of  $\mathcal{Z}_{X_D}$  of the GAF  $X_D$  on  $D \subsetneq \mathbb{C}$ with covariance kernel  $S_D(z, w)$  are given by

$$
\rho_D^n(z_1,\ldots,z_n)=\frac{\text{per}_{1\leq i,j\leq n}\left[(\partial_z\partial_{\overline{w}}S_D^{z_1,\ldots,z_n})(z_i,z_j)\right]}{\det_{1\leq i,j\leq n}\left[S_D(z_i,z_j)\right]},\quad n\in\mathbb{N},\quad z_1,\ldots,z_n\in D,
$$

with respect to a reference measure  $\lambda$ , whenever  $\det_{1 \leq i,j \leq n} [S_D(z_i, z_j)] > 0$ . Here  $\{S_D^{z_1,...,z_n}\}$  $D^{(z_1,...,z_n}(z_i,z_j)\}$ denote the conditional kernels.

Here we abbreviate  $\gamma_f^q$  $\{z_{\ell}\}_{\ell=1}^n$  given by [\(2.53\)](#page-44-6) to  $\gamma_n^q$ . Then [\(2.54\)](#page-44-3) gives

$$
S_{\mathbb{A}_q}^{z_1,...,z_n}(z,w;r) = S_{\mathbb{A}_q}(z,w;r) \prod_{\ell=1}^n |z_\ell|^2 \gamma_n^q(z) \overline{\gamma_n^q(w)}
$$

for  $z, w, z_1, \ldots, z_n \in A_q$ . By [\(2.57\)](#page-45-1) of Lemma [2.7,](#page-45-2) this formula gives

$$
(\partial_z \partial_{\overline{w}} S_{\mathbb{A}_q}^{z_1,\ldots,z_n})(z_i,z_j;r) = S_{\mathbb{A}_q}\Big(z_i,z_j;r \prod_{\ell=1}^n |z_\ell|^2\Big) \gamma_n^{q'}(z_i) \overline{\gamma_n^{q'}(z_j)}.
$$

Therefore, Proposition [2.21](#page-55-0) gives now

<span id="page-55-1"></span>
$$
\rho_{\mathbb{A}_q}^n(z_1,\ldots,z_n;r) = \frac{\text{per}_{1\leq i,j\leq n} \left[S_{\mathbb{A}_q}\left(z_i,z_j;r\prod_{\ell=1}^n|z_{\ell}|^2\right)\right] \prod_{k=1}^n |\gamma_n^{q'}(z_k)|^2}{\det_{1\leq i,j\leq n} \left[S_{\mathbb{A}_q}(z_i,z_j;r)\right]}.
$$
\n(2.88)

By [\(2.49\)](#page-44-0) and [\(2.57\)](#page-45-1) and [\(2.60\)](#page-45-3) of Lemma [2.7,](#page-45-2) we see that

$$
\prod_{i=1}^{n} |\gamma_n^{q'}(z_i)|^2 = \prod_{i=1}^{n} \left| \left( \prod_{1 \le j \le n, j \ne i} h_{z_j}^q(z_i) \right) h_{z_i}^{q'}(z_i) \right|^2 = \prod_{i=1}^{n} \left| \left( \prod_{1 \le j \le n, j \ne i} z_i \frac{\theta(z_j/z_i)}{\theta(\overline{z_j}z_i)} \right) \frac{q_0^2}{\theta(|z_i|^2)} \right|^2
$$
\n
$$
= \left| \frac{q_0^{2n} \prod_{1 \le i < j \le n} z_i \theta(z_j/z_i) \cdot \prod_{1 \le i' < j' \le n} z_j \cdot \theta(z_{i'}/z_{j'})}{\prod_{i=1}^{n} \prod_{j=1}^{n} \theta(z_i \overline{z_j})} \right|^2.
$$
\nBy (2.27),  $z_i \theta(z_j/z_i) = z_i(-z_j/z_i) \theta(z_i/z_j) = -z_j \theta(z_i/z_j)$ . Hence this is written as

<span id="page-56-1"></span>
$$
\prod_{i=1}^{n} |\gamma_n^{q'}(z_i)|^2 = q_0^{4n} \left| \frac{(-1)^{n(n-1)/2} \left( \prod_{1 \le i < j \le n} z_j \theta(z_i/z_j) \right)^2}{\prod_{i=1}^{n} \prod_{j=1}^{n} \theta(z_i \overline{z_j})} \right|^2
$$
\n
$$
= q_0^{4n} \left( \frac{\prod_{1 \le i < j \le n} |z_j|^2 \theta(z_i/z_j, \overline{z_i}/\overline{z_j})}{\prod_{i=1}^{n} \prod_{j=1}^{n} \theta(z_i \overline{z_j})} \right)^2.
$$
\n(2.89)

The following identity is known as the Frobenius determinantal formula, which is an elliptic extension of the Cauchy determinantal formula due Frobenius (see Theorem 1.1 in [\[27\]](#page-88-6), Theorem 66 in  $[45]$ , Corollary 4.7 in  $[67]$ ),

$$
\det_{1\leq i,j\leq n}\left[\frac{\theta(tx_ia_j)}{\theta(t,x_ia_j)}\right]=\frac{\theta(t\prod_{k=1}^nx_ka_k)}{\theta(t)}\frac{\prod_{1\leq i
$$

By [\(2.39\)](#page-42-1) in Proposition [2.4,](#page-42-3) we have

<span id="page-56-0"></span>
$$
q_0^{2n} \frac{\prod_{1 \le i < j \le n} |z_j|^2 \theta(z_i/z_j, \overline{z_i}/\overline{z_j})}{\prod_{i=1}^n \prod_{j=1}^n \theta(z_i \overline{z_j})} = \frac{\theta(-s)}{\theta(-s) \prod_{\ell=1}^n |z_\ell|^2} \det_{1 \le i, j \le n} \left[ S_{\mathbb{A}_q}(z_i, z_j; s) \right], \quad \forall s > 0. \tag{2.90}
$$

Then [\(2.89\)](#page-56-1) is written as

$$
\prod_{i=1}^{n} |\gamma_n^{q'}(z_i)|^2 = \frac{\theta(-r)}{\theta(-r\prod_{\ell=1}^n |z_\ell|^2)} \det_{1 \le i,j \le n} [S_{\mathbb{A}_q}(z_i, z_j; r)]
$$
\n
$$
\times \frac{\theta(-r\prod_{\ell=1}^n |z_\ell|^2)}{\theta(-r\prod_{\ell=1}^n |z_\ell|^4)} \det_{1 \le i,j \le n} \left[ S_{\mathbb{A}_q}(z_i, z_j; r \prod_{\ell=1}^n |z_\ell|^2) \right]
$$
\n
$$
= \frac{\theta(-r)}{\theta(-r\prod_{\ell=1}^n |z_\ell|^4)} \det_{1 \le i,j \le n} [S_{\mathbb{A}_q}(z_i, z_j; r)] \det_{1 \le i,j \le n} \left[ S_{\mathbb{A}_q}(z_i, z_j; r \prod_{\ell=1}^n |z_\ell|^2) \right].
$$

Applying the above to [\(2.88\)](#page-55-1), the correlation functions in Theorem [2.11](#page-49-3) are obtained.

# 2.5 Unfolded 2-correlation function of PDPP

By the determinantal factor in perdet [\(2.67\)](#page-49-0) the PDPP shall be negatively correlated when distances of points are short in the domain  $\mathbb{A}_q$ . The effect of the permanental part in perdet will appear in long distances. In order to clarify this fact, we study the two-point correlation function normalized by the product of one-point functions,

$$
g_{\mathbb{A}_q}(z, w; r) := \frac{\rho_{\mathbb{A}_q}^2(z, w; r)}{\rho_{\mathbb{A}_q}^1(z; r)\rho_{\mathbb{A}_q}^1(w; r)}, \quad (z, w) \in \mathbb{A}_q^2,
$$
\n(2.91)

where  $\rho_{\mathbb{A}_q}^1$  and  $\rho_{\mathbb{A}_q}^2$  are explicitly given by [\(2.69\)](#page-49-5) and [\(2.71\)](#page-49-6), respectively. This function is simply called an intensity ratio in [\[63\]](#page-90-1), but here we call it an unfolded 2-correlation function following a terminology used in random matrix theory [\[21\]](#page-88-7).

We have proved the following (see Fig[.5\)](#page-57-1) [\[36\]](#page-88-1).

- **Theorem 2.22** (i) When  $0 < q < 1$ , in the short distance, the correlation is generally repulsive in common with DPPs.
- (ii) There exists a critical value

<span id="page-57-0"></span>
$$
r_0 = r_0(q) \in (q, 1)
$$
 for each  $q \in (0, 1)$ 

such that if  $r \in (r_0, 1)$  positive correlation emerges between zeros when the distance between them is large enough within  $\mathbb{A}_q$ .

(iii) The limits  $g_{\mathbb{D}}(z,w;r) := \lim_{q\to 0} g_{\mathbb{A}_q}(z,w;r)$ ,  $z, w \in \mathbb{D}^\times$  and  $r_{\mathbb{C}} := \lim_{q\to 0} r_0(q)$  are welldefined, and  $r_c$  is positive. When  $r \in [0, r_c)$  all positive correlations vanish in  $g_{\mathbb{D}}(z, w; r)$ , while when  $r \in (r_c, \infty)$  positive correlations can survive.



<span id="page-57-1"></span>Figure 5: The critical curve  $r = r_0(q)$  mentioned in Theorem [2.22](#page-57-0) (ii) is numerically plotted (in red) in the fundamental cell  $\Omega := \{(q, r) : q \in (0, 1), q \leq r \leq 1\}$  in the parameter space, which is located between the diagonal line  $r = q$  and the horizontal line  $r = 1$ . The parabolic curve  $r_c + cq^2$ given in [\[36\]](#page-88-1) and the line  $1 - (1 - q)/2$  are also dotted, which approximate  $r = r_0(q)$  well for  $q \gtrsim 0$ and  $q \leq 1$ , respectively.

## 2.6 Exercises 2

## 2.6.1 Exercise 2.1

(1) By the definition of the theta function

$$
\theta(z;p) := (z, p/z; p)_{\infty} = (z; p)_{\infty} (p/z; p)_{\infty}
$$

$$
= \prod_{i=1}^{\infty} \left\{ (1 - zp^i)(1 - p^{i+1}/z) \right\},\tag{2.92}
$$

prove the inversion formula

<span id="page-58-0"></span>
$$
\theta(1/z;p) = -\frac{1}{z}\theta(z;p)
$$
\n(2.93)

and the quasi-periodicity property

$$
\theta(pz;p) = -\frac{1}{z}\theta(z;p). \tag{2.94}
$$

(2) By direct calculation following the definition [\(2.92\)](#page-58-0), verify the equalities,

$$
\theta(\zeta;p^3)\theta(\zeta p;p^3)\theta(\zeta p^2;p^3) = \theta(\zeta;p),\tag{2.95}
$$

$$
\theta(\zeta;p)\theta(\zeta\omega_3;p)\theta(\zeta\omega_3^2;p) = \theta(\zeta^3;p^3), \quad \zeta \in \mathbb{C}^\times,
$$
\n(2.96)

where  $\omega_3$  is a primitive 3rd root of unity.

(3) Let  $x = e^{-2ia}$ ,  $y = e^{-2ib}$ ,  $u = e^{-2ic}$ , and  $v = e^{-2id}$  in the Weierstrass addition formula

$$
\theta(xy, x/y, uv, u/v; p) - \theta(xv, x/v, uy, u/y; p) = \frac{u}{y}\theta(yv, y/v, xu, x/u; p).
$$
 (2.97)

Consider the limit  $p \to 0$  and derive the following trigonometric equality,

$$
\sin(a+b)\sin(a-b)\sin(c+d)\sin(c-d) \n-\sin(a+d)\sin(a-d)\sin(c+b)\sin(c-b) \n=\sin(b+d)\sin(b-d)\sin(a+c)\sin(a-c).
$$
\n(2.98)

Show that [\(2.98\)](#page-58-1) is readily verified if we use the following addition formulas of trigonometric functions,

<span id="page-58-1"></span>
$$
\sin(a \pm b) = \sin a \cos b \pm \cos a \sin b. \tag{2.99}
$$

## 2.6.2 Exercise 2.2

Consider the weighted Szegő kernel of  $\mathbb{A}_q := \{z \in \mathbb{C} : q < |z| < 1\}$ ,  $q \in (0,1)$  expressed using the theta function

$$
S_{\mathbb{A}_q}(z, w; r) = \frac{q_0^2 \theta(-rz\overline{w})}{\theta(-r, z\overline{w})}, \quad z, w \in \mathbb{A}_q, \quad r > 0,
$$
\n(2.100)

where

$$
q_0 := \prod_{n \in \mathbb{N}} (1 - q^{2n}) = (q^2; q^2)_{\infty}.
$$
\n(2.101)

Schottky's theorem asserts that the group of conformal transformations from  $\mathbb{A}_q$  to itself is generated by the rotations

$$
R_{\theta}(z) := e^{\sqrt{-1}\theta} z, \quad \theta \in [0, 2\pi), \tag{2.102}
$$

and the q-inversions

$$
T_q(z) := \frac{q}{z}.\tag{2.103}
$$

- (1) Show that the rotational invariance is obvious.
- (2) Prove the following equality,

$$
\sqrt{T_q'(z)}\sqrt{T_q'(w)}S_{\mathbb{A}_q}(T_q(z), T_q(w); r) = \frac{q}{r}S_{\mathbb{A}_q}(z, w; q^2/r).
$$
\n(2.104)

Note that if  $r = q$  the weighted Szegő kernel is reduced to the original Szegő kernel,  $S_{\mathbb{A}_q}(z, w) =$  $S_{\mathbb{A}_q}(z, w; q)$ . In this case the above becomes

$$
\sqrt{T_q'(z)}\sqrt{T_q'(w)}S_{\mathbb{A}_q}(T_q(z),T_q(w)) = S_{\mathbb{A}_q}(z,w).
$$
\n(2.105)

Hence we can conclude that the original Szegő kernel,  $S_{A_q}(z, w) = S_{A_q}(z, w; q)$ , is invariant under conformal transformations from  $\mathbb{A}_q$  to itself.

## 2.6.3 Exercise 2.3

Consider the DPP of Peres and Virág,  $\mathcal{Z}_{X_{\mathbb{D}}}$ , whose correlation kernel with respect to  $m/\pi$  is given by the Bergman kernel of  $\mathbb{D}, K_{\mathbb{D}};$ 

$$
\rho_{\mathbb{D},\text{PV}}^n(z_1,\ldots,z_n) = \det_{1\leq i,j\leq n}[K_{\mathbb{D}}(z_i,z_j)], \quad n \in \mathbb{N}, \quad z_1,\ldots,z_n \in \mathbb{D},
$$
  
with 
$$
K_{\mathbb{D}}(z,w) = \frac{1}{(1-z\overline{w})^2}, \quad z,w \in \mathbb{D}.
$$
 (2.106)

(1) Prove the following equality,

$$
\rho_{\mathbb{D},\text{PV}}^1(z) = \frac{1}{4} \Delta \log S_{\mathbb{D}}(z,z), \quad z \in \mathbb{D},\tag{2.107}
$$

where  $\Delta := 4\partial_z\partial_{\overline{z}}$  and  $S_{\mathbb{D}}$  denotes the Szegő kernel of D,

$$
S_{\mathbb{D}}(z,w) = \frac{1}{1 - z\overline{w}}, \quad z, w \in \mathbb{D}.
$$
 (2.108)

This equality is called the Edelman–Kostlan formula.

(2) Show that the following general equalities are established,

$$
K_{\mathbb{D}}(z,w) = \partial_z \partial_{\overline{w}} \log S_{\mathbb{D}}(z,w) = S_{\mathbb{D}}(z,w)^2, \quad z, w \in \mathbb{D}.
$$
 (2.109)

#### 2.6.4 Exercise 2.4

Consider the GAF and the permanenta-determinantal point process on  $\mathbb{A}_q$  discussed there.

(1) Prove the equalities,

$$
\partial_z \partial_{\overline{w}} \log S_{\mathbb{A}_q}(z, w; r) = \frac{\theta(-r)}{\theta(-r(z\overline{w})^2)} S_{\mathbb{A}_q}(z, w; rz\overline{w})^2, \quad z, w \in \mathbb{A}_q. \tag{2.110}
$$

In particular,

$$
\Delta \log S_{\mathbb{A}_q}(z, z; r) = 4 \frac{\theta(-r)}{\theta(-r|z|^4)} S_{\mathbb{A}_q}(z, z; r|z|^2)^2, \quad z \in \mathbb{A}_q.
$$
 (2.111)

(2) The correlation function of the permanental-detaeminantal point process on  $\mathbb{A}_q$  is generally given by

$$
\rho_{\mathbb{A}_q}^n(z_1,\ldots,z_n;r) = \frac{\theta(-r)}{\theta(-r\prod_{k=1}^n |z_k|^4)} \operatorname{perdet}_{1 \le i,j \le n} \left[ S_{\mathbb{A}_q}\left(z_i,z_j;r\prod_{\ell=1}^n |z_\ell|^2\right) \right] \tag{2.112}
$$

for every  $n \in \mathbb{N}$  and  $z_1, \ldots, z_n \in \mathbb{A}_q$  with respect to  $m/\pi$ . Show that the Edelman–Kostran formula is also satisfied.

# 3 Multiple Schramm–Loewner Evolution/Gaussain Free Field Coupling

## 3.1 Multiple Schramm–Loewner Evolution

## 3.1.1 Loewner equations for single-slit and multi-slit

Let D be a simply connected domain in  $\mathbb C$  which does not complete the plane;  $D \subsetneq \mathbb C$ . Its boundary is denoted by  $\partial D$ . We consider a slit in D, which is defined as a trace  $\eta = {\eta(t) : t \in (0, \infty)}$  of a simple curve  $\eta(t) \in D, 0 < t < \infty$ ;

$$
\eta(s) \neq \eta(t) \quad \text{for } s \neq t.
$$

We assume that the initial point of the slit is located in  $\partial D$ ,  $\exists \eta(0) := \lim_{t \to 0} \eta(t) \in \partial D$ . Let

$$
\eta(0,t]:=\{\eta(s):s\in(0,t]\}\quad\text{and}\quad D_t^\eta:=D\setminus\eta(0,t], t\in(0,\infty).
$$

The Loewner theory [\[50\]](#page-89-5) describes a slit  $\eta$  by encoding the curve into a time-dependent analytic function  $g_{D_t^{\eta}}: t \in (0, \infty)$  such that

$$
g_{D_t^\eta} : \text{conformal map } D_t^\eta \to D, \quad t \in (0,\infty).
$$

By the Riemann mapping theorem, for  $D \subsetneq \mathbb{C}$  and a point  $z_0 \in D$ , there exists a unique analytic function  $\varphi(z)$  in D, specified by  $\varphi(z_0) = 0, \varphi'(z_0) > 0$ , such that

 $\varphi:$  conformal map  $D \to \mathbb{D}$ ,

where  $\mathbb D$  denotes a unit disk;  $\mathbb D := \{z \in \mathbb C : |z| < 1\}$  (see [\(2.8\)](#page-37-3) in the previous section). Loewner gave a differential equation for  $g_{D_t^{\eta}}$  in the case  $D = \mathbb{D}$ , which is called the Loewner equation [\[50\]](#page-89-5).

Since a special case of the Möbius transformation

$$
\mathsf{m}(z):=\sqrt{-1}\frac{\alpha-z}{\alpha+z},\quad |\alpha|=1,
$$

maps D to the upper half complex plane  $\mathbb{H} := \{z \in \mathbb{C} : \text{Im } z > 0\}$  with  $m(0) = \sqrt{-1}, m(\infty) =$ −  $\sqrt{-1}$ , we can apply the Loewner theory to the case with  $D = \mathbb{H}$ , in which [\[46\]](#page-89-6)

$$
\eta := \{ \eta(t) : t \in (0, \infty) \} \text{ is a simple curve,}
$$

$$
\eta(0) := \lim_{t \to 0} \eta(t) = 0 \in \mathbb{R},
$$

$$
\eta(0, t] \subset \mathbb{H}, \quad \forall t \in (0, \infty),
$$

$$
\lim_{t \to \infty} \eta(t) = \infty.
$$

For each time  $t \in (0, \infty)$ ,

<span id="page-61-4"></span>
$$
\mathbb{H}_t^{\eta} := \mathbb{H} \setminus \eta(0, t] \tag{3.1}
$$

is a simply connected domain in  $\mathbb C$  and there exists a unique analytic function  $g_{\mathbb{H}_t^{\eta}}$  such that

 $g_{\mathbb{H}_t^{\eta}}:$  conformal map  $\mathbb{H}_t^{\eta} \to \mathbb{H}$ ,

which satisfies the condition

$$
g_{\mathbb{H}_t^{\eta}}(z) = z + \frac{c_t}{z} + \mathcal{O}(|z|^{-2}) \quad \text{as } z \to \infty
$$

for some  $c_t > 0$ , in which the coefficient of z is unity and no constant term appears. This is called the hydrodynamic normalization. The coefficient  $c_t$  gives the half-plane capacity of  $\eta(0, t]$ and denoted by  $hcap(\eta(0, t])$ . The following has been shown (see [\[46,](#page-89-6) [47\]](#page-89-7)).

<span id="page-61-3"></span>**Theorem 3.1** Let  $\eta$  be a slit in H for which the parameterization by t is arranged so that

$$
c_t = \text{hcap}(\eta(0, t]) = 2t, \quad t \in (0, \infty).
$$

Then there exists a unique continuous **driving function**  $V(t) \in \mathbb{R}$ ,  $t \in (0,\infty)$  such that the solution  $g_t$  of the differential equation

<span id="page-61-0"></span>
$$
\frac{dg_t(z)}{dt} = \frac{2}{g_t(z) - V(t)}, \quad t \ge 0,
$$
\n(3.2)

under the initial condition

$$
g_0(z)=z\in\mathbb{H}
$$

gives  $g_t = g_{\mathbb{H}_t^{\eta}}, t \in (0, \infty)$ .

The equation [\(3.2\)](#page-61-0) is called the **chordal Loewner equation**. Note that at each time  $t \in (0, \infty)$ , the tip of slit  $\eta(t)$  and the value of  $V(t)$  satisfy the following relations,

<span id="page-61-1"></span>
$$
V(t) = \lim_{\substack{z \to 0, \\ \eta(t) + z \in \mathbb{H}_t^{\eta}}} g_{\mathbb{H}_t^{\eta}}(\eta(t) + z) \iff \eta(t) = \lim_{\substack{z \to 0, \\ z \in \mathbb{H}}} g_{\mathbb{H}_t^{\eta}}^{-1}(V(t) + z), \quad t \ge 0.
$$
 (3.3)

Moreover,  $V(t) = \lim_{s \le t, s \to t} g_{\mathbb{H}_s^{\eta}}(\eta(t))$  and  $t \to V(t)$  is continuous (see, for instance, Lemma 4.2 in [\[47\]](#page-89-7)). We write

$$
g_{\mathbb{H}_t^{\eta}}(\eta(t)) = V(t) \in \mathbb{R}, \quad \Longleftrightarrow \quad \eta(t) = g_{\mathbb{H}_t^{\eta}}^{-1}(V(t)) \in \partial \mathbb{H}_t^{\eta}, \quad t \ge 0,
$$
\n(3.4)

<span id="page-61-2"></span>in the sense of [\(3.3\)](#page-61-1).

**Example 3.2** When the driving function is identically zero;  $V(t) \equiv 0, t \in (0,\infty)$ , the chordal Loewner equation  $dg_{\mathbb{H}_t^\eta}(z)/dt = 2/g_{\mathbb{H}_t^\eta}(z), t \geq 0$  is solved under the initial condition  $g_{\mathbb{H}_0^\eta}(z) =$  $z \in \mathbb{H}$  as  $g_{\mathbb{H}_t^{\eta}}(z)^2 = 4t + z^2, t \ge 0$ . In this simple case, [\(3.3\)](#page-61-1) gives  $\eta(t) = 2\sqrt{-1}t^{1/2}, t \ge 0$ . That is, the slit  $\eta(0, t], t > 0$  is a straight line along the imaginary axis starting from the origin,  $\eta(0) = \lim_{t \to 0} \eta(t) = 0$ , and growing upward as time t is passing.

**Example 3.3** The above example can be extended by introducing one parameter  $\alpha \in (0,1)$  as follows. Let  $\kappa = \kappa(\alpha) = 4(1 - 2\alpha)^2/\{\alpha(1 - \alpha)\}\,$ , and consider the case such that

$$
V(t) = \begin{cases} \sqrt{\kappa t}, & \text{if } \alpha \le 1/2, \\ -\sqrt{\kappa t}, & \text{if } \alpha > 1/2. \end{cases}
$$

In this case, the inverse of  $g_t$  is solved as  $g_{\mathbb{H}_t}^{-1}(z) = \left(z + 2\sqrt{\frac{\alpha}{1-\alpha}}\right)$  $\sqrt{t}\bigg)^{1-\alpha}\left(z-2\sqrt{\frac{1-\alpha}{\alpha}}\right)$  $\sqrt{t}$ , and the slit is obtained as

$$
\eta(t) = g_{\mathbb{H}_t^{\eta}}^{-1}(V(t)) = 2\left(\frac{1-\alpha}{\alpha}\right)^{1/2-\alpha} e^{\sqrt{-1}\alpha\pi} t^{1/2}, \quad t \ge 0.
$$

The slit grows from the origin along a straight line in  $\mathbb{H}$  which makes an angle  $\alpha\pi$  with respect to the positive direction of the real axis. When  $\alpha = 1/2$ , this is reduced to the result mentioned in Example [3.2.](#page-61-2) More detail for this example, see Example 4.12 in [\[47\]](#page-89-7) and Section 2.2 in [\[31\]](#page-88-8).

Theorem [3.1](#page-61-3) can be extended to the situation such that  $\eta$  in H is given by a multi-slit [\[68,](#page-90-4) [34\]](#page-88-9). Let  $N \in \mathbb{N} := \{1, 2, \dots\}$  and assume that we have N slits  $\eta_i = \{\eta_i(t) : t \in (0, \infty)\} \subset \mathbb{H}, i = 1, \dots, N$ , which are simple curves, disjoint with each other,  $\eta_i \cap \eta_j = \emptyset, i \neq j$ , starting from N distinct points  $\lim_{t\to 0} \eta_i(t) =: \eta_i(0)$  on  $\mathbb{R}; \eta_1(0) < \cdots < \eta_N(0)$ , and all going to infinity;  $\lim_{t\to \infty} \eta_i(t) = \infty$ ,  $i = 1, \ldots, N$ . A **multi-slit** is defined as a union of them,  $\bigcup_{i=1}^{N} \eta_i$ , and

$$
\mathbb{H}_t^{\eta} := \mathbb{H} \setminus \bigcup_{i=1}^N \eta_i(0, t] \tag{3.5}
$$

for each  $t > 0$  with  $\mathbb{H}_0^{\eta}$  $\eta_0^n := \mathbb{H}$ . For each time  $t \in (0, \infty)$ ,  $\mathbb{H}_t^n$  $t$ <sup>''</sup> is a simply connected domain in  $\mathbb C$  and then there exists a unique analytic function  $g_{\mathbb{H}_t^{\eta}}$  such that

$$
g_{\mathbb{H}_t^{\eta}}: \text{conformal map } \mathbb{H}_t^{\eta} \to \mathbb{H},
$$

satisfying the hydrodynamic normalization condition

$$
g_{\mathbb{H}_t^{\eta}}(z) = z + \frac{\text{hcap}(\bigcup_{i=1}^N \eta_i(0, t])}{z} + O(|z|^{-2}) \text{ as } z \to \infty.
$$

<span id="page-62-0"></span>**Theorem 3.4** For  $N \in \mathbb{N}$ , let  $\bigcup_{i=1}^{N} \eta_i$  be a multi-slit in  $\mathbb{H}$  such that

$$
h\mathrm{cap}\Big(\bigcup_{i=1}^N \eta(0,t]\Big) = 2Nt, \quad t \in (0,\infty).
$$

Then there exists a set of weight functions  $w_i(t) \geq 0, t \geq 0, i = 1, ..., N$  satisfying  $\sum_{i=1}^{N} w_i(t) =$  $1, t \geq 0$  and an N-variate continuous driving function  $\mathbf{V}(t) = (V_1(t), \dots, V_N(t)) \in \mathbb{R}^N, t \in (0, \infty)$ such that the solution  $g_t$  of the differential equation

<span id="page-63-0"></span>
$$
\frac{dg_t(z)}{dt} = \sum_{i=1}^N \frac{2Nw_i(t)}{g_t(z) - V_i(t)}, \quad t \ge 0, \quad g_0(z) = z,
$$
\n(3.6)

gives  $g_t = g_{\mathbb{H}_t^{\eta}}, t \in (0, \infty)$ .

Roth and Schleissinger [\[68\]](#page-90-4) called [\(3.6\)](#page-63-0) the Loewner equation for the multi-slit  $\bigcup_{i=1}^{N} \eta_i$ . Similar to [\(3.3\)](#page-61-1), the following relations hold,

$$
V_i(t) = \lim_{\substack{z \to 0, \\ \eta_i(t) + z \in \mathbb{H}_t^{\eta}}} g_{\mathbb{H}_t^{\eta}}(\eta_i(t) + z) \iff \eta_i(t) = \lim_{\substack{z \to 0, \\ z \in \mathbb{H}}} g_{\mathbb{H}_t^{\eta}}^{-1}(V_i(t) + z), \quad i = 1, \dots, N, \quad t \ge 0,
$$
\n(3.7)

and we write for the **multiple tips**,  $\eta_i(t)$ ,  $i = 1, \ldots, N$ ,  $t \geq 0$ ,

<span id="page-63-1"></span>
$$
g_{\mathbb{H}_t^n}(\eta_i(t))=V_i(t)\in\mathbb{R},\quad i=1,\ldots,N,\quad t\geq 0
$$

in the sense of [\(3.7\)](#page-63-1).

## 3.1.2 Schramm–Loewner evolution with parameter  $\kappa$  (SLE<sub> $\kappa$ </sub>)

So far we have considered the problem such that, given time-evolution of a single slit  $\eta(0, t], t \geq 0$ or a multi-slit  $\bigcup_{i=1}^{N} \eta(0, t], t \ge 0$  in  $\mathbb{H}$ , time-evolution of the conformal map from  $\mathbb{H}_{t}^{\eta}$  $t_t^{\eta}$  to  $\mathbb{H}, t \geq 0$ is asked. The answers are given by the solution of the Loewner equation [\(3.2\)](#page-61-0) in Theorem [3.1](#page-61-3) for a single slit and by the solution of the multiple Loewner equation [\(3.6\)](#page-63-0) in Theorem [3.4](#page-62-0) for a multi-slit, which are driven by a function  $(V(t))_{t>0}$  and by a multi-variate function  $V(t)$  =  $(V_1(t),\ldots,V_N(t))\in\mathbb{R}^N, t\geq 0$ , respectively. The both processes are defined in R and deterministic:

single slit  $\eta(0, t] \in \mathbb{H}, t \ge 0 \implies$  driving function  $(V(t))_{t \ge 0}$  on  $\mathbb{R}$ multi-slit  $\bigcup_{i=1}^{N} \eta(0, t] \in \mathbb{H}, t \ge 0 \implies$  multi-variate driving function  $(\boldsymbol{V}(t))_{t \ge 0}$  on  $\mathbb{R}^N$ 

For  $\mathbb H$  with a single slit, Schramm considered an inverse problem in a probabilistic setting [\[69\]](#page-90-5). He first asked a suitable family of driving stochastic processes  $(Y(t))_{t\geq0}$  on R. Then he asked the probability law of a random slit in  $\mathbb{H}$ , which will be determined by the relations  $(3.3)$  from  $(Y(t))_{t\geq 0}$  and the solution  $g_t = g_{\mathbb{H}_t^n}, t \geq 0$  of the Loewner equation [\(3.2\)](#page-61-0):

**random** curve  $\eta(0, t] \in \mathbb{H}, t \ge 0 \iff$  driving **stochastic** process  $(Y(t))_{t \ge 0}$  on  $\mathbb{R}$ 

Schramm argued that **conformal invariance** implies that the driving process  $(Y(t))_{t\geq0}$  should be a continuous Markov process which has in a particular parameterization independent increments. Hence  $Y(t)$  can be a constant time change of a one-dimensional standard Brownian motion  $(B(t))_{t>0}$ , and it is expressed as

$$
(\sqrt{\kappa}B(t))_{t\geq 0} \stackrel{\text{(law)}}{=} (B(\kappa t))_{t\geq 0} \quad \text{with a parameter } \kappa > 0. \tag{3.8}
$$

The solution of the Loewner equation driven by  $Y(t) = \sqrt{\kappa}B(t), t \ge 0$ ,

<span id="page-64-0"></span>
$$
\frac{dg_{\mathbb{H}_t^{\eta}}(z)}{dt} = \frac{2}{g_{\mathbb{H}_t^{\eta}}(z) - \sqrt{\kappa}B(t)}, \quad t \ge 0, \quad g_{\mathbb{H}_0^{\eta}}(z) = z \in \mathbb{H},
$$
\n(3.9)

is called the **chordal Schramm–Loewner evolution** (chordal SLE) with parameter  $\kappa > 0$  and is written as  $\mathbf{SLE}_\kappa$  for short.

The following was proved by Lawler, Schramm, and Werner [\[48\]](#page-89-8) for  $\kappa = 8$  and by Rohde and Schramm [\[66\]](#page-90-6) for  $\kappa \neq 8$ .

**Proposition 3.5** By [\(3.3\)](#page-61-1), a chordal  $SLE_{\kappa}$   $g_{\mathbb{H}_t^n}$ ,  $t \in (0,\infty)$  determines a continuous curve  $\eta =$  $\{\eta(t): t \in (0,\infty)\}\subset \mathbb{H}$  such that  $\eta(0):= \lim_{t\downarrow 0} \eta(t) = 0$  and  $\lim_{t\to\infty} |\eta(t)| = \infty$  with probability one.

The continuous curve  $\eta$  determined by an  $SLE_{\kappa}$  is called an  $SLE_{\kappa}$  curve (or  $SLE_{\kappa}$  trace) The probability law of an  $SLE_{\kappa}$  curve depends on  $\kappa$ . As a matter of fact,  $SLE_{\kappa}$  curve becomes self-intersecting and can touch the real axis R when  $\kappa > 4$ , so it is no more a slit, since a slit has been defined as a trace of a continuous simple curve. When  $\kappa > 4$ , the domain  $\mathbb{H} \setminus \eta(0, t]$  is divided into many components, only one of which is unbounded. So here we change the definition [\(3.1\)](#page-61-4) of  $\mathbb{H}^\eta_t$  $t, t \geq 0$  as follows,

$$
\mathbb{H}_t^{\eta} := \text{the unbounded component of } \mathbb{H} \setminus \eta(0, t], \quad t \ge 0. \tag{3.10}
$$

Then

$$
g_{\mathbb{H}_t^{\eta}}(z) : \text{conformal map } \mathbb{H}_t^{\eta} \to \mathbb{H}, \quad t \geq 0.
$$

We also define

$$
K_t^{\eta} := \overline{\mathbb{H} \setminus \mathbb{H}_t^{\eta}}, \quad t \ge 0,
$$
\n(3.11)

<span id="page-64-1"></span>and call it the SLE hull.

There are **three phases of an**  $SLE_{\kappa}$  **curve** as shown by the follows (see Fig[.6\)](#page-65-0).

**Proposition 3.6** There are two critical values of  $\kappa$ ;

$$
\kappa_{\rm c} = 4 \quad \text{and} \quad \overline{\kappa}_{\rm c} = 8.
$$

- (a) If  $0 < \kappa \le \kappa_c = 4$ , then the SLE<sub> $\kappa$ </sub> curve is **simple** and  $\eta = \eta(0, \infty) \subset \mathbb{H}$  with probability one.
- (b) If  $\kappa_c = 4 < \kappa < \overline{\kappa}_c = 8$ , the SLE<sub> $\kappa$ </sub> curve is **self-intersecting**,  $\eta \cap \mathbb{R} \neq \emptyset$ , and touch the real R with positive probability. Then

$$
\bigcup_{t\in [0,\infty)} K^\eta_t = \overline{\mathbb{H}},
$$

but  $\eta[0,\infty) \cap \mathbb{H} \neq \mathbb{H}$ .

(c) If  $\kappa \ge \overline{\kappa}_c = 8$ , then  $\eta$  is a **space-filling** curve;

$$
\eta[0,\infty)=\overline{\mathbb{H}}.
$$



<span id="page-65-0"></span>Figure 6: Schematic pictures of  $SLE_{\kappa}$  curves in (a) Phase 1 (0 <  $\kappa \leq \kappa_c = 4$ ), (b) Phase 2  $(\kappa_c = 4 < \kappa < \overline{\kappa}_c = 8)$ , and (c) Phase 3 ( $\kappa \geq \overline{\kappa}_c = 8$ ).

Remark 3.7 Consider BES<sub>D</sub> introduced in Section [1.4.](#page-10-0) We set  $Z^z(t) = X^z(t) + \sqrt{-1}Y^z(t) \in$  $\overline{\mathbb{H}} \setminus \{0\}, t \geq 0$  and complexificate [\(1.45\)](#page-11-0) as

<span id="page-65-2"></span>
$$
dZ^{z}(t) = dB(t) + \frac{D-1}{2} \frac{dt}{Z^{z}(t)}
$$
\n(3.12)

with the initial condition

$$
Z^z(0) = z = x + \sqrt{-1}y \in \overline{\mathbb{H}} \setminus \{0\}.
$$

The crucial point of this **complexification of Bessel flow** is that the BM remains real,  $B(t) \in$  $\mathbb{R}, t \geq 0$ . Then, there is an asymmetry between the real part and the imaginary part of the flow in H,

<span id="page-65-1"></span>
$$
dX^{z}(t) = dB(t) + \frac{D-1}{2} \frac{X^{z}(t)}{(X^{z}(t))^{2} + (Y^{z}(t))^{2}} dt,
$$
\n(3.13)

$$
dY^{z}(t) = -\frac{D-1}{2} \frac{Y^{z}(t)}{(X^{z}(t))^{2} + (Y^{z}(t))^{2}} dt.
$$
\n(3.14)

Assume  $D > 1$ . Then as indicated by the minus sign in RHS of [\(3.14\)](#page-65-1), the flow is downward in  $\overline{\mathbb{H}}$ . If the flow goes down and arrives at the real axis, the imaginary part vanishes,  $Y^z(t) = 0$ , then equation [\(3.13\)](#page-65-1) is reduced to be the same equation as equation [\(1.45\)](#page-11-0) for the  $BES<sub>D</sub>$ , which is now considered for  $\mathbb{R} \setminus \{0\} = \mathbb{R}_+ \cup \mathbb{R}_-$ . If  $D > D_c = 2$ , by Theorem [1.2](#page-14-0) (ii), the flow on  $\mathbb{R}\setminus\{0\}$  is asymptotically outward,  $X^z(t) \to \pm \infty$  as  $t \to \infty$ . Therefore, the flow on  $\overline{\mathbb{H}}$  will be described as shown by Fig. [3.1.2.](#page-66-0) The behavior of flow should be, however, more complicated when  $\overline{D}_{c} = 3/2 < D < D_{c}$  and  $1 < D < \overline{D}_{c} = 3/2$ . For  $z \in \overline{\mathbb{H}} \setminus \{0\}, t \geq 0$ , put

$$
\widehat{g}_t(z) := Z^z(t) + B(t). \tag{3.15}
$$

Then, Eq. [\(3.12\)](#page-65-2) is rewritten for  $\hat{g}_t(z)$  as

<span id="page-65-3"></span>
$$
\frac{\partial \hat{g}_t(z)}{\partial t} = \frac{D-1}{2} \frac{1}{\hat{g}_t(z) - B(t)}, \quad t \ge 0.
$$
\n(3.16)



<span id="page-66-0"></span>Figure 7: A schematic picture of **complexificated Bessel flow** on  $\overline{\mathbb{H}} \setminus \{0\}$  for  $D > 2$ .

Put

$$
\kappa = \frac{4}{D - 1} \quad \Longleftrightarrow \quad D = 1 + \frac{4}{\kappa},\tag{3.17}
$$

and set

$$
g_t(z) = \widehat{g}_{\kappa t}(z)
$$

in [\(3.16\)](#page-65-3). Then we have the equation in the form

$$
\frac{\partial g_t(z)}{\partial t} = \frac{2}{g_t(z) - \sqrt{\kappa}B(t)}, \quad t \ge 0.
$$
\n(3.18)

This is equal to the Schramm–Loewner evolution [\(3.9\)](#page-64-0). As a matter of fact, identification of the  $SLE_{\kappa}$  as a complexification of  $BES_{1+4/\kappa}$  gives the proof of Proposition [3.6,](#page-64-1) in which simple correspondence between two sets of critical values;

$$
D_{\rm c} = 2 \iff \kappa_{\rm c} = 4, \qquad \overline{D}_{\rm c} = \frac{3}{2} \iff \overline{\kappa}_{\rm c} = 8. \tag{3.19}
$$

More detailed description of the probability laws of an  $\text{SLE}_\kappa$  curves at special values of  $\kappa$ , see, for instance [\[47,](#page-89-7) [31,](#page-88-8) [40\]](#page-89-9).

The highlight of the theory of SLE would be that, if the value of  $D$  is properly chosen, the probability law of  $\gamma$  realizes that of the scaling limit of important lattice paths studied in a **statistical** mechanics model exhibiting critical phenomena or describing interesting fractal geometry defined on an infinite discrete lattice. The following is a list of the correspondence (up to a conjecture) between the  $SLE_{\kappa}$  paths with specified values of  $\kappa$ , and the names of lattice paths (with the names of models studied in statistical mechanics and fractal physics), whose scaling limits are described by the  $SLE_{\kappa}$  paths. (See also [\[18,](#page-87-4) [17\]](#page-87-5).)



Moreover, Beffara [\[6\]](#page-87-8) determined the Hausdorff dimensions  $d_{\kappa}^{\text{H}}$  of the  $\text{SLE}_{\kappa}$  curves as

$$
d_{\kappa}^{\mathcal{H}} = \begin{cases} 1 + \frac{\kappa}{8}, & \kappa \in (0, 8) \\ 2, & \kappa \ge 8. \end{cases}
$$
 (3.20)

The relationship between the  $\text{SLE}_\kappa$  and the **conformal field theory** has been clarified [\[4\]](#page-87-9). The central charge c and the scaling dimension (the highest weight) of representation of the Virasoro algebra are related with the parameter  $\kappa$  of  $SLE_{\kappa}$  as

$$
c = c_{\kappa} := \frac{(6 - \kappa)(3\kappa - 8)}{2\kappa}, \quad h = h_{\kappa} := \frac{6 - \kappa}{2\kappa}.
$$
 (3.21)

## 3.1.3 Multiple SLE

For simplicity, we assume that  $w_i(t) \equiv 1/N, t \geq 0, i = 1, ..., N$  in [\(3.6\)](#page-63-0) in Theorem [3.4.](#page-62-0) Then the Loewner equation for the multi-slit in  $\mathbb H$  is written as

$$
\frac{dg_{\mathbb{H}_t^{\eta}}(z)}{dt} = \sum_{i=1}^N \frac{2}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)}, \quad t \ge 0, \quad g_{\mathbb{H}_0^{\eta}}(z) = z \in \mathbb{H}.
$$
\n(3.22)

Then we ask what is the suitable family of driving stochastic processes of N particles on  $\mathbb{R}, Y(t) =$  $(Y_1(t), \ldots, Y_N(t)), t \ge 0$  [\[11,](#page-87-10) [5,](#page-87-11) [44,](#page-89-10) [24\]](#page-88-10);

$$
\begin{aligned}\n\textbf{multiple SLE}_{\kappa} \text{ curves } & \bigcup_{i=1}^{N} \eta_i(0, t] \in \mathbb{H}, t \ge 0 \\
& \Longleftrightarrow \quad \text{driving } \textbf{many-particle stochastic process } (\mathbf{Y}(t))_{t \ge 0} \text{ on } \mathbb{R}^N \text{ ?}\n\end{aligned}
$$

The same argument with Schramm [\[69\]](#page-90-5) will give that  $\mathbf{Y}(t)$  should be a continuous Markov process. Moreover, Bauer, Bernard, and Kytölä  $[5]$ , Graham  $[24]$ , and Dubédat  $[15]$  argued that  $(Y_i(t))_{t>0}, i = 1,\ldots,N$  are semi-martingales and the quadratic variations should be given by  $d\langle Y_i, Y_j\rangle_t = \kappa \delta_{ij} dt, t \geq 0, 1 \leq i, j \leq N$  with  $\kappa > 0$ . Then we will be able to assume that the system of SDEs for  $(Y(t))_{t>0}$  is in the form,

$$
dY_i(t) = \sqrt{\kappa}dB_i(t) + F_i(\mathbf{Y}(t))dt, \quad t \ge 0, \quad i = 1, \dots, N,
$$
\n(3.23)

where  $(B_i(t))_{t\geq 0}, i=1,\ldots,N$  are independent one-dimensional standard Brownian motions,  $\kappa > 0$ , and  $\{F_i(\boldsymbol{x})\}_{i=1}^N$  are suitable functions of  $\boldsymbol{x} = (x_1, \ldots, x_N)$  which do not explicitly depend on t.



Figure 8: Schematic picture of multiple SLE curves  $\eta_i, i = 1, \ldots, N$  in  $\mathbb{H}$  driven by  $\boldsymbol{Y}(t) =$  $(Y_1(t),..., Y_N(t)), t \geq 0$ . Here we assume that they are all simple and avoiding from each other in H.

In the following, We will give a *theory* so that the driving process  $(Y(t))_{t\geq0}$  should be a time change of DYS<sub>β</sub> with  $\beta = 8/\kappa$  to construct a *proper* multiple SLE. Here we define the Gaussian free field (GFF) and its generalization called the imaginary surface with parameter  $\chi$ , which are considered as the distribution-valued random fields on H. Under the relation  $\chi = 2/\sqrt{\kappa} - \kappa/\sqrt{2}$ , we regard the SLE/GFF coupling studied by Dubédat, Sheffield, and Miller as a temporally stationary field, and extend it to multiple cases. We prove that the multiple SLE/GFF coupling is established, if and only if the driving N-particle process on  $\mathbb R$  is identified with  $DYS_{8/\kappa}$ .

## 3.2 Gaussian free field (GFF) with Dirichlet boundary condition

## 3.2.1 Bochner–Minlos Theorem

Here we start with the classical **Bochner theorem**, which states that a probability measure on a finite dimensional Euclidean space is determined by a **characteristic function** which is a Fourier transform of the probability measure. Note that we have identified the Fourier transform of the transition probability density of BM [\(1.11\)](#page-5-0) with the characteristic function of BM [\(1.12\)](#page-5-1) in Section [1.1,](#page-2-0) and the multitime Laplace transforms of the joint probability distribution function [\(1.99\)](#page-25-0) with the generating function of correlation functions for  $DYS_2$  in Section [1.7.1.](#page-25-1) First we define a functional of positive type.

**Definition 3.8** Let V be a finite or infinite dimensional vector space. A function  $\psi : V \to \mathbb{C}$  is said to be a functional of positive type if for arbitrary  $N \in \mathbb{N}$ ,  $\xi_1, \ldots, \xi_N \in \mathcal{V}$ , and  $z_1, \ldots, z_N \in \mathbb{C}$ , we have

$$
\sum_{n=1}^{N} \sum_{m=1}^{N} \psi(\xi_n - \xi_m) z_n \overline{z_m} \ge 0.
$$

Then the following is proved.

**Lemma 3.9** Let  $\psi : \mathcal{V} \to \mathbb{C}$  be a functional of positive type on a vector space  $\mathcal{V}$ . Then it follows that (i)  $\psi(0) \geq 0$ , (ii)  $\psi(\xi) = \overline{\psi(-\xi)}$  for all  $\xi \in \mathcal{V}$ , and (iii)  $|\psi(\xi)| \leq \psi(0)$  for all  $\xi \in \mathcal{V}$ .

<span id="page-69-0"></span>For  $x, y \in \mathbb{R}^N$ , the standard inner product is denoted by  $x \cdot y$  and we write  $|x| := \sqrt{x \cdot x}$ . Let  $\mathcal{B}^N$ be the family of Borel sets in  $\mathbb{R}^N$ . Then the following is known as the Bochner theorem.

**Theorem 3.10 (Bochner theorem)** Let  $\psi : \mathbb{R}^N \to \mathbb{C}$  be a continuous functional of positive type such that  $\psi(0) = 1$ . Then there exists a unique probability measure P on  $(\mathbb{R}^N, \mathcal{B}^N)$  such that

$$
\psi(\xi) = \int_{\mathbb{R}^N} e^{\sqrt{-1}x \cdot \xi} P(dx) \text{ for } \xi \in \mathbb{R}^N.
$$

If we consider the case that  $\psi(\xi)$  is given by  $\Psi(\xi) := e^{-|\xi|^2/2}, \xi \in \mathbb{R}^N$ , then the probability measure P given by the Bochner theorem is the  $N$ -dimensional standard Gaussian measure,

$$
P(dx) = \frac{1}{(2\pi)^{N/2}} e^{-|x|^2/2} dx = \prod_{i=1}^N \lambda_{N(0,1)}(dx_i), \quad x = (x_1, \dots, x_N) \in \mathbb{R}^N.
$$

Hence we can say that the finite-dimensional standard Gaussian measure P is determined by the characteristic function  $\Psi(\xi)$  as

$$
\Psi(\xi) = \int_{\mathbb{R}^N} e^{\sqrt{-1}x \cdot \xi} P(dx) = e^{-|\xi|^2/2} \quad \text{for} \quad \xi \in \mathbb{R}^N.
$$

Now consider the case that  $H$  is an infinite dimensional Hilbert space with inner product  $\langle \cdot, \cdot \rangle = \langle \cdot, \cdot \rangle_{\mathcal{H}}$  with  $||x|| = ||x||_{\mathcal{H}} = \sqrt{\langle x, x \rangle_{\mathcal{H}}}, x \in \mathcal{H}$ . The dual space of  $\mathcal{H}$  will be denoted by  $\mathcal{H}^*$ . Suppose that there were a probability measure P on  $\mathcal H$  with a suitable  $\sigma$ -algebra such that

$$
\psi(\xi) = \int_{\mathcal{H}} e^{\sqrt{-1}\langle x,\xi\rangle} P(dx) = e^{-\|\xi\|^2/2} \text{ for } \xi \in \mathcal{H}.
$$

Let  $\{e_n\}_{n=1}^{\infty}$  be a complete orthonormal system (CONS) of H. If we set  $\xi = te_n, t \in \mathbb{R}$  for an arbitrary  $n \in \mathbb{N}$ , then

$$
\int_{\mathcal{H}} e^{\sqrt{-1}t\langle x, e_n \rangle} P(dx) = e^{-t^2/2}, \quad t \in \mathbb{R}.
$$

Since  $x \in \mathcal{H}$ , we have  $\langle x, e_n \rangle \to 0$  as  $n \to \infty$ . Therefore in the limit  $n \to \infty$ , the above equality gives  $e^{-t^2/2} = 1$ , which is a contradiction. This observation suggests that the application of the Bochner theorem to an infinite dimensional space requires more consideration. The following arguments are base on [\[3\]](#page-87-13) and a note given by Koshida [\[43\]](#page-89-11).

Let  $D \subsetneq \mathbb{C}$  be a simply connected proper domain of  $\mathbb{C}$  that is bounded. We consider the case  $\mathcal{H} = L^2(D, \mu(dz))$  with  $\langle f, g \rangle := \int_D f(z)g(z) d\mu(z)$ ,  $f, g \in L^2(D, \mu(dz))$ , where  $\mu(dz)$  is the Lebesgue measure on  $D \subset \mathbb{C}$ ;  $\mu(dz) = d\text{Re } z d\text{Im } z = \sqrt{-1}dz d\overline{z}/2$ . Let  $\Delta$  be the **Dirichlet Laplacian** acting on  $L^2(D, \mu(dz))$ . Then  $-\Delta$  has positive discrete eigenvalues so that

$$
-\Delta e_n = \lambda_n e_n, \quad e_n \in L^2(D, \mu(dz)), \quad n \in \mathbb{N}.
$$
\n(3.24)

We assume that the eigenvalues are labeled in a non-decreasing order;  $0 < \lambda_1 \leq \lambda_2 \leq \cdots$ . The system of eigenvalue functions  $\{e_n\}_{n\in\mathbb{N}}$  forms a **complete orthonormal system** (CONS) of  $L^2(D)$ . The following is known as the Weyl formula

**Lemma 3.11** Let  $D \subsetneq \mathbb{C}$  be a simply connected finite domain. The eigenvalues  $\{\lambda_n\}_{n\in\mathbb{N}}$  of the operator  $-\Delta$  on D exhibit the following asymptotic behavior,

$$
\lim_{n \to \infty} \frac{\lambda_n}{n} = O(1).
$$

For two functions  $f, g \in \mathcal{C}_c^{\infty}(D)$ , their **Dirichlet inner product** is defined as

<span id="page-70-0"></span>
$$
\langle f, g \rangle_{\nabla} := \frac{1}{2\pi} \int_{D} (\nabla f)(z) \cdot (\nabla g)(z) \mu(dz). \tag{3.25}
$$

The Hilbert space completion of  $\mathcal{C}_{\underline{c}}^{\infty}(D)$  with respect to this Dirichlet inner product will be denoted by  $W(D)$ . We write  $||f||_{\nabla} = \sqrt{\langle f, f \rangle_{\nabla}}, f \in W(D)$ . If we set  $u_n = \sqrt{2\pi/\lambda_n} e_n, n \in \mathbb{N}$ , then by integration by parts, we have

$$
\langle u_n, u_n \rangle_{\nabla} = \frac{1}{2\pi} \langle u_n, (-\Delta)u_m \rangle = \delta_{nm}, \quad n, m \in \mathbb{N}.
$$

Therefore  $\{u_n\}_{n\in\mathbb{N}}$  forms a CONS of  $W(D)$ .

Let  $\widehat{\mathcal{H}}(D)$  be the space of formal real infinite series in  $\{u_n\}_{n\in\mathbb{N}}$ . This is obviously isomorphic to  $\mathbb{R}^{\mathbb{N}}$  by setting  $\widehat{\mathcal{H}}(D) \ni \sum_{n \in \mathbb{N}} f_n u_n \mapsto (f_n)_{n \in \mathbb{N}} \in \mathbb{R}^{\mathbb{N}}$ . As a subspace of  $\widehat{\mathcal{H}}(D)$ ,  $W(D)$  is isomorphic to  $\ell^2(\mathbb{N}) \subset \mathbb{R}^{\mathbb{N}}$ . For two formal series  $f = \sum_{n \in \mathbb{N}} f_n u_n, g = \sum_{n \in \mathbb{N}} g_n u_n \in \widehat{\mathcal{H}}(D)$  such that  $\sum_{n\in\mathbb{N}}|f_ng_n|<\infty$ , we define their **pairing** as

$$
\langle f,g\rangle_\nabla:=\sum_{n\in\mathbb{N}}f_ng_n.
$$

In the case when  $f, g \in W(D)$ , their pairing of course coincides with the Dirichlet inner product  $(3.25).$  $(3.25).$ 

Notice that, for any  $a \in \mathbb{R}$ , the operator  $(-\Delta)^a$  acts on  $\widehat{\mathcal{H}}(D)$  as

$$
(-\Delta)^a \sum_{n \in \mathbb{N}} f_n u_n := \sum_{n \in \mathbb{N}} \lambda_n^a f_n u_n, \quad (f_n)_{n \in \mathbb{N}} \in \mathbb{R}^{\mathbb{N}}.
$$

Using this fact, we define  $\mathcal{H}_a(D) := (-\Delta)^a W(D)$ ,  $a \in \mathbb{R}$ , each of which is a Hilbert space with inner product

$$
\langle f, g \rangle_a := \langle (-\Delta)^{-a} f, (-\Delta)^{-a} g \rangle_{\nabla}, \quad f, g \in \mathcal{H}_a(D).
$$

<span id="page-70-1"></span>We write  $\|\cdot\|_a := \sqrt{\langle \cdot, \cdot \rangle_a}, a \in \mathbb{R}$ .

**Example 3.12** When  $a = 1/2$ , we have

<span id="page-70-3"></span>
$$
\langle f, g \rangle_{1/2} = \left\langle (-\Delta)^{-1/2} f, (-\Delta)^{-1/2} g \right\rangle_{\nabla} = \frac{1}{2\pi} \langle f, g \rangle, \quad f, g \in \mathcal{H}_{1/2}(D).
$$

Therefore  $\mathcal{H}_{1/2}(D) = L^2(D, \mu(dz)).$ 

We can prove the following two lemmas.

<span id="page-70-2"></span>**Lemma 3.13** Assume  $a < b$ . Then  $\mathcal{H}_a(D) \subset \mathcal{H}_b(D)$ .

**Lemma 3.14** Let  $a \in \mathbb{R}$  and fix  $h \in \mathcal{H}_a(D)$ . Then the assignment

 $\langle h, \cdot \rangle_{\nabla} : \mathcal{H}_{-a}(D) \to \mathbb{R}$  such that  $\mathcal{H}_{-a}(D) \ni f \mapsto \langle h, f \rangle_{\nabla} \in \mathbb{R}$ 

is well-defined and continuous. In particular,  $\mathcal{H}_a(D)$  and  $\mathcal{H}_{-a}(D)$  makes a dual pair of Hilbert spaces with respect to the Dirichlet inner product  $\langle \cdot, \cdot \rangle_{\nabla}$ .

**Remark 3.15** Since  $\mathcal{H}_{1/2}(D) = L^2(D, \nu(dz))$  as mentioned in Example [3.12,](#page-70-1) the members of  $\mathcal{H}_a(D)$  with  $a > 1/2$  cannot be functions, but are **distributions**.

Define

$$
\mathcal{E}(D) := \bigcup_{a>1/2} \mathcal{H}_a(D). \tag{3.26}
$$

Then its dual Hilbert space is identified with  $\mathcal{E}(D)^* := \bigcap_{a<-1/2} \mathcal{H}_a(D)$  by Lemma [3.14,](#page-70-2) and

 $\mathcal{E}(D)^* \subset W(D) \subset \mathcal{E}(D)$ 

is established (by definition and Lemma [3.13\)](#page-70-3). Here  $(\mathcal{E}(D)^*, W(D), \mathcal{E}(D))$  is called a Gel'fand triple. We set  $\Sigma_{\mathcal{E}(D)} = \sigma(\{\langle \cdot, f \rangle_{\nabla} : f \in \mathcal{E}(D)^*\})$ . On such a setting, the following is obtained. This theorem is the extension of the Bochner theorem (Theorem [3.10\)](#page-69-0) and is called the **Bochner**– Minlos theorem (see, for instance, [\[25,](#page-88-11) [73,](#page-91-6) [3\]](#page-87-13)).

**Theorem 3.16 (Bochner–Minlos theorem)** Let  $\psi$  be a continuous function of positive type on  $W(D)$  such that  $\psi(0) = 1$ . Then there exists a unique probability measure **P** on  $(\mathcal{E}(D), \Sigma_{\mathcal{E}(D)})$ such that

<span id="page-71-0"></span>
$$
\psi(f) = \int_{\mathcal{E}(D)} e^{\sqrt{-1}\langle h, f \rangle_{\nabla}} \mathbf{P}(dh) \quad \text{for } f \in \mathcal{E}(D)^*.
$$
\n(3.27)

Under certain conditions on  $\psi$ , the domain of function f for [\(3.27\)](#page-71-0) can be extended from  $\mathcal{E}(D)^*$ to  $W(D)$ . It is easy to verify that the functional

$$
\Psi(f) := e^{-\|f\|_{\nabla}^2/2} \tag{3.28}
$$

satisfies the conditions. Then the following holds with a probability measure **P** on  $(\mathcal{E}(D), \Sigma_{\mathcal{E}(D)})$ ,

$$
\Psi(f) = \int_{\mathcal{E}(D)} e^{\sqrt{-1}\langle h, f \rangle_{\nabla}} \mathbf{P}(dh) = e^{-\|f\|_{\nabla}^2/2} \quad \text{for } f \in W(D). \tag{3.29}
$$

Definition 3.17 (Dirichlet boundary GFF) A Gaussian free field (GFF) with Dirich**let boundary condition** is defined as a pair  $((\Omega^{\text{GFF}}, \mathcal{F}^{\text{GFF}}, \mathbb{P}^{\text{GFF}}), H)$  of a probability space  $(\Omega^{\text{GFF}}, \mathcal{F}^{\text{GFF}}, \mathbb{P}^{\text{GFF}})$  and an isometry

$$
H: W(D) \to L^2(\Omega^{\mathrm{GFF}}, \mathcal{F}^{\mathrm{GFF}}, \mathbb{P}^{\mathrm{GFF}})
$$

such that each  $H(f)$ ,  $f \in W(D)$  is a Gaussian random variable.
For each  $f \in W(D)$ , [\(3.29\)](#page-71-0) gives a Gaussian random variable  $H(f) := \langle H, f \rangle_{\nabla}$  for the variable  $h \in \mathcal{E}(D)$  by

$$
\mathcal{E}(D) \ni h \mapsto \langle h, f \rangle_{\nabla} \in L^2(\mathcal{E}(D), \Sigma_{\mathcal{E}(D)}, \mathbf{P}).
$$

In this way, [\(3.29\)](#page-71-0) ensures that the pair of of the probability space

$$
(\Omega^{\mathrm{GFF}}, \mathcal{F}^{\mathrm{GFF}}, \mathbb{P}^{\mathrm{GFF}}) := (\mathcal{E}(D), \Sigma_{\mathcal{E}(D)}, \mathbf{P})
$$

with the isometry  $H$ ,  $((\mathcal{E}(D), \Sigma_{\mathcal{E}(D)}, P), H)$  gives a GFF with Dirichlet boundary condition. We often just call H a **Dirichlet boundary GFF** without referring to the probability space  $(\mathcal{E}(D), \Sigma_{\mathcal{E}(D)}, \mathbf{P}).$ By this definition, [\(3.29\)](#page-71-0) is written as

<span id="page-72-0"></span>
$$
\mathbb{E}^{\text{GFF}}[e^{\sqrt{-1}\langle H,f\rangle_{\nabla}}] = e^{-\|f\|_{\nabla}^2/2}, \quad f \in W(D). \tag{3.30}
$$

This determines all the moments of the family of Gaussian random variables  $\{H, f\}\nabla : f \in W(D)\}.$ For examples, the covariance is given by

$$
Cov[\langle H, f \rangle_{\nabla}, \langle H, g \rangle_{\nabla}] := \mathbb{E}^{\text{GFF}} \Big[ \langle H, f \rangle_{\nabla} \langle H, g \rangle_{\nabla} \Big] = \langle f, g \rangle_{\nabla}, \quad f, g \in W(D). \tag{3.31}
$$

In particular, the variance is written as

$$
\text{Var}[\langle H, f \rangle_{\nabla}] := \mathbb{E}^{\text{GFF}} \Big[ \langle H, f \rangle_{\nabla}^2 \Big] = \|f\|_{\nabla}^2, \quad f \in W(D).
$$

<span id="page-72-2"></span>By introducing a parameter  $\theta \in \mathbb{R}$ , and by replacing  $f \to \theta f$  in [\(3.30\)](#page-72-0), the above result implies the following useful formula.

**Lemma 3.18** For the family of **centered (mean zero) Gaussian random variables**  $\{H, f\}\nabla$ :  $f \in W(D)$  given by the Dirichlet boundary GFF,

<span id="page-72-1"></span>
$$
\mathbb{E}^{\text{GFF}}[e^{\sqrt{-1}\theta \langle H, f \rangle_{\nabla}}] = \exp\left(-\frac{\theta^2}{2} \text{Var}\Big[\langle H, f \rangle_{\nabla}\Big]\right), \quad \theta \in \mathbb{R}, \quad f \in W(D). \tag{3.32}
$$

We call [\(3.32\)](#page-72-1) the characteristic function of Dirichlet boundary GFF.

# 3.2.2 Conformal invariance of GFF

<span id="page-72-3"></span>Assume that  $D, D' \subsetneq \mathbb{C}$  are simply connected domains and let  $\varphi : D' \to D$  be a conformal map.

**Lemma 3.19** The Dirichlet inner product  $(3.25)$  is conformally invariant, that is,

$$
\int_D (\nabla f)(z) \cdot (\nabla g)(z) \mu(dz) = \int_{D'} (\nabla (f \circ \varphi))(z) \cdot (\nabla (g \circ \varphi))(z) \mu(dz) \quad \text{for } f, g \in \mathcal{C}_c^{\infty}(D).
$$

From the above lemma, we see that the **pull-back**  $\varphi^* : W(D) \ni f \mapsto f \circ \varphi \in W(D')$  is an isomorphism. This allows one to consider a GFF on an unbounded domain. Namely, if  $D'$  is bounded on which a Dirichlet GFF  $H$  is defined, but  $D$  is unbounded, we can define a family  $\{\langle \varphi_* H, f \rangle_\nabla : f \in W(D)\}\$  by

$$
\langle \varphi_* H, f \rangle_{\nabla} := \langle H, \varphi^* f \rangle_{\nabla}, \quad f \in W(D).
$$

The transformation rule of covariance is given by

$$
Cov[\langle \varphi_* H, f \rangle_{\nabla}, \langle \varphi_* H, g \rangle_{\nabla}] = \mathbb{E}^{GFF} \Big[ \langle \varphi_* H, f \rangle_{\nabla} \langle \varphi_* H, g \rangle_{\nabla} \Big] = \langle \varphi^* f, \varphi^* g \rangle_{\nabla} = \langle f, g \rangle_{\nabla}
$$
  
= 
$$
\mathbb{E}^{GFF} \Big[ \langle H, f \rangle_{\nabla} \langle H, g \rangle_{\nabla} \Big] = Cov[\langle H, f \rangle_{\nabla}, \langle H, g \rangle_{\nabla}].
$$
 (3.33)

Relying on the following formal computation

<span id="page-73-0"></span>
$$
\langle \varphi_* H, f \rangle_{\nabla} = \langle H, \varphi^* f \rangle_{\nabla} = \frac{1}{2\pi} \int_{D'} (\nabla H)(z) \cdot (\nabla f \circ \varphi)(z) \mu(dz)
$$

$$
= \frac{1}{2\pi} \int_{D} (\nabla H \circ \varphi^{-1})(z) \cdot (\nabla f)(z) \mu(dz)
$$

we understand the equality  $\varphi_* H = H \circ \varphi^{-1}$ . By the fact [\(3.33\)](#page-73-0) such that the covariance structure does not change under a conformal map  $\phi$ , we say that the GFF is conformally invariant.

## 3.2.3 The Green's function of GFF

By a formal integration by parts, we see that

$$
\langle H, f \rangle_{\nabla} = \frac{1}{2\pi} \int_{D} (\nabla H)(z) \cdot (\nabla f)(z) \mu(dz) = \frac{1}{2\pi} \int_{D} H(z) (-\Delta f)(z) \mu(dz)
$$

$$
= \frac{1}{2\pi} \langle H, (-\Delta)f \rangle.
$$

Motivated by this observation, we define

$$
\langle H, f \rangle := 2\pi \langle H, (-\Delta)^{-1} f \rangle_{\nabla} \quad \text{for } f \in \mathsf{D}((-\Delta)^{-1}), \tag{3.34}
$$

where  $\mathsf{D}((-\Delta)^{-1})$  denotes the domain of  $(-\Delta)^{-1}$  in  $W(D)$ . Note that if D is bounded, then  $(-\Delta)^{-1}$ is a bounded operator, but if D is unbounded, then  $(-\Delta)^{-1}$  is not defined on  $W(D)$ . The action of  $(-\Delta)^{-1}$  is expressed as an integral operator and the integral kernel is known as the Green's function. Namely,

$$
((-\Delta)^{-1}f)(z) = \frac{1}{2\pi} \int_D G_D(z, w) f(w) \mu(dw), \quad \text{a.e. } z \in D, \quad f \in D((-\Delta)^{-1}),
$$

where  $G_D(z, w)$  denotes the Green's function of D under the Dirichlet boundary condition. Hence the covariance of  $\langle H, f \rangle$  and  $\langle H, g \rangle$  with  $f, g \in D((-\Delta)^{-1})$  is written as

<span id="page-73-1"></span>
$$
\mathbb{E}^{\text{GFF}}[\langle H, f \rangle \langle H, g \rangle] = \int_{D \times D} f(z) G_D(z, w) g(w) \mu(dz) \mu(dw). \tag{3.35}
$$

When we symbolically write

$$
\langle H, f \rangle = \int_D H(z) f(z) \mu(dz), \quad f \in \mathsf{D}((-\Delta)^{-1}),
$$

the covariance structure can be expressed as

$$
\mathbb{E}^{\text{GFF}}[H(z)H(w)] = G_D(z, w), \quad z, w \in D, \quad n \neq w.
$$

<span id="page-73-2"></span>From the formula [\(3.35\)](#page-73-1), we see that  $\mathcal{C}_c^{\infty}(D) \subset D((-\Delta)^{-1})$ . In the following, we will consider the family of random variables  $\{ \langle H, f \rangle : f \in C_c^{\infty}(D) \}$  to characterize the GFF, H.

**Example 3.20** When D is the upper half plane  $\mathbb{H}$ ,

$$
G_{\mathbb{H}}(z, w) = \log \left| \frac{z - \overline{w}}{z - w} \right| = \log |z - \overline{w}| - \log |z - w|
$$

$$
= \text{Re} \log(z - \overline{w}) - \text{Re} \log(z - w),
$$

 $z, w \in \mathbb{H}, z \neq w$ .

**Example 3.21** When D is the first orthant  $\mathbb{O} := \{z : \text{Re } z > 0, \text{Im } z > 0\},\$ 

$$
G_{\mathbb{O}}(z,w) = \log \left| \frac{(z - \overline{w})(z + \overline{w})}{(z - w)(z + w)} \right|
$$
  
=  $\log |z - \overline{w}| + \log |z + \overline{w}| - \log |z - w| - \log |z + w|$ ,  
= Re  $\log(z - \overline{w}) +$  Re  $\log(z + \overline{w}) -$  Re  $\log(z - w) -$  Re  $\log(z + w)$ ,

 $z, w \in \mathbb{O}, z \neq w$ .

The conformal invariance of GFF implies that for a conformal map  $\varphi : D' \to D$ , we have the equality,

$$
G_{D'}(z,w) = G_D(\varphi(z), \varphi(w)), \quad z, w \in D'. \tag{3.36}
$$

In the following, we will regard the upper half plane H as the representative of the simply connected domain  $D \subsetneq \mathbb{C}$ . Since each  $D' \subsetneq \mathbb{C}$  is specified by the conformal transformation  $\varphi : D' \to \mathbb{H}$ , we put this in the superscript and write

<span id="page-74-1"></span><span id="page-74-0"></span>
$$
G^{\varphi}(z,w) := G_{\varphi^{-1}(\mathbb{H})}(z,w) = G_{\mathbb{H}}(\varphi(z), \varphi(w)).
$$
\n(3.37)

By Example [3.20,](#page-73-2)

$$
G^{\varphi}(z, w) = \log |\varphi(z) - \overline{\varphi(w)}| - \log |\varphi(z) - \varphi(w)|
$$
  
= Re log( $\varphi(z) - \overline{\varphi(w)}$ ) - Re log( $\varphi(z) - \varphi(w)$ ). (3.38)

## 3.3 GFF coupled with stochastic log-gase

## 3.3.1 GFF transformed by multiple SLE

We write the probability space of the multiple SLE as  $(\Omega^{\text{SLE}}, \mathcal{F}^{\text{SLE}}, (\mathcal{F}^{\text{SLE}})_{\geq 0}, \mathbb{P}^{\text{SLE}})$ . Now we define the direct product of this space and that for the Dirichlet boundary GFF,

$$
(\Omega, \mathcal{F}, \mathbb{P}) = (\Omega^{\text{SLE}} \times \Omega^{\text{GFF}}, \mathcal{F}^{\text{SLE}} \times \mathcal{F}^{\text{GFF}}, \mathbb{P}^{\text{SLE}} \times \mathbb{P}^{\text{GFF}}).
$$
(3.39)

We consider the multiple SLE as well as GFF in this extended probability space, and the multiple SLE is assumed to be adapted to the filtration,

$$
\mathcal{F}_t = \mathcal{F}_t^{\text{SLE}} \times \{ \emptyset, \Omega^{\text{GFF}} \}, \quad t \ge 0.
$$

We assume that the multiple SLE [\(3.22\)](#page-67-0) driven by  $\mathbf{Y}(t) = (Y_1(t), \ldots, Y_N(t))$  under the initial condition  $g_0(z) = z \in \mathbb{H}$  has a unique solution  $(g_t)_{t>0}$ . Moreover, we assume that this solution can be extended to R and determines

$$
\eta_i := \{\eta_i(t) : t \ge 0\}, \quad i = 1, \dots, N
$$

by the equations

$$
Y_i(t) = g_t(\eta_i(t)), \quad i = 1, \dots, N. \tag{3.40}
$$

We consider the situation such that each  $\eta_i, i = 1, \ldots, N$  $\eta_i, i = 1, \ldots, N$  $\eta_i, i = 1, \ldots, N$  make a continuous curve in  $\overline{\mathbb{H}}^{-1}$ , and for each  $t \geq 0$ , we define

<span id="page-75-5"></span>
$$
\mathbb{H}_t^{\eta} := \text{the unbounded component of } \mathbb{H} \setminus \bigcup_{i=1}^N \eta_i(0, t], \quad t \ge 0. \tag{3.41}
$$

Then the solution of [\(3.22\)](#page-67-0) gives the conformal transformation from  $\mathbb{H}^{\eta}_t$  $t$ <sup>n</sup> to  $\mathbb H$  and hence it is denoted by  $g_{\mathbb{H}_t^{\eta}}$ .

In summary, we assume the following;

$$
\frac{dg_{\mathbb{H}_{t}^{\eta}}(z)}{dt} = \sum_{i=1}^{N} \frac{2}{g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{i}(t)}, \quad t \ge 0, \quad g_{\mathbb{H}_{0}^{\eta}}(z) = g_{\mathbb{H}}(z) = z \in \mathbb{H},
$$
\n
$$
Y_{i}(t) = g_{\mathbb{H}_{t}^{\eta}}(\eta_{i}(t)) := \lim_{z \to \eta_{i}(t), z \in \mathbb{H}_{t}^{\eta}} g_{\mathbb{H}_{t}^{\eta}}(z), \quad i = 1, \dots, N, \quad t \ge 0,
$$
\n
$$
g_{\mathbb{H}_{t}^{\eta}}: \text{ conformal}: \mathbb{H}_{t}^{\eta} \to \mathbb{H}.
$$
\n(3.42)

By the notation [\(3.37\)](#page-74-0), we put

<span id="page-75-4"></span><span id="page-75-3"></span><span id="page-75-2"></span><span id="page-75-1"></span>
$$
G^{g_{\mathbb{H}_t^{\eta}}}(z, w) := G_{\mathbb{H}}(g_{\mathbb{H}_t^{\eta}}(z), g_{\mathbb{H}_t^{\eta}}(w)), \quad t \ge 0
$$
\n(3.43)

and define the Dirichlet boundary GFF,  $g_{\mathbb{H}^n_{t}} H$  on  $\mathbb{H}^n_t$ <br> $(3.48)$ , W  $t, t \geq 0$ , so that its Green's function is given by [\(3.43\)](#page-75-1). We can prove the following.

# **Proposition 3.22** For  $t \geq 0$ ,

$$
\frac{dG^{\mathcal{G}_{\mathbb{H}_{t}^{\eta}}}(z,w)}{dt} = -4\sum_{i=1}^{N} \text{Im}\left(\frac{1}{\mathcal{G}_{\mathbb{H}_{t}^{\eta}}(z) - Y_{i}(t)}\right) \text{Im}\left(\frac{1}{\mathcal{G}_{\mathbb{H}_{t}^{\eta}}(w) - Y_{i}(t)}\right), \quad z, w \in \mathbb{H}_{t}^{\eta}.
$$
\n(3.44)

Proof. By [\(3.38\)](#page-74-1),

$$
\frac{dG^{\mathcal{G}_{\mathbb{H}_{t}^{\eta}}}(z,w)}{dt} = \frac{d}{dt} \Big[ \text{Re}\,\log(g_{\mathbb{H}_{t}^{\eta}}(z) - \overline{g_{\mathbb{H}_{t}^{\eta}}(w)}) - \text{Re}\,\log(g_{\mathbb{H}_{t}^{\eta}}(z) - g_{\mathbb{H}_{t}^{\eta}}(w)) \Big] \n= \text{Re}\,\frac{1}{g_{\mathbb{H}_{t}^{\eta}}(z) - \overline{g_{\mathbb{H}_{t}^{\eta}}(w)}} \left( \frac{dg_{\mathbb{H}_{t}^{\eta}}(z)}{dt} - \frac{dg_{\mathbb{H}_{t}^{\eta}}(w)}{dt} \right) \n- \text{Re}\,\frac{1}{g_{\mathbb{H}_{t}^{\eta}}(z) - g_{\mathbb{H}_{t}^{\eta}}(w)} \left( \frac{dg_{\mathbb{H}_{t}^{\eta}}(z)}{dt} - \frac{dg_{\mathbb{H}_{t}^{\eta}}(w)}{dt} \right), \quad t \ge 0.
$$

Since the multiple SLE [\(3.42\)](#page-75-2) gives

$$
\frac{dg_{\mathbb{H}_t^{\eta}}(z)}{dt} - \frac{d\overline{g_{\mathbb{H}_t^{\eta}}(w)}}{dt} = \sum_{i=1}^N \frac{2}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} - \sum_{i=1}^N \frac{2}{g_{\mathbb{H}_t^{\eta}}(w) - Y_i(t)}
$$
\n
$$
= -2(g_{\mathbb{H}_t^{\eta}}(z) - \overline{g_{\mathbb{H}_t^{\eta}}(w)}) \sum_{i=1}^N \frac{1}{(g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t))(\overline{g_{\mathbb{H}_t^{\eta}}(w)} - Y_i(t))},
$$

<span id="page-75-0"></span><sup>1</sup>This is proved by  $\boxed{34}$ . See Theorem [3.33](#page-84-0) shown in Section [3.3.5.](#page-82-0)

the above is written as

$$
\frac{dG^{\mathcal{G}_{\mathbb{H}_{t}^{n}}}(z,w)}{dt} = -2\sum_{i=1}^{N} \text{Re} \frac{1}{(g_{\mathbb{H}_{t}^{n}}(z) - Y_{i}(t))(\overline{g_{\mathbb{H}_{t}^{n}}(w)} - Y_{i}(t))} + 2\sum_{i=1}^{N} \text{Re} \frac{1}{(g_{\mathbb{H}_{t}^{n}}(z) - Y_{i}(t))(\overline{g_{\mathbb{H}_{t}^{n}}(w)} - Y_{i}(t))}, \quad t \ge 0.
$$

In general, the equality

<span id="page-76-2"></span><span id="page-76-1"></span> $\text{Re}\,\zeta\overline{w} - \text{Re}\,\zeta w = 2\text{Im}\,\zeta\,\text{Im}\,w$ 

holds for  $\zeta, w \in \mathbb{C}$ , and hence [\(3.44\)](#page-75-3) is proved.

For a given domain  $A \subset \mathbb{H}$ , we assume that supp  $(f)$  of  $f \in C_c^{\infty}(\mathbb{H})$  satisfies supp  $(f) \subset A$ . Then we define

$$
E_A^{g_{\mathbb{H}_t^{\eta}}}(f) := \int_{A \times A} f(z) G^{g_{\mathbb{H}_t^{\eta}}}(z, w) f(w) \mu(dz) \mu(dw).
$$
 (3.45)

By Proposition [3.22,](#page-75-4) we have

$$
\frac{dE_A^{\mathcal{G}_{\mathbb{H}_t^{\eta}}}(f)}{dt} = -\sum_{i=1}^N \left( \int_A \text{Im} \, \frac{2}{\mathcal{G}_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} f(z) \mu(dz) \right)^2, \tag{3.46}
$$

It implies that  $E_A^{\mathcal{G}_{\mathbb{H}_l^{\eta}}}(f)$  is non-increasing function of t.

## 3.3.2 Imaginary surface and extended GFF

Consider a simply connected domain  $D \subsetneq \mathbb{C}$  and write  $\mathcal{C}_c^{\infty}(D)$  for the space of real smooth functions on D with compact support. Assume  $h \in C_c^{\infty}(D)$  and consider a smooth vector field  $e^{\sqrt{-1}(h/\chi+\theta)}$ with parameters  $\chi, \theta \in \mathbb{R}$ . Then a flow line along this vector field,  $\eta : (0, \infty) \ni t \mapsto \eta(t) \in D$ , starting from  $\lim_{t\to 0} \eta(t) =: \eta(0) = x \in \partial D$  is defined (if exists) as the solution of the ordinary differential equation (ODE) [\[74,](#page-91-0) [54\]](#page-90-0)

<span id="page-76-0"></span>
$$
\frac{d\eta(t)}{dt} = e^{\sqrt{-1}\{h(\eta(t))/\chi + \theta\}}, \quad t \ge 0, \quad \eta(0) = x.
$$
 (3.47)

Let  $\widetilde{D} \subsetneq \mathbb{C}$  be another simply connected domain and consider a conformal map  $\varphi : \widetilde{D} \to D$ . Then we define the pull-back of the flow line  $\eta$  by  $\varphi$  as  $\tilde{\eta}(t) = (\varphi^{-1} \circ \eta)(t)$ . That is,  $\varphi(\tilde{\eta}(t)) = \eta(t)$ , and the derivatives with respect to  $t$  of the both sides of this equation gives

$$
\varphi'(\widetilde{\eta}(t))\frac{d\widetilde{\eta}(t)}{dt} = \frac{d\eta(t)}{dt}
$$

with the notation

$$
\varphi'(z) := \frac{d\varphi(z)}{dz}.
$$

We use the polar coordinate  $\varphi'(\cdot) = |\varphi'(\cdot)|e^{\sqrt{-1}\arg \varphi'(\cdot)}$ , where  $\arg \zeta$  of  $\zeta \in \mathbb{C}$  is a priori defined up to additive multiples of  $2\pi$ , and hence we have  $d\tilde{\eta}(t)/dt = e^{\sqrt{-1}\{(h\circ\varphi - \chi\arg\varphi')(\tilde{\eta}(t))/\chi + \theta\}}/|\varphi'(\tilde{\eta}(t))|, t \ge 0$ .

If we perform a time change  $t \to \tau = \tau(t)$  by putting  $t = \int_0^{\tau} ds / |\varphi'(\tilde{\eta}(s))|$  and  $\hat{\eta}(t) := \tilde{\eta}(\tau(t))$ , then the above equation becomes

$$
\frac{d\widehat{\eta}(t)}{dt} = e^{\sqrt{-1}\{(h\circ\varphi - \chi\arg\varphi')(\widehat{\eta}(t))/\chi + \theta\}}, \quad t \ge 0.
$$

Since a time change preserves the image of a flow line, we can identify h on D and  $h \circ \varphi - \chi \text{arg }\varphi'$ on  $\tilde{D} = \varphi^{-1}(D)$ . In [\[74,](#page-91-0) [54,](#page-90-0) [55,](#page-90-1) [56,](#page-90-2) [57\]](#page-90-3), such a flow line is considered also in the case that h is given by an instance of a GFF defined as follows.

**Definition 3.23** Let  $D \subsetneq \mathbb{C}$  be a simply connected domain and H be a Dirichlet boundary GFF on  $D$ . A GFF on  $D$  is a random distribution  $h$  of the form

<span id="page-77-2"></span>
$$
h = H + u,
$$

where  $u$  is a deterministic harmonic function on  $D$ .

Since a GFF is not function-valued, but it is a **distribution-valued random field** (see Remark 4.1 in Section [3.2\)](#page-68-0), the ODE in the form [\(3.47\)](#page-76-0) no longer makes sense mathematically in the classical sense. Using the theory of SLE, however, the notion of flow lines was generalized as follows [\[71,](#page-91-1) [72\]](#page-91-2).

Consider the collection

$$
\mathsf{S} := \left\{ (D, h) \middle| \begin{array}{c} D \subsetneq \mathbb{C} : \text{ simply connected} \\ h \colon \text{GFF on } D \end{array} \right\}.
$$

Fixing a parameter  $\chi \in \mathbb{R}$ , we define the following equivalence relation in S.

**Definition 3.24** Two pairs  $(D, h)$  and  $(\widetilde{D}, \widetilde{h}) \in S$  are  $\chi$ -equivalent if there exists a conformal map  $\varphi : \widetilde{D} \to D$  and

<span id="page-77-0"></span>
$$
\widetilde{h} \stackrel{\text{(law)}}{=} \varphi_* h - \chi \text{arg } \varphi'.
$$

In this case, we write  $(D, h) \sim (\widetilde{D}, \widetilde{h}).$ 

We call each element belonging to the equivalence class  $S / \sim$  an **imaginary surface** [\[54\]](#page-90-0) (or an AC surface [\[74\]](#page-91-0); (AC means a combination of an altimeter and a compass.) That is, in this equivalence class, a conformal map  $\varphi$  causes not only a coordinate change of a GFF as  $h \mapsto \varphi_* h$ associated with changing the domain of definition of the field as  $D \mapsto \varphi^{-1}(D)$ , but also an addition of a deterministic harmonic function  $-\chi \text{arg }\varphi'$  to the field. Notice that this definition depends on one parameter  $\chi \in \mathbb{R}$ .

#### <span id="page-77-3"></span>3.3.3 Complex-valued logarithmic potentials and local martingales

Assume that the driving process  $(Y(t))_{t>0}$  of the multiple SLE is given by [\(3.23\)](#page-67-1), that is

<span id="page-77-1"></span>
$$
dY_i(t) = \sqrt{\kappa}dB_i(t) + F_i(\mathbf{Y}(t))dt, \quad t \ge 0, \quad i = 1, \dots, N,
$$
\n(3.48)

Now for  $z \in \mathbb{C}, y_i \in \mathbb{R}, i = 1, \ldots, N$ , we consider a sum of **complex-valued logarithmic** potentials [\[32,](#page-88-1) [33,](#page-88-2) [34,](#page-88-0) [42\]](#page-89-0),  $\lambda$ r

$$
\Phi(z, y) := \sum_{i=1}^{N} \log(z - y_i), \tag{3.49}
$$

and consider a stochastic process  $(\Phi(g_{\mathbb{H}_t^n}(z), Y(t)))_{t\geq 0}$ . It should be noted that introduction of a complex-valued logarithmic potential in the form  $log(z - y)$  has been reported in the previous papers [\[16,](#page-87-0) [72,](#page-91-2) [29\]](#page-88-3) in order to establish the coupling between the original  $SLE_{\kappa}$  and the extended GFFs. In this sense, the present generalization for the multiple SLE/GFF coupling seems to be very straightforward. But this generalization is nontrivial. Moreover, it was shown by [\[42\]](#page-89-0) that this choice is the only possibility to obtain the multiple SLE/GFF coupling when  $\kappa \neq 4$ . See also [\[29\]](#page-88-3) for the consideration from the viewpoint of conformal field theory.

Since

$$
\frac{\partial}{\partial z}\log(z-x) = \frac{1}{z-x}, \quad \frac{\partial}{\partial x}\log(z-x) = -\frac{1}{z-x}, \quad \frac{\partial^2}{\partial x^2}\log(z-x) = -\frac{1}{(z-x)^2},
$$

Itô's formula [\(1.19\)](#page-7-0) with the multiple  $SLE_{\kappa}$  [\(3.22\)](#page-67-0) gives

$$
d \log(g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)) = \frac{1}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} \left[ dg_{\mathbb{H}_t^{\eta}}(z) - dY_i(t) \right]
$$
  
= 
$$
\frac{1}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} \left[ 2 \sum_{j=1}^N \frac{1}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} dt - \sqrt{\kappa} dB_i(t) - F_i(\mathbf{Y}(t)) dt \right]
$$
  
- 
$$
\frac{1}{2} \frac{1}{(g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t))^2} \kappa dt,
$$

and hence we have

$$
d\Phi(g_{\mathbb{H}_t^{\eta}}(z), \mathbf{Y}(t)) = -\sqrt{\kappa} \sum_{i=1}^N \frac{1}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} dB_i(t) + 2\left(\sum_{i=1}^N \frac{1}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)}\right)^2 dt - \sum_{i=1}^N \frac{1}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} F_i(\xi) dt - \frac{\kappa}{2} \sum_{i=1}^N \frac{1}{(\varphi_t(z) - Y_i(t))^2} dt.
$$
 (3.50)

The factor appearing in the second term of RHS is written as

$$
\left(\sum_{i=1}^{N} \frac{1}{g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{i}(t)}\right)^{2} = \sum_{i=1}^{N} \sum_{1 \leq j \leq N, j \neq i} \frac{1}{(g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{i}(t))(g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{j}(t))} + \sum_{i=1}^{N} \frac{1}{(g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{j}(t))^{2}}
$$

$$
= 2 \sum_{i=1}^{N} \frac{1}{g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{i}(t)} \sum_{\substack{1 \leq j \leq N, \\ j \neq i}} \frac{1}{Y_{i}(t) - Y_{j}(t)} + \sum_{i=1}^{N} \frac{1}{(g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{j}(t))^{2}}.
$$
(3.51)

By differentiate the multiple SLE  $(3.22)$  with respect to z, we have

<span id="page-78-0"></span>
$$
dg'_{\mathbb{H}_t^{\eta}}(z) = -2g'_{\mathbb{H}_t^{\eta}}(z) \sum_{i=1}^N \frac{1}{(g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t))^2} dt,
$$

and hence the equality

$$
d \log g'_{\mathbb{H}_t^\eta}(z) = -2 \sum_{i=1}^N \frac{1}{(g_{\mathbb{H}_t^\eta}(z) - Y_i(t))^2} dt
$$

is established. Therefore, if [\(3.50\)](#page-78-0) is multiplied by  $2/\sqrt{\kappa}$ , then we obtain

$$
\frac{2}{\sqrt{\kappa}} d\Phi(g_{\mathbb{H}_{t}^{\eta}}(z), \mathbf{Y}(t)) = -2 \sum_{i=1}^{N} \frac{1}{g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{i}(t)} dB_{i}(t) \n+ \frac{2}{\sqrt{\kappa}} \sum_{i=1}^{N} \frac{1}{g_{\mathbb{H}_{t}^{\eta}}(z) - Y_{i}(t)} \left[ 4 \sum_{\substack{1 \leq j \leq N, \\ j \neq i}} \frac{1}{Y_{i}(t) - Y_{j}(t)} - F_{i}(\mathbf{Y}(t)) \right] dt \n- \frac{1}{\sqrt{\kappa}} \left( 2 - \frac{\kappa}{2} \right) d \log g'_{\mathbb{H}_{t}^{\eta}}(z).
$$
\n(3.52)

Comparing the above calculation result with the  $\chi$ -equivalence

<span id="page-79-0"></span>
$$
H_{D_1} \stackrel{\text{(law)}}{=} \varphi_* H_{D_2} - \chi \text{arg}\,\varphi'
$$

defined by Definition [3.24,](#page-77-0) we find the following:

(i) if 
$$
\chi = \frac{2}{\sqrt{\kappa}} - \frac{\sqrt{\kappa}}{2}
$$
,

then the imaginary part of the last term in RHS of  $(3.52) = -d(\chi \arg g'_{\mathbb{H}_t} (z)),$ 

(ii) if 
$$
F_i(\mathbf{Y}) = 4 \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{1}{Y_i(t) - Y_j(t)}, \quad i = 1, ..., N,
$$
 (3.53)

then the second term in RHS of  $(3.52) = 0$ .  $(3.54)$ 

Remark 3.25 If we choose the drift term as [\(1.65\)](#page-17-0), the system of SDEs of the driving process of the multiple SLE [\(3.48\)](#page-77-1) is determined as

<span id="page-79-1"></span>
$$
dY_i(t) = \sqrt{\kappa} dB_i(t) + 4 \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{1}{Y_i(t) - Y_j(t)} dt, \quad t \ge 0, \quad i = 1, ..., N,
$$
 (3.55)

As we have already claimed that  $(\sqrt{\kappa}B(t))_{t\geq0}$   $\stackrel{\text{(law)}}{=} (B(\kappa t))_{t\geq0}$ , we perform a time change  $\kappa t \to t$ and define  $\mathbf{X}(t) = \mathbf{Y}(\kappa t), t \geq 0$ . Then we have the following set of SDEs for  $(\mathbf{X}(t))_{t \geq 0}$ ,

$$
dX_i(t) = dB_i(t) + \frac{4}{\kappa} \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{1}{X_i(t) - X_j(t)} dt
$$
  

$$
\iff dX_i(t) = dB_i(t) + \frac{\beta}{2} \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{1}{X_i(t) - X_j(t)} dt, \quad t \ge 0, \quad i = 1, ..., N,
$$
 (3.56)

with

$$
\beta = \frac{8}{\kappa}.\tag{3.57}
$$

Hence, we can say that the N-particle system  $(Y(t))_{t\geq 0}$  satisfying [\(3.55\)](#page-79-1) is a time change of the Dyson model with  $\beta = 8/\kappa$ ; DYS<sub>8/κ</sub>.

We introduce the following notation,

$$
\mathfrak{h}_t(\cdot) := -\frac{2}{\sqrt{\kappa}} \mathrm{Im} \,\Phi(g_{\mathbb{H}_t^{\eta}}(\cdot), \mathbf{Y}(t)) - \chi \mathrm{Im} \, \log g_{\mathbb{H}_t^{\eta}}'(\cdot)
$$

$$
= -\frac{2}{\sqrt{\kappa}} \sum_{i=1}^N \arg \left( g_{\mathbb{H}_t^{\eta}}(\cdot) - Y_i(t) \right) - \chi \arg g_{\mathbb{H}_t^{\eta}}'(\cdot). \tag{3.58}
$$

<span id="page-80-0"></span>The above observation implies the following.

Proposition 3.26 Assume  $(3.55)$ , that is, the driving process of the multiple SLE is given by DYS<sub>8/κ</sub>. If  $\chi = \frac{2}{\sqrt{2}}$  $\frac{1}{\sqrt{\kappa}}$  –  $\sqrt{\kappa}$  $\frac{1}{2}$ , then  $(\mathfrak{h}_t(\cdot))_{t\geq 0}$  is a continuous local martingale and satisfies

$$
d\mathfrak{h}_t(z) = 2 \sum_{i=1}^N \text{Im} \, \frac{1}{g_{\mathbb{H}_t^{\eta}}(z) - Y_i(t)} dB_i(t), \quad z \in \mathbb{H}_t^{\eta}, \quad t \in [0, \infty). \tag{3.59}
$$

**Remark 3.27** Between the central charge  $c = c_{\kappa}$  given by [\(3.21\)](#page-67-2) and  $\chi = \chi_{\kappa} := \frac{2}{\sqrt{\kappa}}$  $\frac{1}{\sqrt{\kappa}}$  –  $\sqrt{\kappa}$  $\frac{1}{2}$ , the following simple relation is established,

<span id="page-80-4"></span>
$$
c_{\kappa} = 1 - 6\chi_{\kappa}^2. \tag{3.60}
$$

By this formula, for the SLE parameter  $\kappa > 0$ ,  $c_{\kappa} \leq 1$ , and the maximum central charge  $c_{\kappa} = 1$  is realized if and only if  $\chi_{\kappa} = 0 \Longleftrightarrow \kappa = 4$ .

Comparing Proposition [3.26](#page-80-0) with [\(3.44\)](#page-75-3) in Proposition [3.22,](#page-75-4) we see that

$$
d\langle \mathfrak{h}.(z),\mathfrak{h}.(w)\rangle_t=-dG^{g_{\mathbb{H}_t^{\eta}}}(z,w),\quad z,w\in \mathbb{H}_t^{\eta},\quad t\in [0,\infty).
$$

Moreover, with [\(3.46\)](#page-76-1) we have the equality,

$$
d\langle\langle \mathfrak{h}, f \rangle, \langle \mathfrak{h}, f \rangle\rangle_t = -dE_A^{g_{\mathbb{H}_t^{\eta}}}(f) \quad \text{for } f \in \mathcal{C}_c^{\infty}(\mathbb{H}) \text{ with } \text{supp}(f) \subset A \subset \mathbb{H}. \tag{3.61}
$$

#### 3.3.4 Extended GFF-valued stochastic process and its stationarity

Put

<span id="page-80-3"></span><span id="page-80-1"></span>
$$
\chi = \frac{2}{\sqrt{\kappa}} - \frac{\sqrt{\kappa}}{2},
$$

and consider the following time-evolution of extended GFF,

$$
H_t := g_{\mathbb{H}_t^\eta} H + \mathfrak{h}_t, \quad t \ge 0. \tag{3.62}
$$

<span id="page-80-2"></span>Remark 3.28 By Definition [3.23,](#page-77-2) a GFF added by a harmonic function has been called an extended GFF. The additional harmonic function in [\(3.62\)](#page-80-1) is

$$
\arg\left(g_{\mathbb{H}_t^{\eta}}(\cdot) - Y_i(t)\right) = \text{Im}\,\log(g_{\mathbb{H}_t^{\eta}}(\cdot) - Y_i(t)),
$$



<span id="page-81-0"></span>Figure 9: Consider a domain  $A \subset \mathbb{H}$  such that there is a gap between A and the real axis R. For such A, we define  $\tau_A := \sup \left\{ t \geq 0 : A \subset \mathbb{H}_t^n \right\}$ t  $\geq 0$  and consider  $H_t$  during the time period  $t \in [0, \tau_A].$ 

which is not well-defined on the SLE curves  $\eta_i$ ,  $i = 1, \ldots, N$ . Hence,  $H_t$  can be regarded as an extended GFF not in the whole  $\mathbb{H}$ , but only on  $\mathbb{H}^{\eta}_{t}$  $t_t^{\eta}, t \geq 0$ , where we remember that  $\mathbb{H}_t^{\eta}$  $t''_t, t > 0$  is defined by [\(3.41\)](#page-75-5), that is

$$
\mathbb{H}_t^{\eta} := \text{the unbounded component of } \mathbb{H} \setminus \bigcup_{i=1}^N \eta_i(0, t], \quad t \ge 0.
$$

We consider a domain  $A \subset \mathbb{H}$  such that there is a gap between A and the real axis R. In this case, if we set  $v_A := \inf \{\text{Im } z : z \in A\}$ , then  $v_A \geq \exists \delta > 0$ . For such A, we define an  $(\mathcal{F}_t)_{t \geq 0}$ -stopping time,

<span id="page-81-1"></span>
$$
\tau_A := \sup \left\{ t \ge 0 : A \subset \mathbb{H}_t^{\eta} \right\}.
$$
\n(3.63)

Then,  $\tau_A > 0$  and during the time period  $t \in [0, \tau_A]$ , A is separated from any SLE curve and the multiple SLE  $g_{\mathbb{H}_t^n}$  is well defined in A, and hence [\(3.62\)](#page-80-1) can be regarded as an extended GFF-valued process. See Fig[.3.3.4.](#page-81-0) Note that test functions  $f \in C_c^{\infty}(\mathbb{H})$  should be considered with the condition  $supp(f) \subset A$ .

<span id="page-81-2"></span>Under the restrictions mentioned in Remark [3.28,](#page-80-2) the following Proposition asserts that a kind of stationarity is established if we couple the multiple SLE and GFF properly.

#### Proposition 3.29 Assume the following.

(A)  $(h_t)_{t>0}$  is a continuous local martingale and its quadratic covariation satisfies the equality,

$$
d\langle \mathfrak{h}_{\cdot}(z),\mathfrak{h}_{\cdot}(w)\rangle_{t} = -dG^{g_{\mathbb{H}_{t}^{\eta}}}(z,w), \quad z,w \in \mathbb{H}_{t}^{\eta}, \quad t \in [0,\infty).
$$
 (3.64)

Let A be an arbitrary domain in  $\mathbb H$  such that  $v_A \geq \frac{1}{\delta} > 0$ . With the  $(\mathcal F_t)_{t>0}$ -stopping time  $\tau_A$ given by [\(3.63\)](#page-81-1), put  $0 < T < \tau_A$ . Then for any  $f \in C_c^{\infty}(\mathbb{H})$  with supp  $(f) \subset A$ ,

<span id="page-82-1"></span>
$$
\langle H_0, f \rangle \stackrel{\text{(law)}}{=} \langle H_t, f \rangle, \quad t \in [0, T]. \tag{3.65}
$$

Proof. Introduce a real parameter  $\theta \in \mathbb{R}$  and consider the characteristic function for  $\langle H_t, f \rangle$ ,  $\mathbb{E}[e^{\sqrt{-1}\theta\langle H_t,f\rangle}].$  Here E denotes the expectation with respect to the joint probability law P of the multiple SLE and GFF.  $\mathfrak{h}_t(\cdot)$  is  $\mathcal{F}_t$ -measurable,

$$
\mathbb{E}\Big[e^{\sqrt{-1}\theta\langle H_t,f\rangle}\Big]=\mathbb{E}\Bigg[\mathbb{E}\Big[\exp(\sqrt{-1}\theta\langle g_{\mathbb{H}_t^\eta*}H,f\rangle)\Big|\mathcal{F}_t\Big]e^{\sqrt{-1}\theta\langle \mathfrak{h}_t,f\rangle}\Bigg].
$$

Since supp  $(f) \subset A$ , by the definition [\(3.45\)](#page-76-2) of  $E_A^{\mathcal{G}_{\mathbb{H}_t}^n}(f)$   $Var[\langle \mathcal{G}_{\mathbb{H}_t^n} H, f \rangle] = E_A^{\mathcal{G}_{\mathbb{H}_t}^n}(f)$ . Hence the formula  $(3.32)$  of Lemma [3.18,](#page-72-2) which was concluded from the Bochner–Minlos theorem, gives

$$
\mathbb{E}\Big[\exp(\sqrt{-1}\theta\langle g_{\mathbb{H}^\eta_t} H,f\rangle)\Big|\mathcal{F}_t\Big] = \exp\left(-\frac{\theta^2}{2}E_A^{\mathcal{G}_{\mathbb{H}^\eta_t}}(f)\right)
$$

Then we obtain

$$
\mathbb{E}\left[e^{\sqrt{-1}\theta\langle H_t,f\rangle}\right] = \mathbb{E}\left[\exp\left(-\frac{\theta^2}{2}E_A^{\mathcal{G}_{\mathbb{H}_t^{\eta}}}(f) + \sqrt{-1}\theta\langle \mathfrak{h}_t,f\rangle\right)\right].
$$
\n(3.66)

By Itô's formula,

$$
d \exp \left(-\frac{\theta^2}{2} E_A^{\mathcal{G}_{\mathbb{H}_t^{\eta}}}(f) + \sqrt{-1}\theta \langle \mathfrak{h}_t, f \rangle \right)
$$
  
= 
$$
\left\{ \sqrt{-1}\theta d \langle \mathfrak{h}_t, f \rangle - \frac{\theta^2}{2} \left( d E_A^{\mathcal{G}_{\mathbb{H}_t^{\eta}}}(f) + d \langle \langle \mathfrak{h}, f \rangle, \langle \mathfrak{h}, f \rangle \rangle_t \right) \right\} \exp \left(-\frac{\theta^2}{2} E_A^{\mathcal{G}_{\mathbb{H}_t^{\eta}}}(f) + \sqrt{-1}\theta \langle \mathfrak{h}_t, f \rangle \right).
$$

By the assumption (A), [\(3.61\)](#page-80-3) holds and we can conclude that  $\exp\left(-\frac{\theta^2}{2}E_A^{\mathcal{G}_{\mathbb{H}_t}^{\eta}}(f) + \sqrt{-1}\theta\langle\mathfrak{h}_t, f\rangle\right)$ is an  $\mathcal{F}_t$ -local martingale and that the characteristic function is given by its expectation. Therefore, the characteristic function is time independent and given by its initial value,

$$
\mathbb{E}\left[\exp\left(-\frac{\theta^2}{2}E_A^{g_{\mathbb{H}}}(f)+\sqrt{-1}\theta\langle \mathfrak{h}_0, f\rangle\right)\right]=\mathbb{E}\left[e^{\sqrt{-1}\theta\langle H_0, f\rangle}\right].
$$

The proof of the equality  $(3.65)$  is complete.

## <span id="page-82-0"></span>3.3.5 Multiple SLE/GFF coupling realized by  $\text{DYS}_{8/\kappa}$

**Definition 3.30 (multiple SLE/GFF coupling)** For any domain  $A \subset \mathbb{H}$  such that  $v_A \geq \frac{1}{2}\delta$ 0, define the  $(\mathcal{F}_t)_{t\geq0}$ -stopping time  $\tau_A$  by [\(3.63\)](#page-81-1) and assume  $0 < T < \tau_A$ . Given the driving process  $(Y(t))_{t\geq0}$  for the multiple SLE, assume the equality

$$
\langle H_0, f \rangle \stackrel{\text{(law)}}{=} \langle H_t, f \rangle, \quad t \in [0, T]
$$

for any  $f \in \mathcal{C}_c^{\infty}(\mathbb{H})$  with supp  $(f) \subset A$ . Then we say that the **multiple SLE/GFF coupling** is established.

The following is the main result in this section [\[32,](#page-88-1) [33,](#page-88-2) [34\]](#page-88-0).

#### Theorem 3.31 Assume

<span id="page-83-0"></span>
$$
\chi = \frac{2}{\sqrt{\kappa}} - \frac{\sqrt{\kappa}}{2}.\tag{3.67}
$$

Then, if and only if the driving process of multiple SLE  $(Y(t))_{t>0}$  is given by the solution of the SDEs

$$
dY_i(t) = \sqrt{\kappa} dB_i(t) + 4 \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{1}{Y_i(t) - Y_j(t)} dt, \quad t \ge 0, \quad i = 1, ..., N,
$$
 (3.68)

the multiple SLE/GFF coupling is established. In other words, under [\(3.67\)](#page-83-0), the driving process of multiple SLE is uniquely determined to be the Dyson model with

<span id="page-83-1"></span>
$$
\beta = \frac{8}{\kappa},\tag{3.69}
$$

 $DYS_{8/\kappa}$ , if the multiple SLE/GFF coupling holds.

*Proof.* Assume that  $(\mathfrak{h}_t)_{t\geq0}$  given by [\(3.58\)](#page-80-4) satisfies the assumption (A) in Proposition [3.29.](#page-81-2) Then by Propositions [3.26](#page-80-0) and [3.29,](#page-81-2) the present theorem will be proved if we can derive  $DYS_{8}/\kappa$  as  $(Y(t))_{t\geq0}$  under the conditions. Owing to the implicit function theorem, we can say that each  $Y_i(t), t \geq 0, i = 1, \ldots, N$  is a  $\mathcal{C}^{\infty}$  function of a continuous local martingale  $\mathfrak{h}_t(z_i)$ , multiple SLE  $g_t(z_j)$  solving [\(3.22\)](#page-67-0), and  $\log g'_t(z_j)$ ,  $j = 1, \ldots, N$ , which satisfy

$$
\frac{d}{dt}\log g_t'(z) = -\sum_{i=1}^N \frac{2}{(g_t(z) - Y_i(t))^2}, \quad z \in \mathbb{H}_t^{\eta}, \quad t \ge 0.
$$

Then by Itô's formula they are semi-martingales and given in the forms,

$$
Y_i(t) = M_i(t) + F_i(Y_t), \quad t \ge 0, \quad i = 1, ..., N,
$$
\n(3.70)

where  $M_i$ ,  $i = 1, ..., N$  denote the local martingale parts and  $F_i$ ,  $i = 1, ..., N$  do the finite-variation parts. Using these representations, as already seen in Section  $3.3.3$ , Itô's formula gives

$$
d\mathfrak{h}_t(z) = \text{Im} \sum_{i=1}^N \frac{1}{(g_t(z) - Y_i(t))^2} \left\{ \left( -\frac{4}{\sqrt{\kappa}} + 2\chi \right) dt + \frac{1}{\sqrt{\kappa}} d \langle M_i(\cdot), M_i(\cdot) \rangle_t \right\} + \frac{2}{\sqrt{\kappa}} \text{Im} \frac{1}{g_t(z) - Y_i(t)} \left( dF_i(\mathbf{Y}(t)) - \sum_{\substack{1 \le j \le N, \\ j \ne i}} \frac{4}{Y_i(t) - Y_j(t)} dt \right) + \frac{2}{\sqrt{\kappa}} \text{Im} \sum_{i=1}^N \frac{1}{g_t(z) - Y_i(t)} dM_i(t), \quad z \in \mathbb{H}_t^{\eta}, \quad t \ge 0.
$$
 (3.71)

By the assumption  $(A)$ , the first and the second lines in RHS should vanish. The each term in the first line represents the contribution from the pole of second order at  $g_t(z) = Y_i(t)$ ,  $i = 1, ..., N$ . Under [\(3.67\)](#page-83-0), vanishing of all these contribution implies

$$
d\langle M_i(\cdot), M_i(\cdot)\rangle_t = \kappa dt, \quad i = 1, \ldots, N.
$$

The each term in the second line represents the contribution from the pole of first order at  $g_t(z)$  $Y_i(t)$ ,  $i = 1, \ldots, N$ , and vanishing of all these contribution implies

$$
dF_i(\mathbf{Y}_t) = \sum_{\substack{1 \leq j \leq N, \\ j \neq i}} \frac{4}{Y_i(t) - Y_j(t)} dt, \quad t \geq 0, \quad i = 1, \dots, N.
$$

Moreover, the quadratic covariations of  $(\mathfrak{h}_t(\cdot))_{t\geq0}$  is calculated as

$$
d\langle \mathfrak{h}_{\cdot}(z),\mathfrak{h}_{\cdot}(w)\rangle_{t} = -dG^{g_{\mathbb{H}_{t}^{n}}}(z,w)
$$
  
+ 
$$
\frac{4}{\kappa} \sum_{\substack{1 \leq i,j \leq N, \\ i \neq j}} \text{Im}\left(\frac{1}{g_{t}(z) - Y_{i}(t)}\right) \text{Im}\left(\frac{1}{g_{t}(w) - Y_{j}(t)}\right) d\langle M_{i}(\cdot), M_{j}(\cdot)\rangle_{t},
$$

 $z, w \in \mathbb{H}^{\eta}_t$  $t_t^{\eta}, t \geq 0$ . By the assumption  $(\mathbf{A}),$ 

$$
d\langle M_i(\cdot), M_j(\cdot)\rangle_t = 0, \quad t \ge 0, \quad 1 \le i \ne j \le N.
$$

Hence the proof is complete.

Remark 3.32 Remember the relation [\(1.66\)](#page-17-1) in Section [1.5.1,](#page-16-0)

$$
\beta=D-1,
$$

and the relation [\(3.17\)](#page-66-0) in Section [3.1.2,](#page-63-0)

$$
\kappa = \frac{4}{D-1}.
$$

If we simply combine these two relations, we will have

<span id="page-84-0"></span>
$$
\beta = \frac{4}{\kappa}.
$$

The relation [\(3.69\)](#page-83-1) which establishes the multiple SLE/GFF coupling is different from such a simple-minded result. The multiple SLE driven by  $DYS_\beta$  with the true relation [\(3.69\)](#page-83-1),  $\beta = 8/\kappa$ , inherits many properties from the original  $\text{SLE}_\kappa$  with a single SLE curve. Actually, we have proved that our **multiple SLE**<sub>κ</sub> also shows **phase transitions** at  $\kappa_c = 4$  and  $\overline{\kappa}_c = 8$  [\[34\]](#page-88-0).

**Theorem 3.33** For each  $i = 1, ..., N$ , the limit  $\eta_i(t) = \lim_{\varepsilon \downarrow 0} g_{\mathbb{H}_t^{\eta}}(Y_i(t) + \sqrt{-1}\varepsilon)$  exists for all  $t \geq 0$ and  $\lim_{t\to\infty}\eta_i(t) = \infty$  with probability one. Moreover, the following **three phases** are observed.

- (a) If  $0 < \kappa \leq \kappa_c = 4$ ,  $\eta_i(0,\infty), i = 1,\ldots,N$  are simple disjoint curves such that  $\eta_i \subset \mathbb{H}, i =$  $1, \ldots, N$  with probability one.
- (b) If  $\kappa_c = 4 < \kappa < \overline{\kappa}_c = 8$ ,  $\eta_i$ ,  $i = 1, \ldots, N$  are continuous curves with probability one, and they intersect themselves and R with positive probability.
- (c) If  $\kappa = 8$ ,  $\eta_i$ ,  $i = 1, ..., N$  are space filling continuous curves with probability one.

# 3.4 Exercises 3

## 3.4.1 Exercise 3.1

Let  $\alpha \in (0,1)$  and

$$
\kappa = \kappa(\alpha) = \frac{4(1 - 2\alpha)^2}{\alpha(1 - \alpha)}.
$$
\n(3.72)

Then consider the case such that

$$
V(t) = \begin{cases} \sqrt{\kappa t}, & \text{if } \alpha \le 1/2, \\ -\sqrt{\kappa t}, & \text{if } \alpha > 1/2. \end{cases}
$$
 (3.73)

(1) Show that the inverse of  $g_t$  is solved as

$$
g_{\mathbb{H}_t^{\eta}}^{-1}(z) = \left(z + 2\sqrt{\frac{\alpha}{1-\alpha}}\sqrt{t}\right)^{1-\alpha} \left(z - 2\sqrt{\frac{1-\alpha}{\alpha}}\sqrt{t}\right)^{\alpha}, \quad z \in \mathbb{H}.
$$
 (3.74)

(2) Derive that the following formula for the slit,

$$
\eta(t) = g_{\mathbb{H}_t^{\eta}}^{-1}(V(t)) = 2\left(\frac{1-\alpha}{\alpha}\right)^{1/2-\alpha} e^{\sqrt{-1}\alpha\pi} t^{1/2}, \quad t \ge 0. \tag{3.75}
$$

The slit grows from the origin along a straight line in  $\mathbb H$  which makes an angle  $\alpha\pi$  with respect to the positive direction of the real axis. When  $\alpha = 1/2$ , this is reduced to Example [3.2.](#page-61-0)

#### 3.4.2 Exercise 3.2

Assume that  $D, D' \subsetneq \mathbb{C}$  are simply connected domains and let

$$
\varphi: D' \to D \quad \text{conformal map.} \tag{3.76}
$$

(1) We write the conformal map  $\varphi$  using the real functions (harmonic functions) u, v as

$$
\varphi(z) = u(x, y) + \sqrt{-1}v(x, y).
$$
 (3.77)

Using the Cauchy–Riemann equations,

<span id="page-85-0"></span>
$$
\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}, \quad \frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x},\tag{3.78}
$$

show that the Jacobian of  $\varphi$  is given by

$$
\frac{\partial(u,v)}{\partial(x,y)} = \frac{\partial u}{\partial x}\frac{\partial v}{\partial y} - \frac{\partial u}{\partial y}\frac{\partial v}{\partial x} = \left(\frac{\partial u}{\partial x}\right)^2 + \left(\frac{\partial u}{\partial y}\right)^2\tag{3.79}
$$

(2) Assume that  $f, g \in C_c^{\infty}(D)$  Using the chain rule of differentials and the Cauchy–Riemann equations [\(3.78\)](#page-85-0), show the equality,

$$
(\nabla f \circ \varphi)(z) \cdot (\nabla g \circ \varphi)(z) = \left(\frac{\partial f}{\partial u}\frac{\partial g}{\partial u} + \frac{\partial f}{\partial v}\frac{\partial g}{\partial v}\right) \left\{ \left(\frac{\partial u}{\partial x}\right)^2 + \left(\frac{\partial u}{\partial y}\right)^2 \right\}.
$$
 (3.80)

(3) Prove the equality,

$$
\int_{D} (\nabla f)(z) \cdot (\nabla g)(z) \mu(dz) = \int_{D'} (\nabla (f \circ \varphi))(z) \cdot (\nabla (g \circ \varphi))(z) \mu(dz)
$$
\n(3.81)

\nfor  $f, g \in \mathcal{C}^{\infty}(D)$ .

## 3.4.3 Exercise 3.3

Assume that

$$
GFF[e^{\sqrt{-1}\theta \langle H,f\rangle_{\nabla}}] = e^{-\theta^2 \|f\|_{\nabla}^2/2}, \quad \theta \in \mathbb{R}, \quad f \in W(D). \tag{3.82}
$$

(1) From [\(3.82\)](#page-86-0), derive the following equalities,

<span id="page-86-0"></span>E

$$
\text{Var}[\langle H, f \rangle_{\nabla}] := \mathbb{E}^{\text{GFF}} \left[ \langle H, f \rangle_{\nabla}^2 \right] = \|f\|_{\nabla}^2, \quad f \in W(D), \tag{3.83}
$$

$$
Cov[\langle H, f \rangle_{\nabla}, \langle H, g \rangle_{\nabla}] := \mathbb{E}^{\text{GFF}} \Big[ \langle H, f \rangle_{\nabla} \langle H, g \rangle_{\nabla} \Big] = \langle f, g \rangle_{\nabla}, \quad f, g \in W(D). \tag{3.84}
$$

(2) Let  $\varphi : D' \to D$  be a conformal map, and denote its pull-back as

$$
\varphi^* : W(D) \ni f \mapsto f \circ \varphi \in W(D'). \tag{3.85}
$$

Then the equality

$$
\int_{D} (\nabla f)(z) \cdot (\nabla g)(z) \mu(dz) = \int_{D'} (\nabla (f \circ \varphi))(z) \cdot (\nabla (g \circ \varphi))(z) \mu(dz) \tag{3.86}
$$

in Lemma [3.19](#page-72-3) is simply written as

$$
\langle \varphi^* f, \varphi^* g \rangle_{\nabla} = \langle f, g \rangle_{\nabla}.
$$
\n(3.87)

Prove the conformal invariance of GFF,

$$
Cov[\langle \varphi_* H, f \rangle_{\nabla}, \langle \varphi_* H, g \rangle_{\nabla}] = Cov[\langle H, f \rangle_{\nabla}, \langle H, g \rangle_{\nabla}].
$$
\n(3.88)

Acknowledgements This manuscript was prepared for the lectures given in the 4th ZiF Summer School 'Randomness in Physics and Mathematics: From Integrable Probability to Disordered Systems' held at ZiF–Center for Interdisciplinary Research, Bielefeld University, Germany, from 1st to 13th August 2022, which was organized by Gernot Akemann and Friedrich Götze. The present author would like to thank Gernot and Friedrich for kind hospitality and well-organization of the summer school. The lectures were based on the joint work with Hideki Tanemura (Keio University), Tomoyuki Shirai (Kyushu University) and Shinji Koshida (Aalto University). The author is grateful to them for fruitful collaborations. The present study has been supported by the Grant-in-Aid for Scientific Research (C) (No.19K03674), (B) (No.18H01124), (S) (No.16H06338), and (A) (No.21H04432) of Japan Society for the Promotion of Science (JSPS). This work is also supported by JSPS Grant-in-Aid for Transformative Research Areas (A) JP22H05105.

# References

- [1] G. Akemann, J. Baik, P. Di Francesco: The Oxford Handbook of Random Matrix Theory, Oxford University Press, Oxford (2011).
- [2] G. W. Anderson, A. Guionnet, O. Zeitouni: An Introduction to Random Matrices, Cambridge University Press, Cambridge (2010).
- [3] A. Arai, Functional Integral Methods in Quantum Mathematical Physics, (in Japanese), Kyoritsu Shuppan, Tokyo, 2010.
- [4] M. Bauer, D. Bernard:  $SLE_{\kappa}$  growth processes and conformal field theories, Phys. Lett. B 543 (2002) 135–138.
- [5] M. Bauer, D. Bernard, K. Kytölä: Multiple Schramm–Loewner evolutions and statistical mechanics martingales, J. Stat. Phys. 120 (2005) 1125–1163.
- [6] V. Beffara: The dimension of the SLE curves, Ann. Probab. 36 (2008) 1421–1452.
- [7] S. R. Bell: Simplicity of the Bergman, Szegő and Poisson kernel functions. Math. Res. Lett. 2 (1995) 267–277.
- [8] S. R. Bell: The Cauchy Transform, Potential Theory and Conformal Mapping, 2nd ed. CRC Press, Boca Raton, FL (2016)
- [9] S. Bergman: The Kernel Function and Conformal Mapping, 2nd ed. Amer. Math. Soc., Providence, RI (1970).
- [10] A. Borodin, E. M. Rains: Eynard-Mehta theorem, Schur process, and their Pfaffian analog. J. Stat. Phys. 121 (2005) 291–317.
- [11] J. Cardy: Stochastic Loewner evolution and Dyson's circular ensemble. J. Phys. A Math. Gen. 36 (2003) L379-L386; Corrigendum. ibid., 12343.
- [12] D. Chelkak, H. Duminil-Copin, C. Hongler, A. Kemppainen, S. Smirnov: Convergence of Ising interfaces to Schramm's SLE curves, C. R. Acad. Sci. Paris, Sér. I Math. 352 (2014) 157-161.
- [13] D. Chelkak, S. Smirnov: Universality in the 2D Ising model and conformal invariance of fermionic observables, Inv. Math. 189 (2012) 515–580.
- [14] S. Cooper: The development of elliptic functions according to Ramanujan and Venkatachaliengar. Res. Lett. Inf. Math. Sci. 1 (2000) 65–78 (2000); available at https://mro.massey.ac.nz/handle/10179/4390
- [15] J. Dubédat: Commutative relations for Schramm–Loewner evolutions, Commun. Pure Appl. Math. 60 (2007) 1792–1847.
- <span id="page-87-0"></span>[16] J. Dubédat: SLE and the free field: partition functions and couplings, J. Amer. Math. Soc., 22 (2009) 995–1054.
- [17] H. Duminil-Copin: Parafermionic observables and their applications to planar statistical physics models, Ensaios Math. vol. 25 (2013), pp. 1-371.
- [18] H. Duminil-Copin, S. Smirnov: Conformal invariance of lattice models, Clay Mathematics Proceedings, vol. 15 (2012), pp.213–276.
- [19] F. J. Dyson: A Brownian-motion model for the eigenvalues of a random matrix, J. Math. Phys. 3 (1962) 1191–1198.
- [20] A. Edelman, E. Kostlan: How many zeros of a random polynomial are real? Bull. Amer. Math. Soc. (N.S.) 32 (1) 1–37 (1995) 1–17.
- [21] P. J. Forrester: Log-gases and Random Matrices, London Math. Soc. Monographs, Princeton University Press, Princeton (2010).
- [22] G. Gasper, M. Rahman: Basic Hypergeometric Series, 2nd ed., Cambridge University Press, Cambridge (2004).
- [23] P. Graczyk, J. Małecki: Multidimensional Yamada–Watanabe theorem and its applications to particle systems, J. Math. Phys. 54 (2013) 021503/1–15.
- [24] K. Graham: On multiple Schramm–Loewner evolutions, J. Stat. Mech. Theory Exp. 2007 (2007) P03008.
- [25] T. Hida: Brownian Motion, Application of Mathematics, vol.11, Springer, Heidelberg (1980).
- [26] J. B. Hough, M. Krishnapur, Y. Peres, B. Virág: *Zeros of Gaussian Analytic Functions and* Determinantal Point Processes, University Lecture Series, Vol. 51, Amer. Math. Soc., Providence, RI (2009).
- [27] Y. Kajihara, M. Noumi: Multiple elliptic hypergeometric series. An approach from the Cauchy determinant. Indag. Math. (N.S.) 14 (2003) 395–421.
- [28] O. Kallenberg: Random Measures, Theory and Applications, Probability Theory and Stochastic Modelling Vol.77, Springer, Switzerland (2017).
- <span id="page-88-3"></span>[29] N.-G. Kang, N. G. Makarov: *Gaussian Free Field and Conformal Field Theory*, Astérisque **353**, Société Mathématique de France, Paris (2013).
- [30] M. Katori: Determinantal martingales and noncolliding diffusion processes, Stoch. Proc. Appl. 124 (2014) 3724–3768.
- <span id="page-88-1"></span>[31] M. Katori: Bessel Processes, Schramm–Loewner Evolution, and the Dyson Model, Springer-Briefs in Mathematical Physics 11, Springer, Tokyo (2016).
- [32] M. Katori, S. Koshida: Conformal welding problem, flow line problem, and multiple Schramm– Loewner evolution, J. Math. Phys. 61 (2020) 083301/1–25.
- <span id="page-88-2"></span>[33] M. Katori, S. Koshida: Gaussian free fields coupled with multiple SLEs driven by stochastic log-gases, Advanced Study in Pure Mathematics 87 (2021) 315–340.
- <span id="page-88-0"></span>[34] M. Katori, S. Koshida: Three phases of multiple SLE driven by non-colliding Dyson's Brownian motions, J. Phys. A: Math. Theor. 54 (2021) 325002 (19 pages).
- [35] M. Katori, T. Shirai: Partial isometries, duality, and determinantal point processes, Random Matrices: Theory and Applications 11 (2022) 2250025 (70 pages).
- [36] M. Katori, T. Shirai: Zeros of the i.i.d. Gaussian Laurent series on an annulus: weighted Szegő kernels and permanental-determinantal point processes, Commun. Math. Phys. 392 (2022) 1099–1151.
- [37] M. Katori, H. Tanemura: Symmetry of matrix-valued stochastic processes and noncolliding diffusion particle systems, J. Math. Phys. 45 (2004) 3058–3085.
- [38] M. Katori, H. Tanemura: Noncolliding Brownian motion and determinantal processes, J. Stat. Phys. 129 (2007) 1233–1277.
- [39] M. Katori, H. Tanemura: Non-equilibrium dynamics of Dyson's model with an infinite number of particles, Commun. Math. Phys. 293 (2010) 469–497.
- [40] A. Kemppainen: Schramm–Loewner Evolution, Springer Briefs in Mathematical Physics 24, Springer (2017).
- <span id="page-89-0"></span>[41] T. H. Koornwinder: On the equivalence of two fundamental theta identities. Anal. Appl. (Singap.) 12 (2014) 711–725.
- [42] S. Koshida: Multiple backward Schramm–Loewner evolution and coupling with Gaussian free field, Lett. Math. Phys. 111 (2021) 30 (41 pages).
- [43] S. Koshida: private communication.
- [44] M. J. Kozdron, G. F. Lawler: The configurational measure on mutually avoiding SLE paths, In Universality and renormalization, Fields Inst. Commun., Vol. 50, Amer. Math. Soc., Providence, RI (2007), pp.199–224.
- [45] C. Krattenthaler: Advanced determinant calculus: a complement. Linear Algebra Appl. 411 (2005) 68–166.
- [46] P. P. Kufarev, V. V. Sobolev, L. V. Sporyševa: A certain method of investigation of extremal problems for functions that are univalent in the half-plane (in Russian), Trudy Tomsk. Gos. Univ. Ser. Meh.-Mat., 200 (1968) 142–164.
- [47] G. F. Lawler: Conformally Invariant Processes in the Plane, Amer. Math. Soc., Providence, RI (2005).
- [48] G. Lawler, O. Schramm, W. Werner: Conformal invariance of planar loop-erased random walks and uniform spanning trees, Ann. Probab. 32 (2004) 939–995.
- [49] E. H. Lieb: Proofs of some conjectures on permanents. J. Math. Mech. 16 (1966) 127–134.
- [50] K. Löwner: Untersuchungen über schlichte konforme Abbildungen des Einheitskreises I, Math. Ann. 89 (1923) 103–121.
- [51] M. Marcus, H. Minc: A Survey of Matrix Theory and Matrix Inequalities, Dover Publications, New York (1992).
- [52] S. McCullough, L. -C. Shen: On the Szegő kernel of an annulus. Proc. Amer. Math. Soc. 121 (1994) 1111–1121.
- [53] M. L. Mehta: Random Matrices, 3rd edn, Pure and Applied Mathematics (Amsterdam), Vol. 142 Elsevier/Academic Press, Amsterdam (2004).
- <span id="page-90-1"></span><span id="page-90-0"></span>[54] J. Miller, S. Sheffield: Imaginary geometry I : Interacting SLEs, Probab.Theory Relat. Fields 164 (2016) 553–705.
- <span id="page-90-2"></span>[55] J. Miller, S. Sheffield: Imaginary geometry II : Reversibility of  $\text{SLE}_\kappa(\rho_1, \rho_2)$  for  $\kappa \in (0, 4)$ , Ann. Probab. 44 (2016) 1647–1722.
- <span id="page-90-3"></span>[56] J. Miller, S. Sheffield: Imaginary geometry III : Reversibility of  $SLE_{\kappa}$  for  $\kappa \in (4,8)$ , Ann. Math. **184** (2016) 455–486.
- [57] J. Miller, S. Sheffield: Imaginary geometry IV : Interior rays, whole-plane reversibility, and space-filling trees, Probab.Theory Relat. Fields 169 (2017) 729–869.
- [58] H. Minc: Permanent, Encyclopedia of Mathematics and its Applications, Vol. 6, Addison– Wesley, MA (1978).
- [59] F. W. J. Olver, D. W. Lozier, R. F. Boisvert, C. W. Clark (eds.) : NIST Handbook of Mathematical Functions, U.S. Department of Commerce, National Institute of Standards and Technology, Washington, DC/ Cambridge University Press, Cambridge (2010); available at <http://dlmf.nist.gov>
- [60] H. Osada: Infinite-dimensional stochastic differential equations related to random matrices, Probab. Theory Relat. Fields 153 (2012) 471–509.
- [61] H. Osada: Interacting Brownian motions in infinite dimensions with logarithmic potentials, Ann. Probab. 41 (2013) 1–49.
- [62] H. Osada, H. Tanemura: Infinite-dimensional stochastic differential equations and tail  $\sigma$ -fields, Probab. Theory Relat. Fields 177 (2020) 1137–1242.
- [63] Y. Peres, B. Virág: Zeros of the i.i.d. Gaussian power series. A conformally invariant determinantal process. Acta Math. 194 (2005) 1–35.
- [64] D. Revuz, M. Yor: Continuous Martingales and Brownian Motion. 3rd edition, Springer, Berlin (1999)
- [65] L. C. G. Rogers, Z. Shi: Interacting Brownian particles and the Wigner law, Probab. Theory Relat. Fields 95 (1993) 555–570.
- [66] S. Rohde, O. Schramm: Basic properties of SLE, Ann. Math. 161 (2005) 883–924.
- [67] H. Rosengren, M. Schlosser: Elliptic determinant evaluations and the Macdonald identities for affine root systems. Compositio Math. 142 (2006) 937–961.
- [68] O. Roth, S. Schleissinger: The Schramm–Loewner equation for multiple slits, J. Anal. Math. 131 (2017) 73–99.
- [69] O. Schramm: Scaling limits of loop-erased random walks and uniform spanning trees, Israel J. Math. 118 (2000) 221–288.
- [70] O. Schramm, S. Sheffield: The harmonic explorer and its convergence to SLE(4), Ann. Probab. 33 (2005) 2127–2148.
- <span id="page-91-2"></span><span id="page-91-1"></span>[71] O. Schramm, S. Sheffield: Contour lines of the two-dimensional discrete Gaussian free field, Acta Math. 202 (2009) 21–137.
- [72] O. Schramm, S. Sheffield: A contour line of the continuum Gaussian free field, Probab. Theory Relat. Fields 157 (2013) 47–80.
- <span id="page-91-0"></span>[73] S. Sheffield: Gaussian free fields for mathematicians, Probab. Theory Relat. Fields 139 (2007) 521–541.
- [74] S. Sheffield: Conformal weldings of random surfaces: SLE and the quantum gravity zipper, Ann. Probab. 44 (2016) 3474–3545.
- [75] T. Shirai: Limit theorems for random analytic functions and their zeros. In: Functions in Number Theory and Their Probabilistic Aspects – Kyoto 2010, RIMS Kôkyûroku Bessatsu 34 (2012) 335–359.
- [76] T. Shirai, Y. Takahashi: Fermion process and Fredholm determinant. In: Begehr, H. G. W., Gilbert, R. P., Kajiwara, J. (eds.), Proceedings of the Second ISAAC Congress, Vol. 1, pp. 15–23, Kluwer Academic Publishers, Dordrecht (2000)
- [77] T. Shirai, Y. Takahashi: Random point fields associated with certain Fredholm determinants I: fermion, Poisson and boson point process. J. Funct. Anal. 205 (2003) 414–463.
- [78] T. Shirai, Y. Takahashi: Random point fields associated with certain Fredholm determinants II: fermion shifts and their ergodic and Gibbs properties. Ann. Probab. 31 (2003) 1533–1564.
- [79] S. Smirnov: Critical percolation in the plane: conformal invariance, Cardy's formula, scaling limits, C. R. Acad. Sci. Paris, Sér. I Math. **333** (2001) 239–244.
- [80] S. Smirnov: Conformal invariance in random cluster models. I. Holomorphic fermions in the Ising model, Ann. Math. 172 (2010) 1435–1467.
- [81] A. Soshnikov, Determinantal random point fields. Russian Math. Surveys 55 (2000) 923–975.
- [82] T. J. Tegtmeyer, A. D. Thomas: The Ahlfors map and Szegő kernel for an annulus. Rocky Mountain J. Math. 2 (1999) 709–723.
- [83] L. -C. Tsai: Infinite dimensional stochastic differential equations for Dyson's model, Probab. Theory Relat. Fields 166 (2016) 801–850.
- [84] K. Venkatachaliengar, S. Cooper (ed): Development of Elliptic Functions According to Ramanujan, Monograph in Number Theory, Vol.6, World Scientific, Singapore (2012).